



UNESCO/IUPAC Postgraduate Course in Polymer Science

Lecture:

Solution properties of polymers

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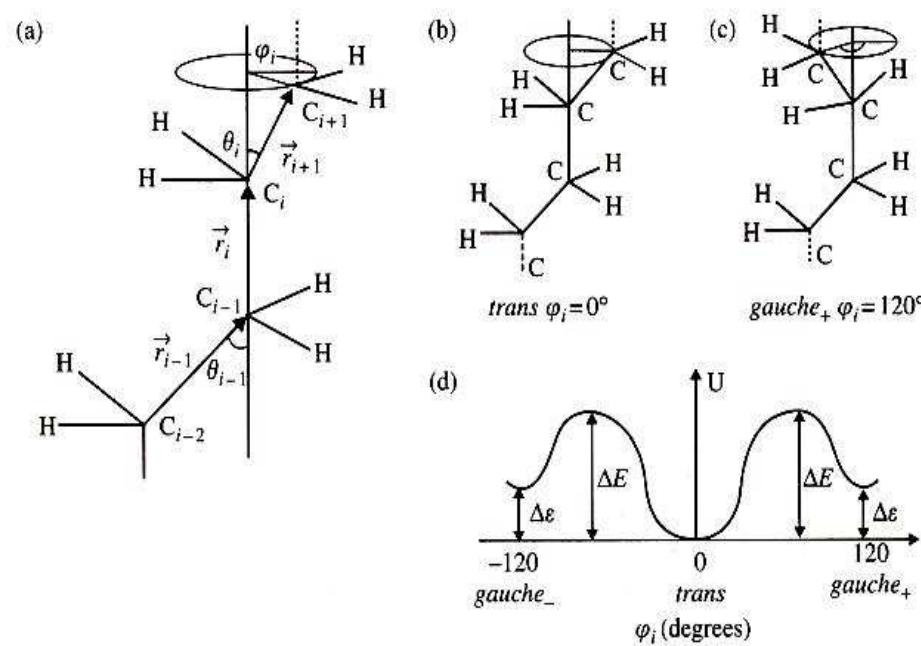
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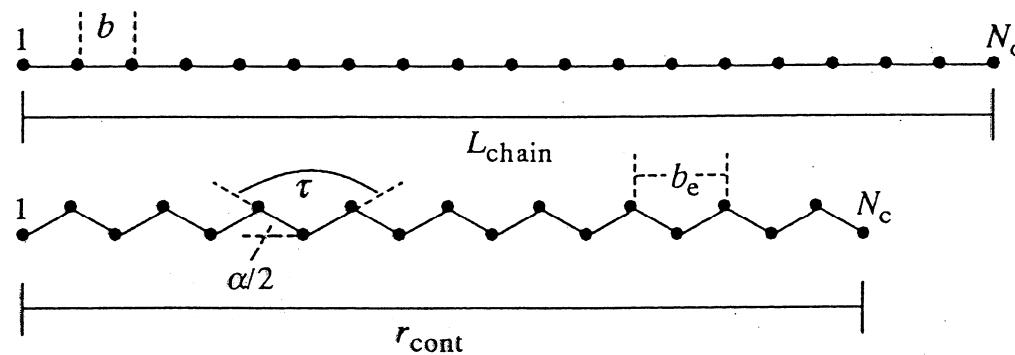
Solution properties of polymers

- Properties of ideal polymer chain
- Thermodynamics of mixing (binary systems containing polymers)
- Osmotic pressure
- Interactions in dilute solutions
- Properties of real polymer chain
- Viscosity
- Light scattering

