

Prague, September 6 2012

Neutrinos from the Beyond-Standard-Model perspective

Michal Malinský

AHEP group of IFIC, CSIC/University of Valencia

Standard model matter fields

$$\begin{pmatrix} u \\ d \end{pmatrix}_L$$

$$u_R \\ d_R$$

$$\begin{pmatrix} c \\ s \end{pmatrix}_L$$

$$c_R \\ s_R$$

$$\begin{pmatrix} t \\ b \end{pmatrix}_L$$

$$t_R \\ b_R$$

$$\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$$

$$e_R$$

$$\begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}_L$$

$$\mu_R$$

$$\begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}_L$$

$$\tau_R$$

Neutrino mass generation in the SM?

Weinberg's d=5 operator

S. Weinberg, Phys. Rev. Lett. 43, 1566 (1979)

$$\mathcal{L} \ni \frac{LLHH}{\Lambda}$$

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- violates the lepton number global symmetry of the SM!
- good to have the “complete Higgs doublet” :-)

Baryon and lepton number violation in the SM

Anomalies:

$$\mathcal{A} \propto \frac{1}{32\pi^2} \text{Tr} (\{T_a, T_b\} T) \tilde{F}_{\mu\nu}^a F^{b\mu\nu}$$

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Standard model matter fields + RH neutrinos

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Standard model matter fields + RH neutrinos

$$Y_{Dij} \overline{Q_{Li}} \langle H \rangle D_{Rj} + Y_{Uij} \overline{Q_{Li}} \langle \tilde{H} \rangle U_{Rj} + Y_{Eij} \overline{L_{Li}} \langle H \rangle E_{Rj} + Y_{Nij} \overline{L_{Li}} \langle \tilde{H} \rangle N_{Rj} + h.c.$$
$$+ \frac{1}{2} \textcolor{red}{M_{Rij}} \overline{N_{Ri}^c} N_{Rj} + h.c.$$

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$$M_\nu = \begin{pmatrix} 0 & Y_N v \\ Y_N^T v & \textcolor{red}{M_R} \end{pmatrix}$$

$$m_\nu = Y_N \textcolor{red}{M_R}^{-1} v^2 Y_N^T$$

Seesaw systematics

Renormalizable “openings” of the Weinberg operator

$$\mathcal{L} \ni \frac{LLHH}{\Lambda}$$

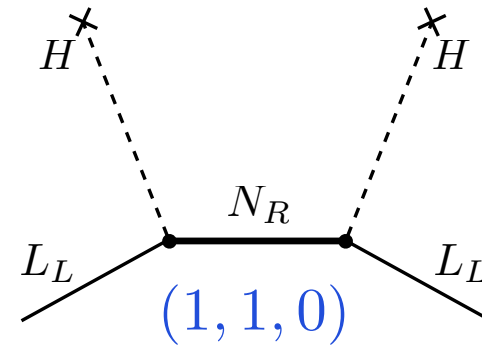


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type-I seesaw

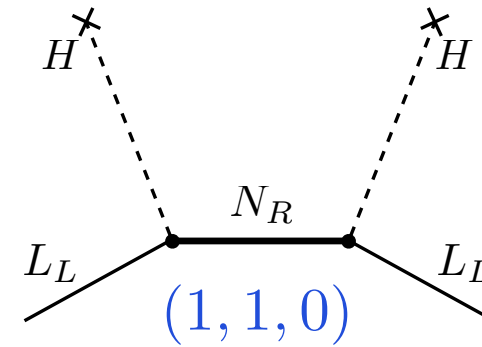


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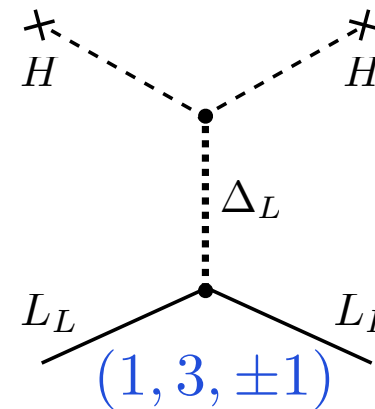
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type-II seesaw



