

# TB-LMTO METHOD FOR EMBEDDED CLUSTER

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# MOTIVATION

- electronic and magnetic structure of impurities, finite clusters, complexes, etc.
- formation energies of impurities
- formation energies of complexes
- structural stability: effective interatomic interactions (Ising type Hamiltonian)
- effective exchange interactions (Heisenberg Hamiltonian)
- parameters of Bruno, Zingales, and Wang model

# OUTLINE

- TB-LMTO method
- Embedding of a cluster
- Electrostatic energy and potential
- Total energy of a cluster
- Applications
  - magnetism of non-magnetic impurities in oxides
  - impurity formation energy
  - parameters of BZW model
- Conclusions and outlook

## TB-LMTO method

$$H = C + \sqrt{\Delta}S\sqrt{\Delta}, \quad G(z) = (z - H)^{-1} = \lambda(z) + \mu(z)g(z)\mu(z)$$

$$g(z) = (P(z) - S)^{-1}, \quad \lambda(z) = \frac{\gamma - \alpha}{\Delta + (\gamma - \alpha)(z - C)}$$

$$\mu(z) = \frac{\sqrt{\Delta}}{\Delta + (\gamma - \alpha)(z - C)}, \quad P(z) = \frac{z - C}{\Delta + (\gamma - \alpha)(z - C)}$$

$S$  ... matrix of screened structure constants

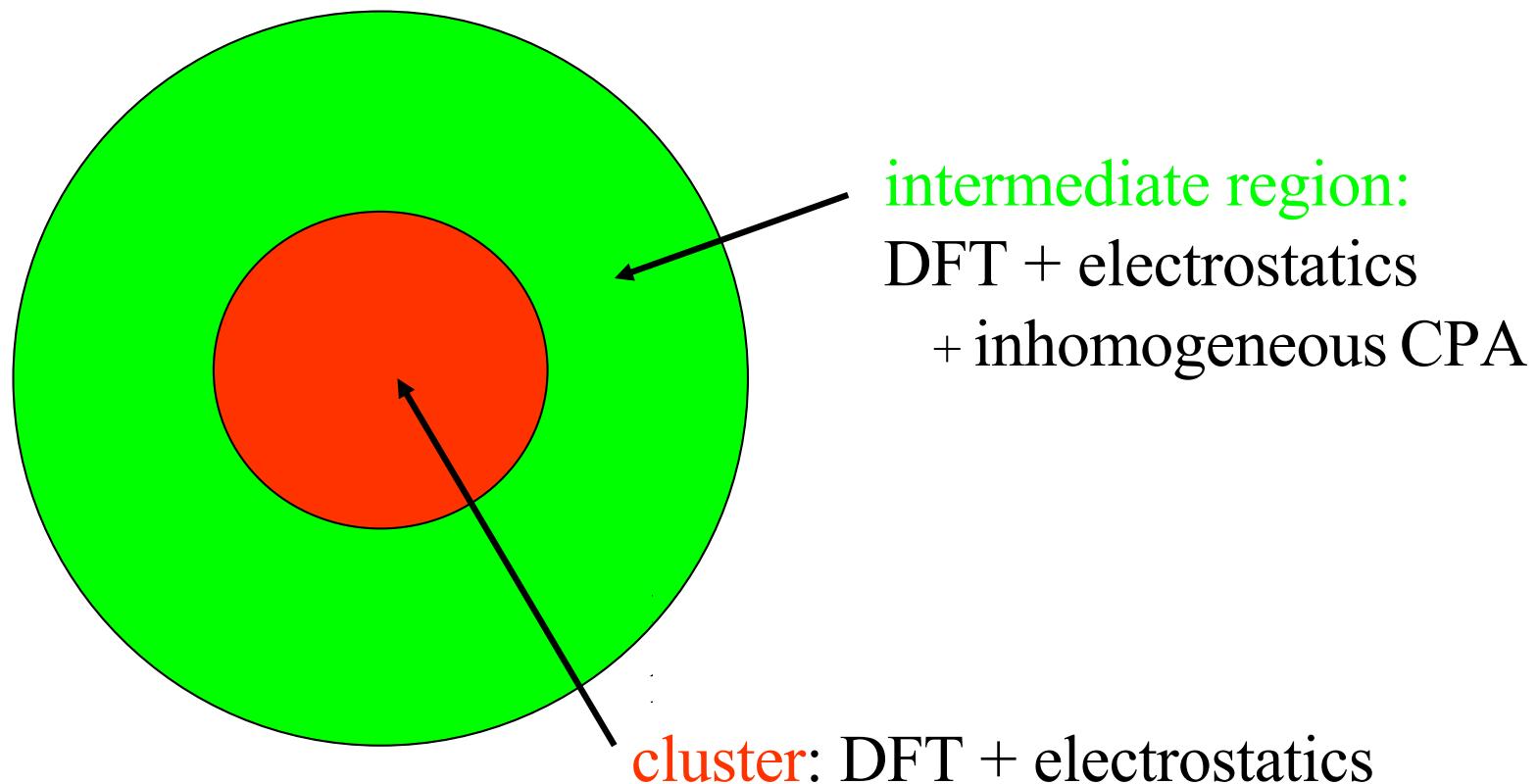
$C, \Delta, \gamma$  ... potential parameters, site-diagonal matrices