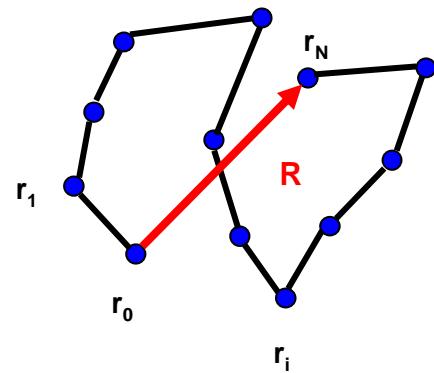
Properties of ideal polymer chain

Thermal motion of molecules: translation
rotation
intramolecular rotation





End-to-end distance



$$\langle \mathbf{R}^2 \rangle \equiv \langle (\mathbf{r}_N - \mathbf{r}_0)^2 \rangle = \left\langle \sum_{i=1}^n \vec{l}_i \sum_{j=1}^n \vec{l}_j \right\rangle = Nl^2 + 2 \sum_{i=1}^n \sum_{j \neq i}^n \langle \vec{l}_i \vec{l}_j \rangle$$

$$\langle \vec{l}_i \vec{l}_j \rangle = l^2 \langle \cos \alpha_{ij} \rangle$$

Freely jointed chain

$$\langle \mathbf{R}^2 \rangle = Nl^2$$

$$\lim_{|i-j| \rightarrow \infty} \langle \cos \alpha_{ij} \rangle = 0$$

$$\langle \mathbf{R}^2 \rangle = l^2 \sum_{i=1}^N \sum_{j=1}^N \langle \cos \alpha_{ij} \rangle = l^2 \sum_{i=1}^N C_i = C_n Nl^2$$

Flory's characteristic ratio C_∞

$$\langle \mathbf{R}^2 \rangle \equiv C_\infty Nl^2$$

Polymer	Structure	C_∞	b (Å)	ρ (g cm ⁻³)	M_0 (g mol ⁻¹)
1,4-Polyisoprene (PI)	-(CH ₂ CH=CHCH(CH ₃))-	4.6	8.2	0.830	113
1,4-Polybutadiene (PB)	-(CH ₂ CH=CHCH ₂)-	5.3	9.6	0.826	105
Polypropylene (PP)	-(CH ₂ CH ₂ (CH ₃))-	5.9	11	0.791	180
Poly(ethylene oxide) (PEO)	-(CH ₂ CH ₂ O)-	6.7	11	1.064	137
Poly(dimethyl siloxane) (PDMS)	-(OSi(CH ₃) ₂)-	6.8	13	0.895	381
Polyethylene (PE)	-(CH ₂ CH ₂)-	7.4	14	0.784	150
Poly(methyl methacrylate) (PMMA)	-(CH ₂ C(CH ₃)(COOCH ₃))-	9.0	17	1.13	655
Atactic polystyrene (PS)	-(CH ₂ CHC ₆ H ₅)-	9.5	18	0.969	720

Equivalent freely jointed chain

Contour lenght

$$nl = R_{\max}$$

Equality of the mean-square end-to-end distance

$$\langle \mathbf{R}^2 \rangle = nb^2 = bR_{\max} = C_\infty Nl^2$$

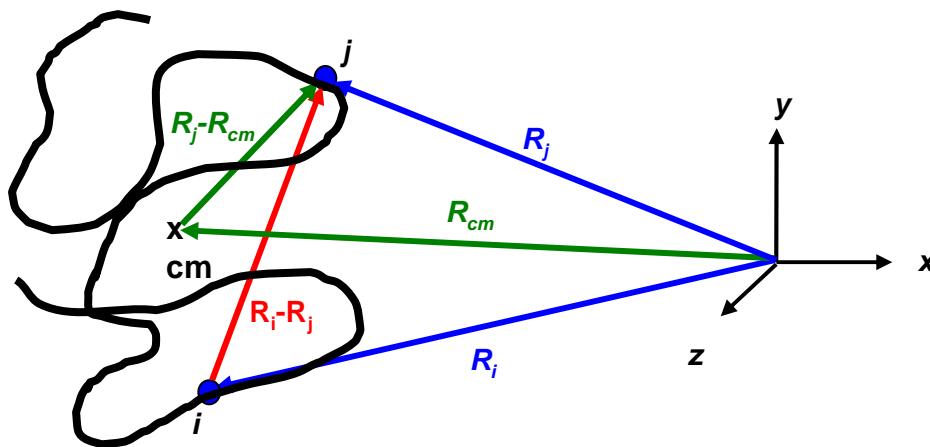
$$n = \frac{R_{\max}^2}{C_\infty Nl^2}$$

$$b = \frac{\langle \mathbf{R}^2 \rangle}{R_{\max}} = \frac{C_\infty Nl^2}{R_{\max}}$$