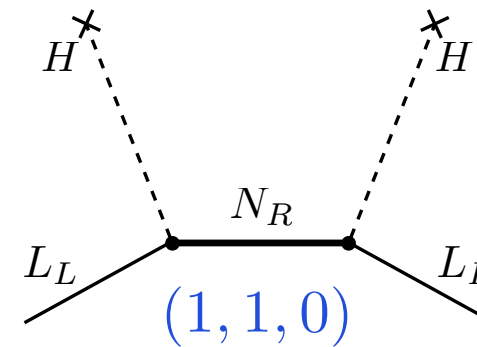
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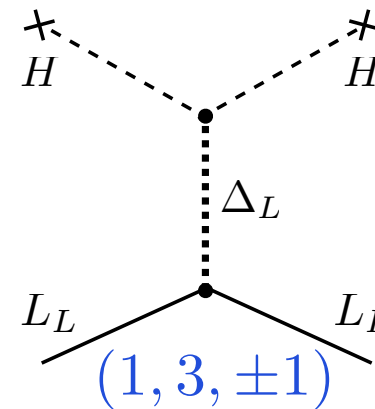
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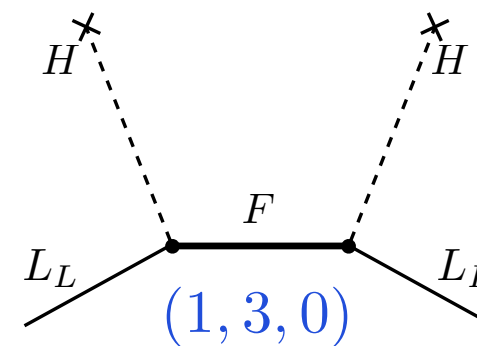
type-I seesaw



type-II seesaw



type-III seesaw



Seesaw scale?

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Lower bounds:

- Neutrino oscillations: $\Delta m_{\odot}^2 = (8.0 \pm 0.3) \times 10^{-5} \text{eV}^2$
 $|\Delta m_A^2| = (2.5 \pm 0.3) \times 10^{-3} \text{eV}^2$

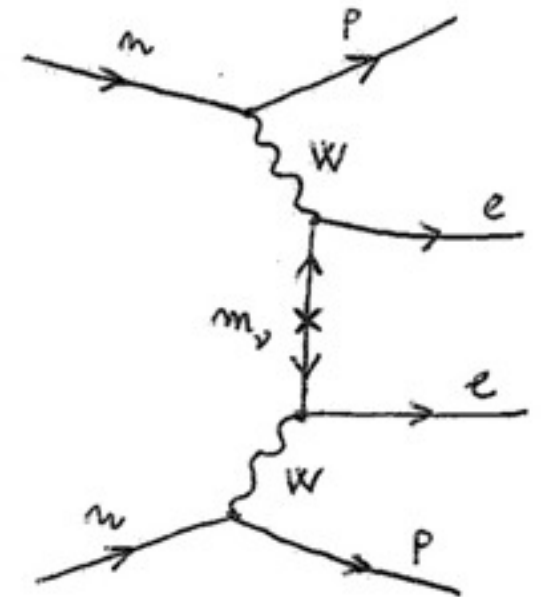
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Upper bounds:

- Cosmology (structure): $\sum_i m_i \lesssim 1 \text{eV}$
- $0\nu 2\beta$: $\langle m^{ee} \rangle \lesssim 1 \text{eV}$



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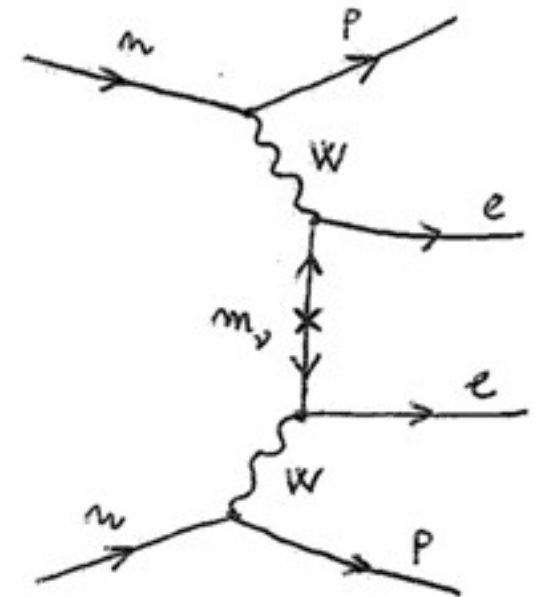
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$$\Lambda \sim (10^{12} - 10^{14}) \text{ GeV}$$