$P(z)$  ... potential function, site-diagonal matrix

$G(z)$  ... GF,  $g(z)$  ... auxiliary GF

# EMBEDDING

reference system: fixed



reference system: infinite homogeneous crystal, superscript 0  
 cluster  $C \dots \Pi_C$ , transition region  $T \dots \Pi_T$   
 inhomogeneous region  $I = C \cup T \dots \Pi_I = \Pi_C + \Pi_T$   
 homogeneous region  $H \dots \Pi_H$

$$g^0 = (P^0 - S)^{-1}, \quad g_{II}^0 = \frac{\Pi_I}{P_{II}^0 - S_{II} - S_{IH} \frac{H}{P^0 - S} S_{HI}}$$

$$g = (P - S)^{-1}, \quad g_{II} = \frac{\Pi_I}{P_{II} - S_{II} - S_{IH} \frac{H}{P - S} S_{HI}}$$

embedding potential:  $\Gamma_{II} = S_{II} - S_{IH} \frac{H}{P^0 - S} S_{HI} = \frac{\Pi_{II}}{P_{II}^0 - \frac{\Pi_I}{g_{II}^0}}$

(note that  $P_{HH} = P_{HH}^0$ )

# Madelung potential and Madelung energy

Madelung matrix

$$M_{ij} = \begin{cases} e^2/R_{ij} & \text{for } i \neq j \\ 0 & \text{for } i = j \end{cases}$$

reference system

$$v_i^0 = \sum_j M_{ij} q_j^0, \quad E^0 = \frac{1}{2} \sum_{ij} M_{ij} q_i^0 q_j^0 = \frac{1}{2} \sum_i v_i^0 q_i^0$$

system with embedded cluster

$$\varepsilon_i = q_i - q_i^0, \quad \varepsilon_i \neq 0 \quad \text{only for } i \in I,$$

$$v_i = v_i^0 + \sum_j^I M_{ij} \varepsilon_j, \quad \varphi_i = \sum_j^I M_{ij} \varepsilon_j,$$

$$E^{\text{Mad}} = \frac{1}{2} \sum_{ij} M_{ij} q_i q_j = E^0 + \frac{1}{2} \sum_i \varepsilon_i (2v_i^0 + \varphi_i)$$

# TOTAL ENERGY

$$E^{\text{tot}} = E^{\text{band}} + E^{\text{core}} + E^{\text{dc}} + E^{\text{xc}} + E^{\text{Mad}}$$

$$n(E) = -\frac{1}{\pi} \operatorname{Im} \operatorname{Tr} G(E + i0), \quad G(z) = (z - H)^{-1}$$

$$E^{\text{band}} = \int^{E_F} dE E n(E) = \frac{1}{2\pi i} \int_{C_R} dz z \operatorname{Tr} G(z)$$

$$C_R = \{z : z = E_F - R + Re^{i\phi}, \phi \in (0, 2\pi)\}$$

$H_0$  ... reference medium,  $V$  ... perturbation by cluster

$$G_0(z) = (z - H_0)^{-1}, \quad G(z) = (z - H_0 - V)^{-1}$$

although  $V$  is limited to a small subspace, it changes  $G(z)$  in the whole space  
and therefore it need not be sufficient to calculate the DOS on cluster atoms

# Tr log G

$$E^{\text{band}} = E_F N_{el} + \frac{1}{2\pi i} \int_{C_R} dz z \operatorname{Tr} \log(z - H)$$

two subspaces:  $\Pi_I$  ... cluster,  $\Pi_H$  ... rest of the system,  $\Pi_I + \Pi_H = 1$

$$\operatorname{Tr} \log(z - H) = \operatorname{Tr}_I \log \left( z\Pi_I - H_{II} - H_{IH} \frac{\Pi_H}{z - H} H_{HI} \right) + \operatorname{Tr}_H \log(z - H_{HH})$$

terms  $\operatorname{Tr}_H \log(z - H_{HH})$  are identical for reference system and for the cluster and thus will mutually subtract

two problems:

1. logarithm is a multivalued function
  2. potential function in the TB-LMTO has poles
- use the trick with  $\operatorname{Tr} \log G(z)$
  - calculate  $E_I^{\text{band}} = \int^{E_F} dE E n_I(E), \quad n_I(E) = -\frac{1}{\pi} \operatorname{Im} \operatorname{Tr}_I G(E + i0)$

# REMARKS

- reference system:
  - ideal crystal
  - disordered alloy
  - semi-infinite solid
  - empty spheres
- short-ranged structure constants needed  $\Rightarrow$  TB-LMTO
- calculation proceeds in two steps:
  - embedding potential  $\Gamma_{II}$  and Madelung potential of reference medium for inhomogeneous region
  - selfconsistency in the inhomogeneous region
- convergence of results with size of cluster important

# Magnetism of non-magnetic impurities in non-magnetic oxides

magnetism without d-electrons, or  $d^0$ –magnetism

- vacancies at cation sites in  $\text{TiO}_2$ ,  $\text{ZrO}_2$ ,  $\text{HfO}_2$  can induce magnetic moments on neighboring oxygen atoms
- formation energy of such vacancies is very high
- substitution of TM atom by an atom of alkali metal Li, Na, K, Rb ... or alkaline earths Be, Mg, Ca, ...

# Results for ZrO<sub>2</sub>

fluorite structure, cluster 59 sites

shell	n	atom
0	1	Zr, Vac, K, etc.
1	8	O ... O(1)
2	6	empty
3	12	Zr
4	24	O ... O(2)
5	8	E