Freely rotating chain

$$\langle \mathbf{R}^2 \rangle = Nl^2 + 2Nl^2 \frac{\cos\theta}{1-\cos\theta} = Nl^2 \left(\frac{1+\cos\theta}{1-\cos\theta} \right)$$

Radius of gyration



Definition

$$\vec{R}_g^2 \equiv \frac{1}{N} \sum_{i=1}^N (\vec{R}_i - \vec{R}_{cm})^2$$

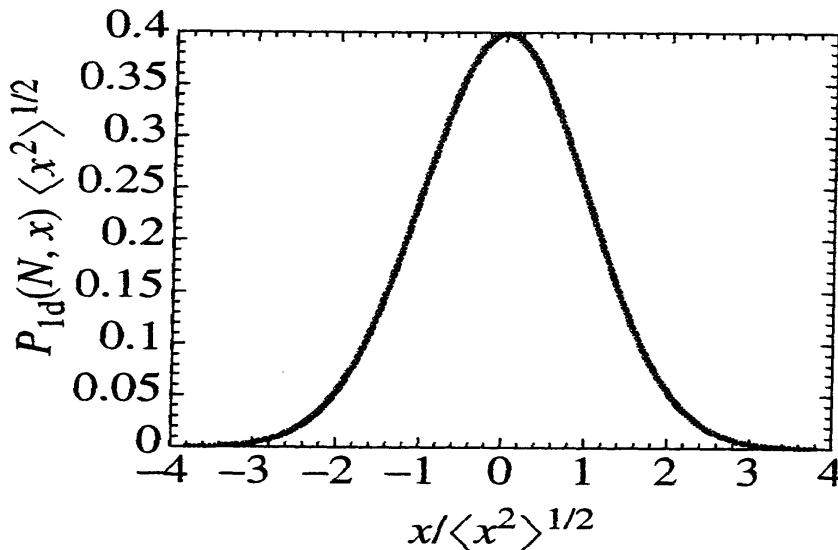
Relation between radius of gyration and end-to-end distance

$$\langle \vec{R}_g^2 \rangle = \frac{b^2 N}{6} = \frac{\langle \vec{R}^2 \rangle}{6}$$

Distribution of end-to-end vectors

Model random walk

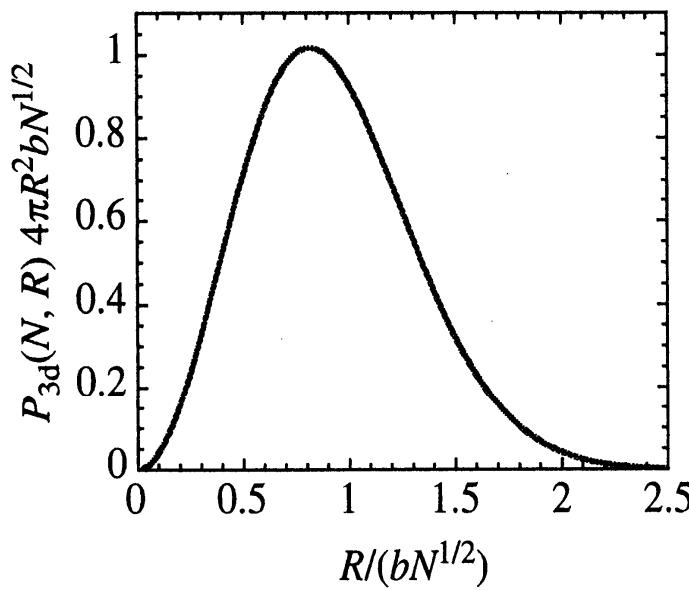
$$P_{1d}(N, x) = \frac{1}{\sqrt{2\pi \langle x^2 \rangle}} \exp\left(-\frac{x^2}{2\langle x^2 \rangle}\right)$$