Leptogenesis?

$$\frac{n_B}{n_\gamma} \equiv \eta_B = (6.1 \pm 0.3) \times 10^{-10}$$

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$$\partial^\mu J_\mu^{B+L} \neq 0$$

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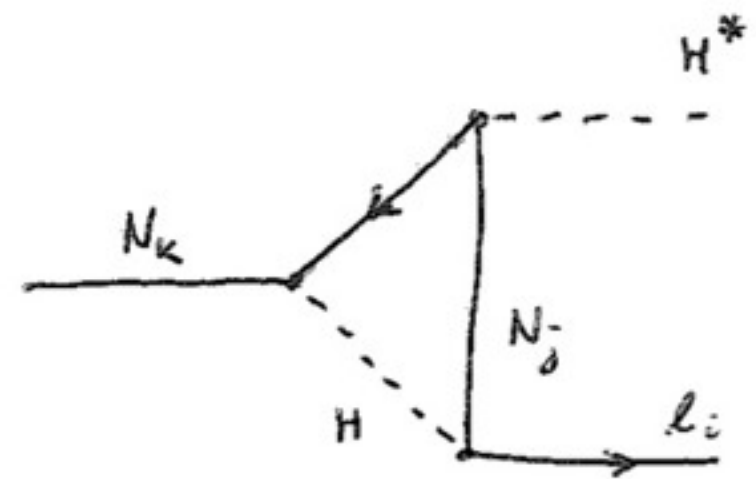
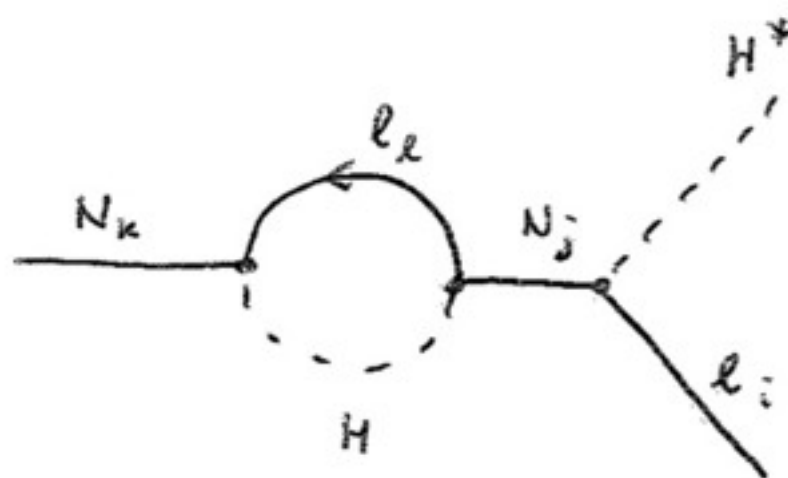
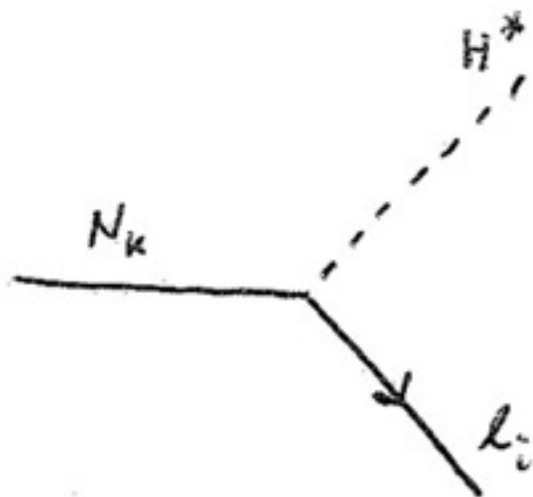
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In the hot early Universe one can transfer L into B

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Generating net L: $\epsilon_1 = \frac{\sum_\alpha [\Gamma(N_1 \rightarrow \ell_\alpha H) - \Gamma(N_1 \rightarrow \bar{\ell}_\alpha \bar{H})]}{\sum_\alpha [\Gamma(N_1 \rightarrow \ell_\alpha H) + \Gamma(N_1 \rightarrow \bar{\ell}_\alpha \bar{H})]}$



Leptogenesis?

