# Particle beams and hydrogen emission in solar flares J. Kašparová<sup>1</sup>, M. Varady<sup>1,2</sup>, P. Heinzel<sup>1</sup>, M. Karlický<sup>1</sup>, Z. Moravec<sup>2</sup>

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R(M)HD seminar, Astronomický ústav, Ondřejov, 22 Oct 2009

## FLARE MODEL



- standard two-ribbon flare
- energy release during mag. reconnection in corona
- plasma heating, particle acceleration
- beams: form of flare energy transport
  - energy loss via Coulomb collision and return current ⇒ heating, excitation, ionisation
  - X-ray, radio emission
- increased emission in EUV, UV, optical, IR bands

- compute time evolution of hydrogen continuum and line profiles
- study influence of particle beams on hydrogen emission
- hydrodynamic and radiative response of solar atmosphere to heating by particle beams: radiative hydro code
- response of hydrostatic VAL C atmosphere to beam heating: HD code
- propagation and energy losses of the beam: test particle code
- ionisation and hydrogen emission: NLTE radiative transfer code





- standard set of 1D HD equations in one fluid approximation describes the state and evolution of plasma along magnetic field lines
- included processes
  - thermal conduction (Spitzer's classical approx.)
  - optically thin (Rosner et al., 1978) and thick (Peres et al., 1982) radiative losses (approx. expressions)
  - heating given by beam energy deposit into the atmosphere and return current calculated by a particle code
  - ionisation calculated by a NLTE radiative transfer code
- numerical methods
  - LCPFCT algorithm (for generalised continuity equations)
  - timestep splitting method
  - Crank-Nicholson algorithm for conduction

### PARTICLE BEAM HEATING (Varady, Karlický, Moravec)

- beam energy deposit is calculated by a test particle code for the instant properties of the atmosphere
- the code includes
  - Coulomb collisions with neutrals and electrons (Emslie, 1978)
  - electron scattering (Bai, 1982)
  - optionally return current (runaway approx., Varady et al., 2005)
- power-law particle beams
  - power-law index  $\delta = 3 7$
  - low-energy cutoff  $\approx 10~{\rm keV}~{\rm (MeV)}$  and high-energy cutoff  $\approx 100~{\rm keV}~{\rm (MeV)}$
  - time modulation of energy flux F(t)
  - properties could be obtained from X-rays



- NLTE radiative transfer for hydrogen is calculated in the lower part of the loop using the instant values of T,  $n_{\rm H}$  and the energy deposit to hydrogen  $E_{\rm H}$
- 5-level + continuum model of hydrogen
- time dependent ESE

$$\frac{\partial n_i}{\partial t} = \sum_{j \neq i} n_j P_{ji} - n_i \sum_{j \neq i} P_{ij}$$

• excitation and ionisation of hydrogen by the particle beam is taken into account by nonthermal collisional rates  $C_{1i}^{\text{nt}}$  (Fang et al., 1993)

$$C_{1j}^{\rm nt} = k_{1j} \frac{E_{\rm H}}{n_1}$$

•  $C_{1j}^{\text{nt}}$  included into transition rates  $P_{ij}$ 

$$P_{1j} = R_{1j} + C_{1j} + C_{1j}^{\text{nt}}$$

• numerics: MALI, linearisation of ESE, Newton-Raphson, Crank-Nicholson schemes

# Effects of the $C^{nt}$ - ionisation

- NLTE ionisation lags behind the time evolution of T due to time evolution of the ratio of the number of recombinations to photoionisations
- $C^{\rm nt}$  increase ionisation  $\approx 1000$  km where temperature increase itself does not completely ionise the plasma
  - stronger effect for high beam energy flux F and low index  $\delta$



## Effects of the $C^{nt}$ - Balmer lines

- $C^{\rm nt}$  influence is strongly linked to  $E_{\rm H}$  as a function of height
  - line intensities are affected according to their formation heights



• new Hlpha wing formation region at  $E_{
m H}$  maximum, stronger for lower  $\delta$ 



Effects of the return current (RC) -  $H\alpha$  line

• return current in runaway approximation:

 $j_{\rm b} = j_{\rm RC}$   $\alpha = n_{\rm RC}/n_{\rm e} = 0.1$ 

- increase in the H  $\alpha$  line centre intensity at  $\sim$  0.4 s
  - result of the higher total energy deposit and subsequent temperature increase at  $\sim 2000~\rm km$
  - $C^{\rm nt}$  again create new formation region at  $E_{\rm H}$  peak location





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$$\delta = 3, F = 10^{12} \text{ erg cm}^{-2} \text{ s}^{-1}$$

- $E_1 = 5, 20 \text{ MeV}$ 
  - same hard X-rays as electron beams (deka-MeV protons)
  - higher energy deposit, temperature at  $z \ge 500$  km for  $E_1 = 5$  MeV





## Electron beams - ${\rm Ly}\alpha$

• during heating Ly $\alpha$  is formed in lower atmospheric layers



• abrupt drop at  $\lambda = 0.45$  Å is due to narrowing of formation region;  $\eta, \tau$  decrease due to decrease in  $n_1, n_2$  caused by  $C^{\text{nt}}$  and heating

# Electron beams - far ir and MM continua $(35\mu$ M - 1cm)

- $\bullet\,$  assumed to be of thermal origin and due to H I and H^- free-free processes
  - Planck source function (LTE), opacity  $\kappa_{\nu}$  calculated using non-LTE populations



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•  $C^{\rm nt}$  affect  $\lambda < 0.2$  mm since larger  $\lambda$  originates at layers above significant  $E_{\rm H}$ 

### DIAGNOSTIC TOOLS - TIME CORRELATION

- can we clearly recognise  $C^{nt}$  effects in hydrogen emission?
  - fast hydrogen lines/continua variations exhibit a good correlation with beam flux variations
  - Ly $\alpha$  wings show anti-correlation





#### DIAGNOSTIC TOOLS - BEAM PARAMETERS

- despite clear influence of  $C^{nt}$  on hydrogen emission, unambiguous diagnostic tool has not been found
  - line/intensity ratios R or wavelength-integrated intensity I<sub>tot</sub> do not exhibit unique and systematic behaviour with beam parameters





- electron/proton beam heating significantly affects hydrogen emission on time scale of the heating
- correlation of lines/continua variations with hard X-rays presents only an indirect indication of pulse beam heating

 $\downarrow$  comparison of simulations with observations is needed

- test flare, electron beam parameters from Yohkoh, however no hydrogen data
- more RHESSI flares, Hα data available (Wrocław, Ondřejov)



# The End