Normalized one-dimensional Gaussian probability distribution function for occupying position x after random N steps from origin $x=0$

Three-dimensional distribution function

$$P_{3d}(N, R) 4\pi R^2 dR = 4\pi \left(\frac{3}{2\pi N b^2} \right)^{3/2} \exp\left(-\frac{3R^2}{2Nb^2}\right) R^2 dR$$



Normalized distribution function of end-to-end distances for an ideal linear chain

Free energy of an ideal chain

$$S = k \ln \Omega \quad S(N, \vec{R}) = -\frac{3}{2}k \frac{\vec{R}^2}{Nb^2} + S(N, 0)$$

Helmholtz free energy of the stretched chain

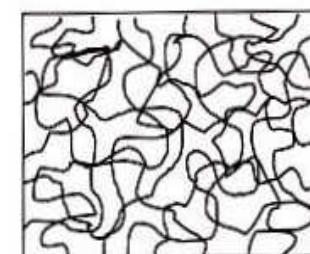
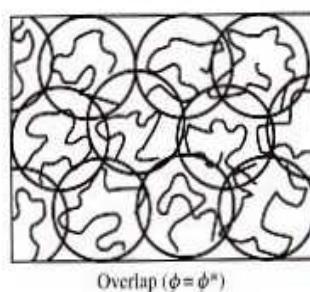
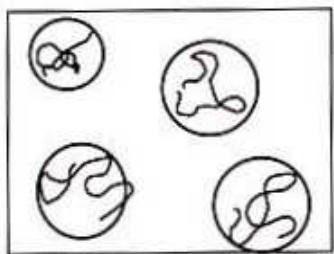
$$F(N, \vec{R}) = \frac{3}{2} kT \frac{\vec{R}^2}{Nb^2} + F(N, 0)$$