CP asymmetry:

$$\epsilon_1 \approx -\frac{3}{8\pi} \frac{1}{(Y_N Y_N^\dagger)_{11}} \sum_{i=2,3} \text{Im} \left[(Y_N Y_N^\dagger)_{1i}^2 \right] \frac{M_1}{M_i}$$

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Davidson-Ibarra bound:

S. Davidson and A. Ibarra, Phys. Lett. B535, 25 (2002)

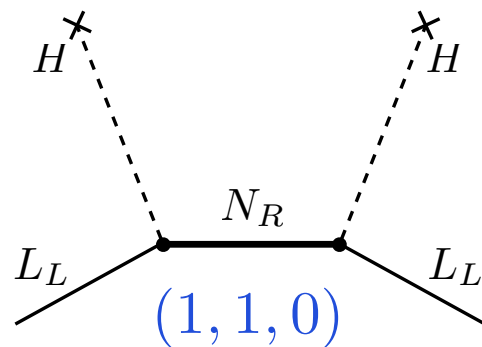
$$|\epsilon_1| \leq \frac{3}{16\pi} \frac{M_1(m_3 - m_2)}{v^2}$$

$$M_1 \gtrsim 10^9 \text{ GeV}$$

“Light” seesaw?

LHC phenomenology(?) - mediators review: arXiv:1001.2693 [hep-ph]

Type-I seesaw: - generally problematic, Yukawa's are too small!

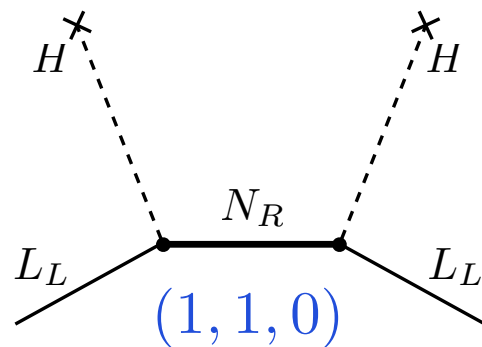


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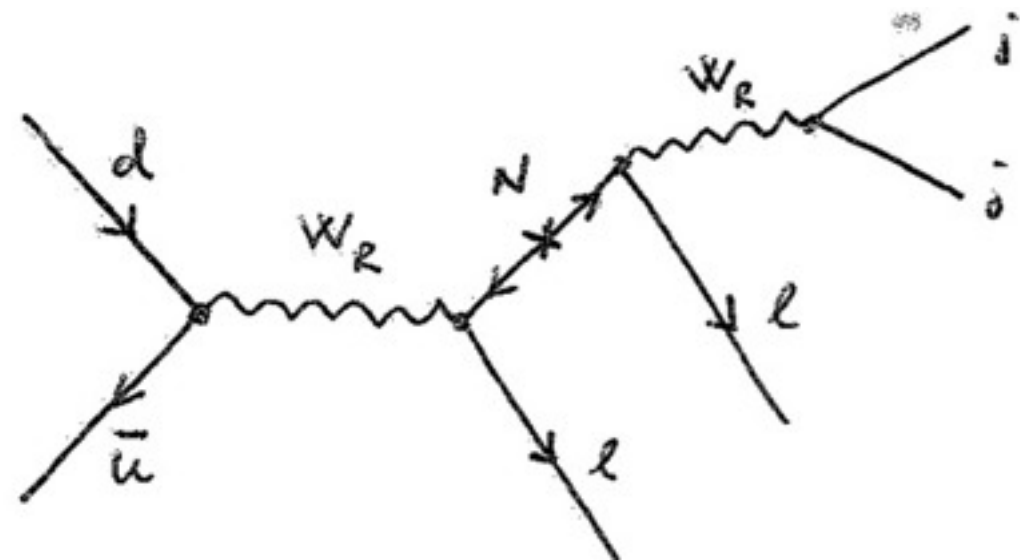
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Type-I seesaw:

- generally problematic, Yukawa's are too small!
- Much better with **extended gauge symmetries**