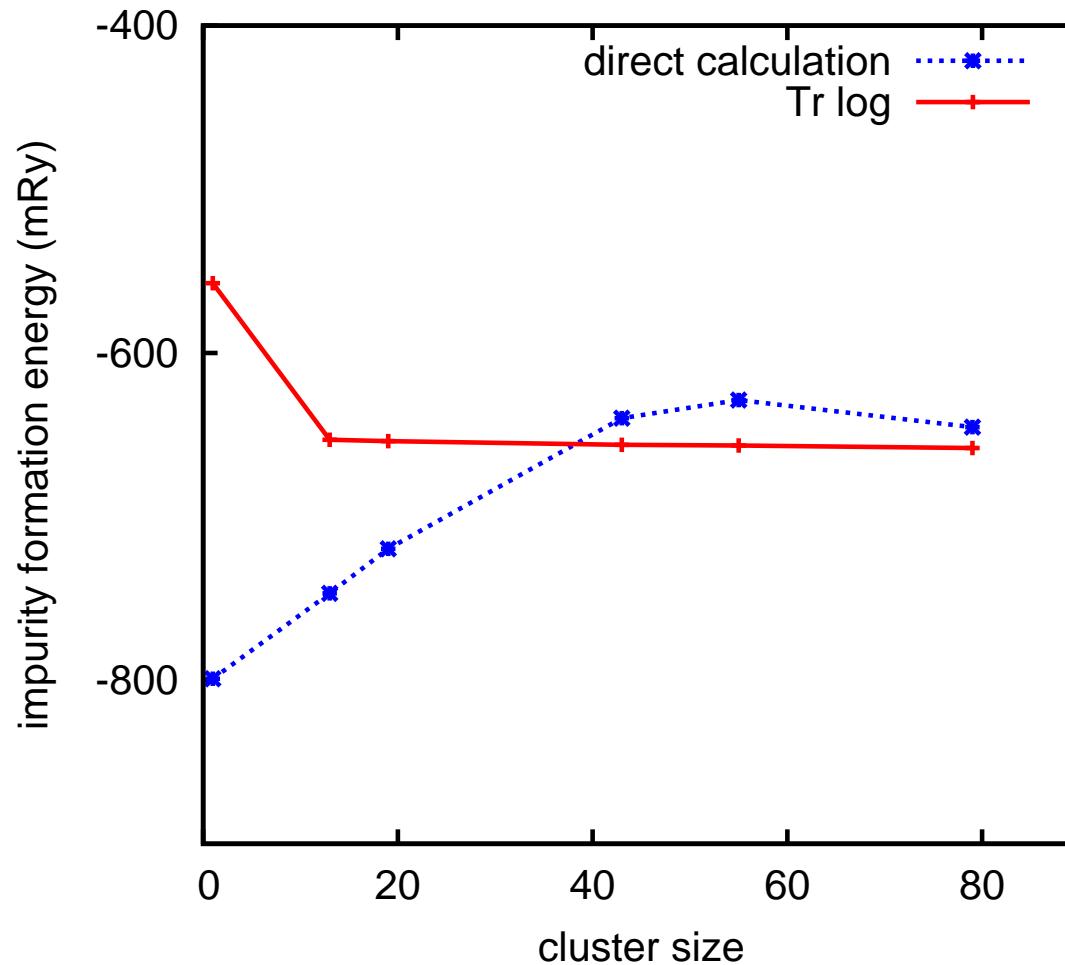
# Results for ZrO<sub>2</sub>

comparison with FPLAPW calculations for XZr<sub>7</sub>O<sub>16</sub> (12.5 % X) and XZr<sub>31</sub>O<sub>64</sub>(3.125 %), where X=Vac, K, etc.

cluster	12.5 %		3.125 %				
	$M[\mu_B]$	$M[\mu_B]$	$M[\mu_B]$	$M[\mu_B]$			
Vac	0.005	0.00	0.00	K	0.158	0.10	0.08
O(1)	0.421	0.40	0.41	O(1)	0.244	0.30	0.24
O(2)	0.027	0.06	0.03	O(2)	0.023	0.05	0.03
total	3.927	3.86	3.99	total	2.616	3.00	2.80

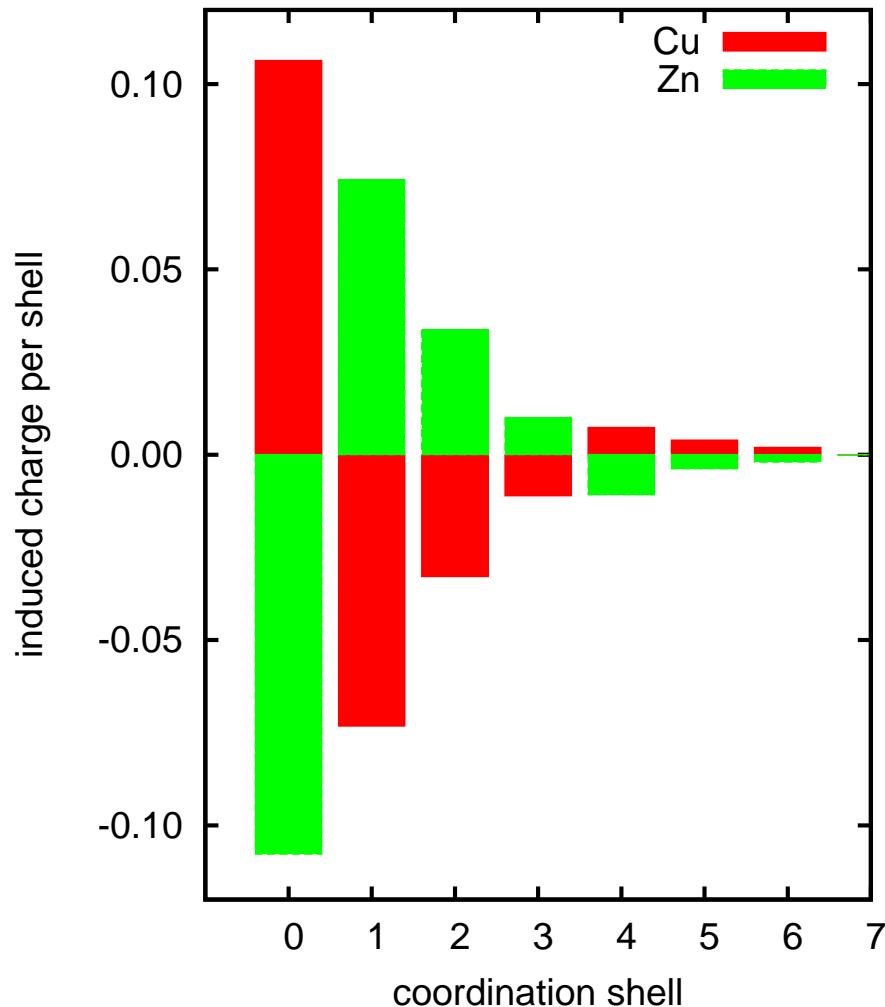
# Zn impurity in fcc-Cu

impurity formation energy  $\Delta E = \left( E[Cu_{N-1}Zn] + \epsilon[Cu] \right) - \left( E[Cu_N] + \epsilon[Zn] \right)$



# Impurity-induced charges

charge induced in the neighborhood of Cu and Zn atom in bcc-Cu<sub>50</sub>Zn<sub>50</sub> alloy