Deformation force

$$\vec{f} = \frac{3kT}{Nb^2} \vec{R}$$

Dependence on temperature and number of segments

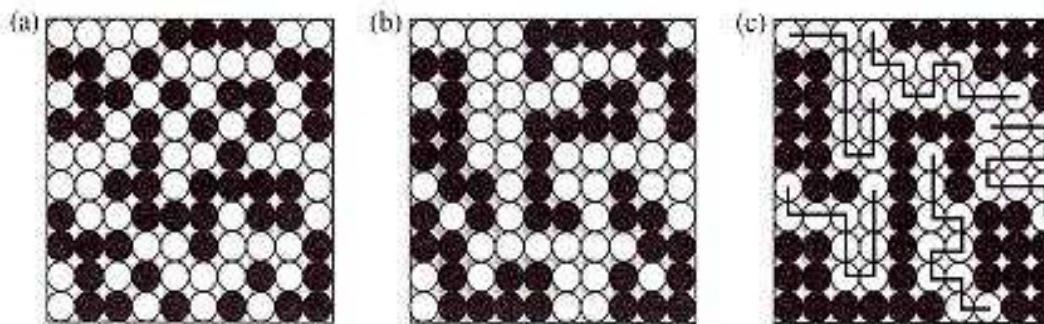
Polymer solutions



$$\varphi^* = \frac{NV_0}{V}$$

Thermodynamics of mixing

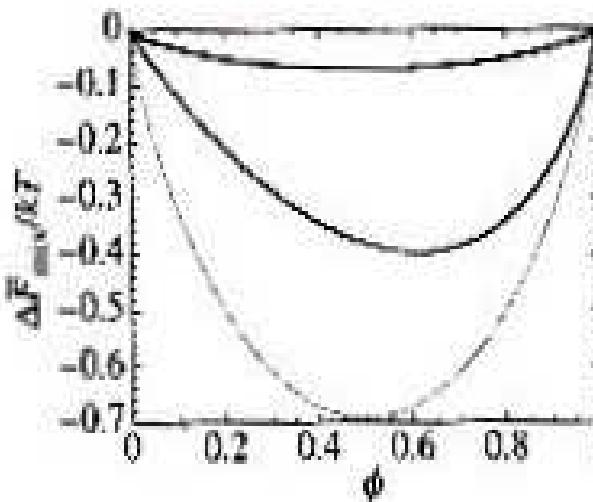
Flory's lattice model



Entropy of mixing

$$\Delta \bar{S}_{mix} = k \left[\frac{\phi_A}{N_A} \ln \phi_A + \frac{\phi_B}{N_B} \ln \phi_B \right]$$

$$\Delta \bar{F}_{mix} = -T \Delta \bar{S}_{mix} = kT \left[\frac{\phi}{N_A} \ln \phi + \frac{1-\phi}{N_B} \ln (1-\phi) \right]$$



The mixing free energy of an ideal mixture is always favourable (negative) and all compositions are stable. The bottom curve is a regular solution with $N_A = N_B = 1$. The middle curve is a polymer solution with $N_A = 10$ and $N_B = 1$. The top curve is a polymer blend with $N_A = N_B = 10$.

Energy of mixing

$$\Delta \bar{U}_{mix} = \frac{z}{2} \phi (1 - \phi) (2u_{AB} - u_{AA} - u_{BB})$$

Flory's interaction parameter χ

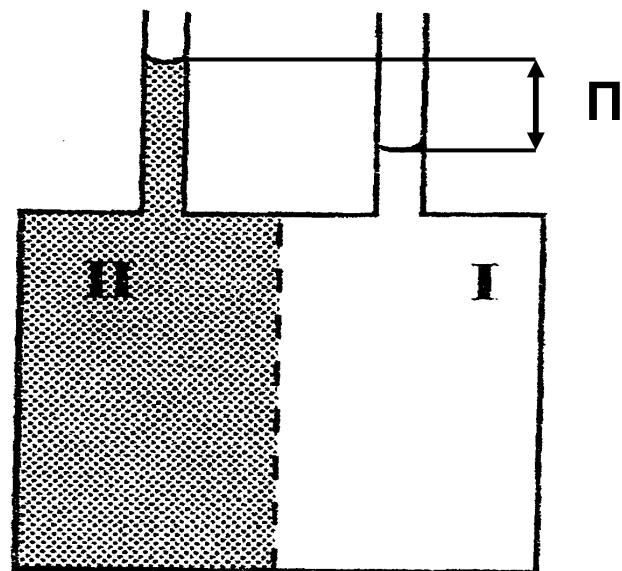
$$\chi \equiv \frac{z}{2} \frac{(2u_{AB} - u_{AA} - u_{BB})}{kT}$$

$$\Delta \bar{F}_{mix} = \Delta \bar{U}_{mix} - T \Delta \bar{S}_{mix} = kT \left[\frac{\phi}{N_A} \ln \phi + \frac{1-\phi}{N_B} \ln (1-\phi) + \chi \phi (1-\phi) \right]$$