W.Y. Keung and G. Senjanovic, Phys. Rev. Lett. 50, 1427 (1983)

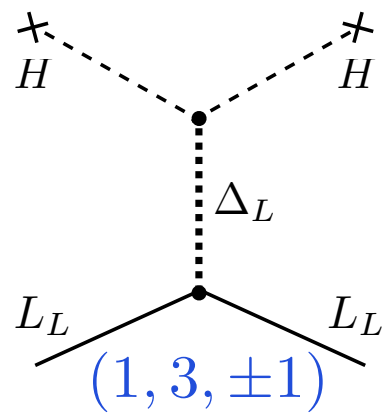


- “essentially background-free” same-sign dilepton signal

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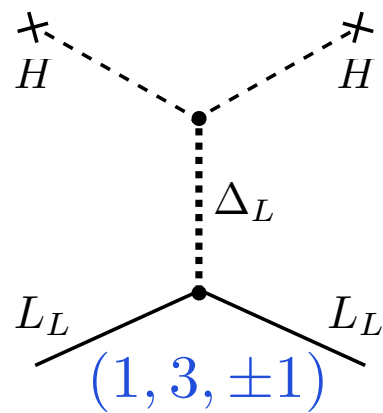
Type-II seesaw: - doubly-charged scalar in the spectrum!



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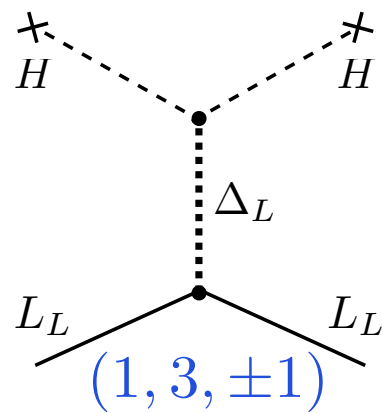
- same sign dilepton pairs (boosted)

$$Z^* \rightarrow \Delta^{++} \Delta^{--} \rightarrow (l^+ l^+) (l^- l^-)$$

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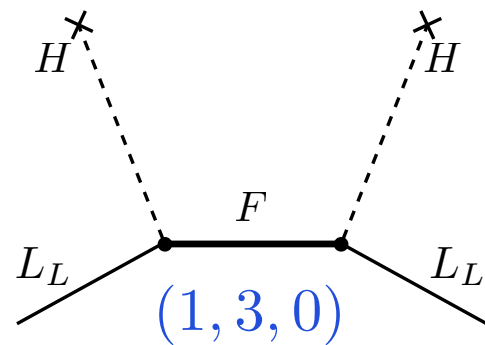
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- decays rely on the Yukawa couplings

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LHC phenomenology(?) - mediators review: [arXiv:1001.2693 \[hep-ph\]](#)

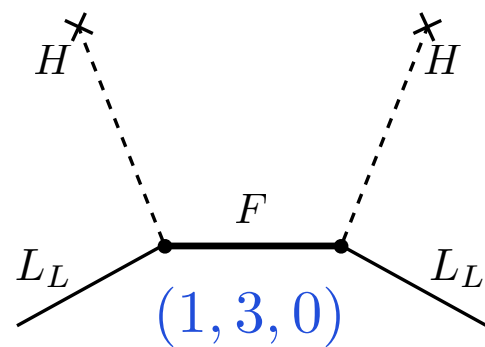
- Type-III seesaw:**
- neutral and charged fermions
 - triplet feels the SM gauge interactions - better!



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Type-III seesaw:



- neutral and charged fermions
- triplet feels the SM gauge interactions - better!
- multi-lepton channels as in type-II, different kinematics!

$$E^+ \rightarrow Z^* l^+ \rightarrow (l^+ l^-) l^+$$

$$E^- \rightarrow Z^* l^- \rightarrow (\nu \bar{\nu}) l^-$$

The fate of B-L (global/gauged?)