# Model of Bruno, Zingales, and Wang

E. Bruno, L. Zingales, and Y. Wang: PRL **91** (2003), 166401.

$$E(\{q_i\}) = E^{\text{loc}}(\{q_i\}) + E^{\text{Mad}}(\{q_i\}) = \sum_i \frac{a_i}{2} (q_i - b_i)^2 + \frac{1}{2} \sum_{i,j} M_{ij} q_i q_j ,$$

$q_i$  ... net charge

$b_i$  ... bare charge

$a_i$  ... strength of electron-electron interaction

minimization of energy:

$$a_i(q_i - b_i) + \sum_j M_{ij} q_j = 0 , \quad V_i = \sum_j M_{ij} q_j$$

$$a_i q_i + V_i = a_i b_i = k_i \quad \dots \text{linear q-V relation}$$

parameters  $a_i$ ,  $b_i$  determined from supercell calculations or from local-field CPA

# qV-relation

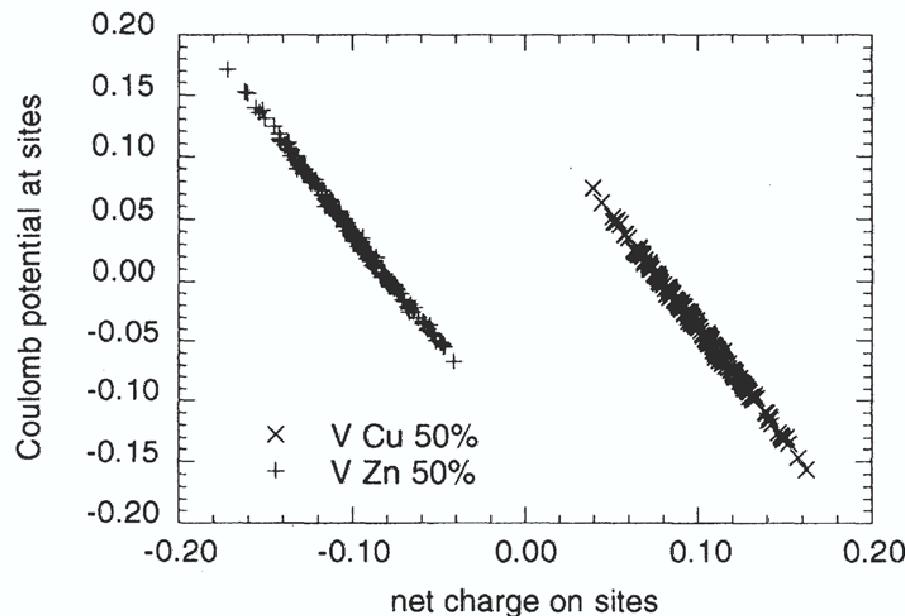


FIG. 1. The Coulomb potential  $V^i$  versus the charges  $q^i$  on the 432 sites in the cell for a 50% Cu-Zn alloy with the bcc crystal structure. The plus signs are the data points for sites that have a Zn atom on them and the crosses are for Cu. The potentials and charges are in dimensionless atomic units as described in the text.

Faulkner, Wang, and Stocks, PRB **52** (1995), 17106.

# BZW model

- screened Madelung interactions  $M_{ij}^{\text{scr}}$
- corrections to effective interatomic interactions

$$V_{ij}^{\text{scr}} = M_{ij}^{\text{scr}} (q_i^A - q_i^B)(q_i^A - q_i^B)$$

- statistics of net charges on atoms in an alloy

V. Drchal, R. Hammerling, and P. Weinberger: Phys. Rev. B **74**, 214204.

# Parameters of BZW model

add external Madelung potential  $V^{\text{ext}}$  on the site at origin

$$aq + V^{\text{tot}} = k = ab, \quad V^{\text{tot}} = V^{\text{ext}} + V^{\text{induced}}$$

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	TB-LMTO-CPA		KKR-CPA	
	Cu	Zn	Cu	Zn
$a$	1.26	1.24	1.84	1.82
$b$	0.0711	-0.0726	0.079	-0.078
$k$	0.0896	-0.0903	0.145	-0.144

- KKR-CPA: J.S. Faulkner, Y. Wang, and G.M. Stocks: Phys. Rev. B **55** (1997), 7492.
- TB-LMTO-CPA: present calculations

# CONCLUSIONS AND OUTLOOK

- flexible, numerically medium-cost method
- electronic structure of clusters (charge distribution, LDOS)
- magnetic structure of clusters (magnetic moments)
- energetics:
  - impurity/complex formation energies
  - effective interactions
  - parameters of BZW model, etc.