Temperature dependence of χ

$$\chi(T) \equiv A + \frac{B}{T}$$

Osmotic pressure



I solvent

II solution

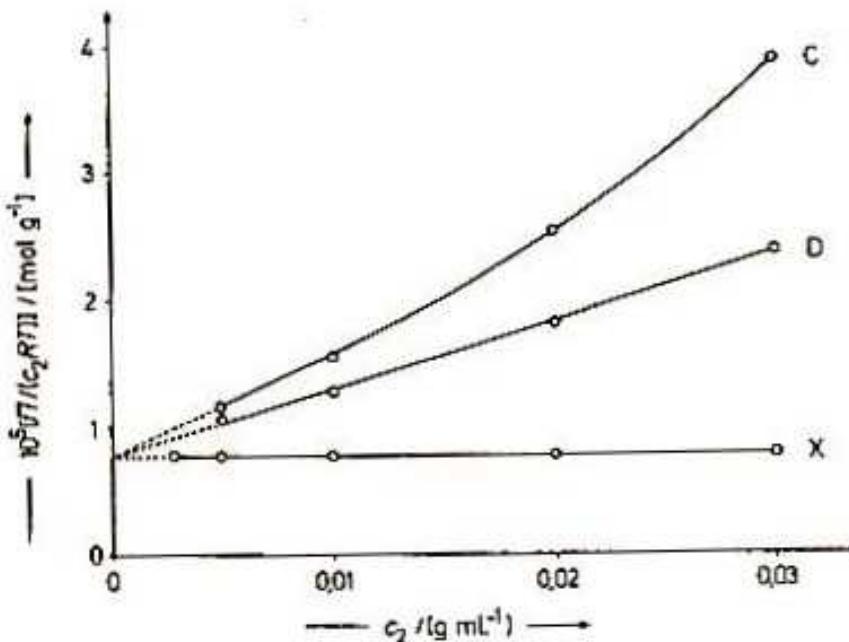
Definition of osmotic pressure

$$\Pi \equiv - \left(\frac{\partial \Delta F_{mix}}{\partial V} \right)_{n_A}$$

$$\Pi = \frac{kT}{b^3} \left[\frac{\phi}{N_A} + (1 - 2\chi) \frac{\phi^2}{2} + \dots \right]$$

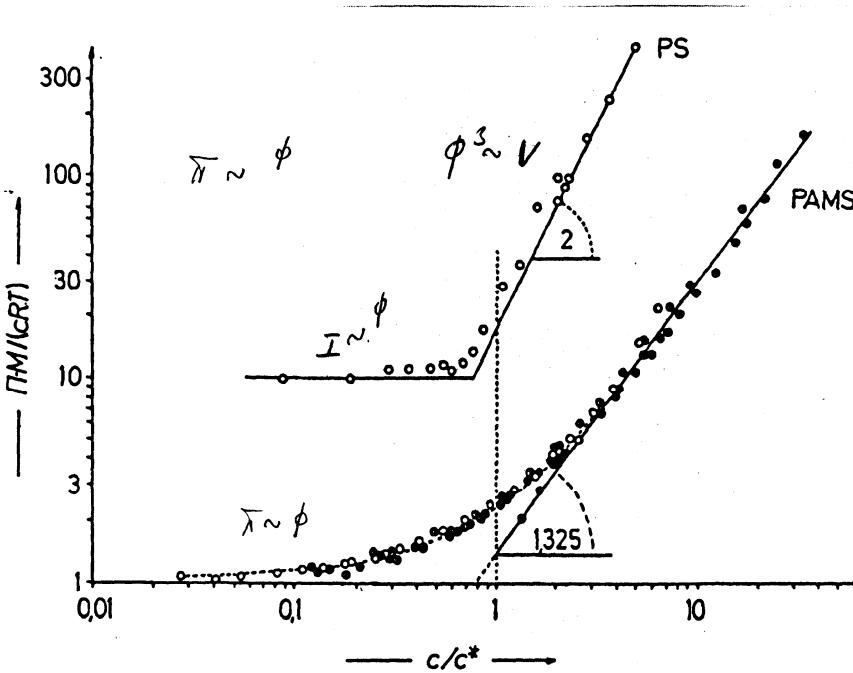
$$\Pi = RT \left(\frac{c}{M} + A_2 c^2 + \dots \right)$$

$$A_2 = \frac{v N_{Avogadro}}{2 M^2} \approx 1 - 2\chi \approx \frac{T - \Theta}{T}$$



Concentration dependence of reduced osmotic pressures of a poly(methyl methacrylate) in various solvents at 20°C (1). X – m-xylene is θ solvent ($A_2 = 0$), 1,4-dioxane D leads to a positive second virial coefficient $A_2 > 0$ (good solvent), and chloroform C furnishes both second and third virial coefficients. The common intercept at $c_2 \rightarrow 0$ equals $1/M_n$

Good solvent $A_2 > 0$; θ solvent $A_2 = 0$; bad solvent $A_2 < 0$



Dependence of reduced osmotic pressures, ($\Pi M/cRT$), on normalized overlap concentration c/c^* (c^* from viscosity) for polystyrenes PS in the Θ -solvent cyclohexane at 34 °C and poly(α -methylstyrene)s PAMS in the good solvent toluene at 25 °C(2). The osmotic overlap concentration is at $c_{\Pi}^* = c_v^*$.

Properties of a real polymer chain in a good solvent

Swelling of polymer coil due to osmotic pressure

$$F_{\text{int}} \approx kT\nu \frac{N^2}{R^3}$$

Total energy

$$F = F_{\text{int}} + F_{\text{ent}} \approx kT \left(\nu \frac{N^2}{R^3} + \frac{R^2}{Nb^2} \right)$$

Equilibrium state

$$\frac{\partial F}{\partial R} = 0 = kT \left(-3 \frac{N^2}{R_F^4} + 2 \frac{R_F}{Nb^2} \right)$$

$$R_F^5 \approx \nu b^2 N^3$$

$$R_F \approx \nu^{\frac{1}{5}} b^{\frac{2}{5}} N^{\frac{3}{5}}$$

Expansion coefficient

$$\alpha = \frac{R_F}{b N^{\frac{1}{2}}} \approx \left(\frac{\nu}{b^3} N^{\frac{1}{2}} \right)^{1/5} \approx z^{\frac{1}{5}}$$

Flory – Fox equation

$$\alpha^2 \equiv \frac{\langle \mathbf{R}^2 \rangle}{\langle \mathbf{R}_0^2 \rangle}$$

$$\alpha^5 - \alpha^3 = 1.276z$$

Viscosity

Increase of solution viscosity caused by polymer dissolved

$$\eta = \eta_0 (1 + 2.5\phi) = \eta_0 \left(1 + 2.5 \frac{V_{pol} c}{M} \right)$$

Definition of the intrinsic viscosity

$$[\eta] = \lim_{c \rightarrow 0} \frac{\eta - \eta_0}{\eta_0} \frac{1}{c} \approx \frac{V_{pol}}{M}$$

Relation between intrinsic viscosity and critical overlap concentrtrion

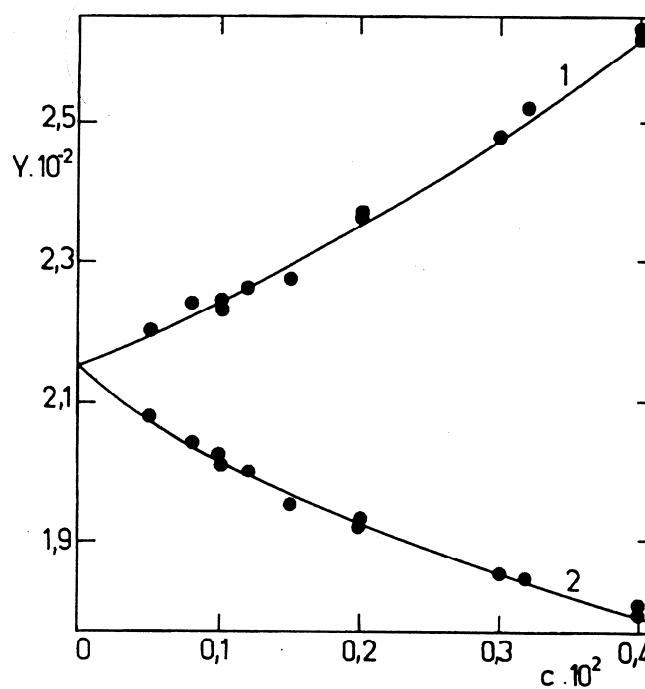
$$[\eta] \approx \frac{V_{pol}}{M} \approx \frac{R^3}{M} \approx \frac{1}{c^*}$$

Huggins equation

$$\frac{\eta - \eta_0}{\eta_0 c} = [\eta] + k_H [\eta]^2 c + \dots$$

Kraemer equation

$$\frac{\ln(\eta/\eta_0)}{c} = [\eta] + \left(k_H - \frac{1}{2} \right) [\eta]^2 c + \dots$$



Determination of the intrinsic viscosity for polymethylmethacrylate in benzene at 25 °C. 1 - Huggins plot 2 - Kraemer plot that extrapolate to the intrinsic viscosity at zero concentration.

Dependence of intrinsic viscosity on molar mass

$$[\eta] = KM^a$$

Light scattering

Intensity of light scattered by one small particle

$$\frac{I}{I_0} = \frac{16\pi^4}{\lambda^4 r^2} \alpha^2$$

Intensity of light scattered by dilute gas

$$\frac{\bar{I}}{I_0} = \frac{4\pi^2 n^2}{\lambda^4 r^2} \left(\frac{dn}{dc} \right)^2 \frac{cM}{N_{Avogadro}}$$

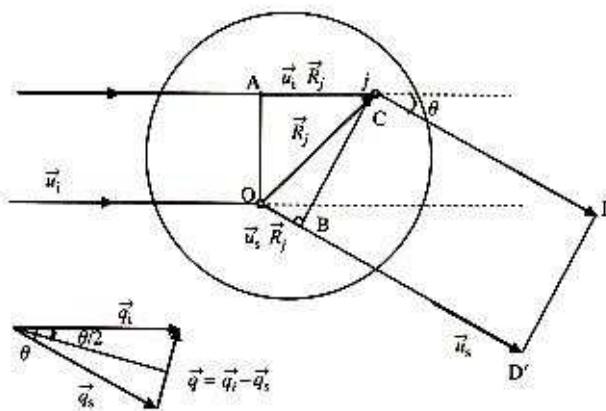
Rayleigh ratio

$$R_\theta = \frac{\bar{I}r^2}{I_0} = \frac{4\pi^2 n^2}{\lambda^4} \left(\frac{dn}{dc} \right)^2 \frac{cM}{N_{Avogadro}} = KcM$$

Role of concentrations fluctuations

$$\frac{Kc}{R_\theta} = \frac{1}{RT} \left(\frac{\partial \Pi}{\partial c} \right)_T = \frac{1}{M} + 2A_2 c + \dots$$

Light scattering by large particles



Phase shift

$$\varphi_j = \frac{2\pi n}{\lambda} (\vec{u}_i - \vec{u}_s) \vec{R}_j = (\vec{q}_i - \vec{q}_s) \vec{R}_j = \vec{q} \vec{R}_j$$

Scattering wave vector

$$q \equiv |\vec{q}| = 2|\vec{q}_i| \sin\left(\frac{\theta}{2}\right) = \frac{4\pi n}{\lambda} \sin\left(\frac{\theta}{2}\right)$$

Form function

$$P(\vec{q}) \equiv \frac{I_s(\vec{q})}{I_s(0)}$$

$$P(q) = \frac{1}{N^2} \sum_{i=1}^N \sum_{j=1}^N \frac{\sin(q R_{ij})}{q R_{ij}}$$

$$P(q) = 1 - \frac{16\pi^2 n^2}{3\lambda^2} \langle R_g^2 \rangle \sin^2\left(\frac{\theta}{2}\right) + \dots$$

$$\left(\frac{Kc}{R_{\theta}}\right)_{c \rightarrow 0} = \frac{1}{M_w}\Bigg[1 + \frac{16\pi^2n^2}{3\lambda}\Big\langle R_g^2\Big\rangle\sin^2\Bigg(\frac{\theta}{2}\Bigg) + ... \Bigg]$$

**The presentation is based on the book *Polymer physics*
by M. Rubinstein and R.H. Colby, Oxford University Press Inc.,
New York, 2003**

- 1) G.V.Schulz, H. Doll, Ber. Deutch. Chem. Ges. 80, (1947), 232**
- 2) P. Štěpánek, K. Perzynski, M. Delsanti, M. Adam,
Macromolecules 17, (1984), 2340**



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