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spontaneous breaking

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$U(1)$

$$J_\mu = \Phi^* \partial_\mu \Phi - \Phi \partial_\mu \Phi^* + \bar{\psi} \gamma_\mu \psi$$

Gelmini, Roncandelli 1980

$$\partial^\mu J_\mu = 0$$

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$$\partial^\mu J_\mu = 0$$

$$U(1) \rightarrow \emptyset \quad \langle \Phi \rangle = \Lambda \quad G = \text{Im} \Phi$$

Goldstone boson!

$$\square G + \frac{1}{\Lambda} \partial^\mu (\bar{\psi} \gamma_\mu \psi) = 0$$

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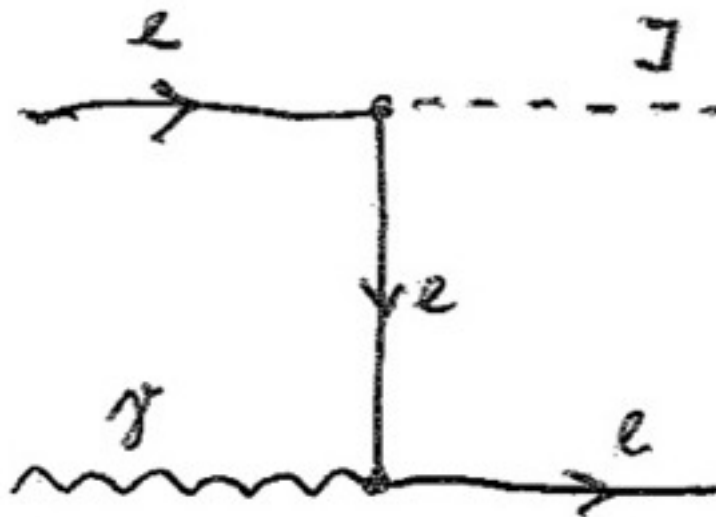
$$\square G + \frac{1}{\Lambda} \partial^\mu (\bar{\psi} \gamma_\mu \psi) = 0$$

$$\mathcal{L} \ni \partial^\mu G \partial_\mu G - \frac{G}{\Lambda} \partial^\mu (\bar{\psi} \gamma_\mu \psi)$$

Majoron

Stellar photoproduction of Majorons (J)

$$\gamma + e \rightarrow e + J$$

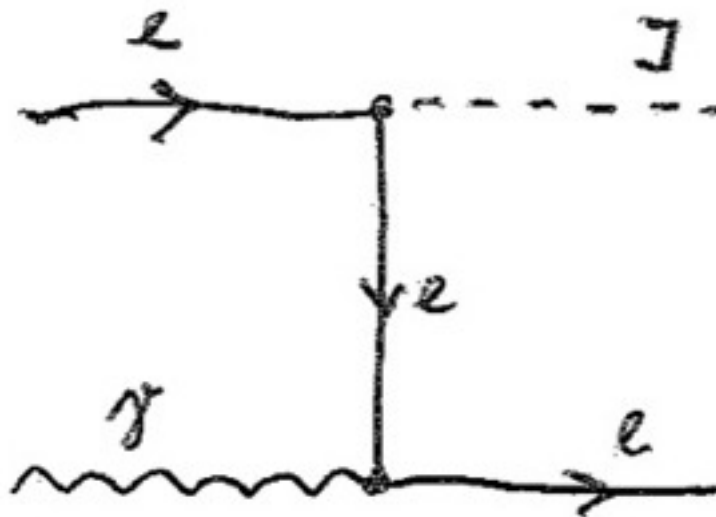


$$L_J \propto \frac{\alpha g^2 T^6}{m_e^4 m_p}$$

Majoron

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$$L_J \propto \frac{\alpha g^2 T^6}{m_e^4 m_p}$$

Red giants: $g < 10^{-12}$

Chanda, Nieves, Pal, PRD37, 1988

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In order to ever gauge B-L three RH neutrinos are needed!

Standard model matter fields + 3 RH neutrinos

	T_L^3	Y	Q	$(B - L)/2$
$\begin{pmatrix} u \\ d \end{pmatrix}_L$	$+\frac{1}{2}$ $-\frac{1}{2}$	$+\frac{1}{6}$	$+\frac{2}{3}$ $-\frac{1}{3}$	$+\frac{1}{6}$
u_R	0	$+\frac{2}{3}$	$+\frac{2}{3}$	$+\frac{1}{6}$
d_R	0	$-\frac{1}{3}$	$-\frac{1}{3}$	
$\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$	$+\frac{1}{2}$ $-\frac{1}{2}$	$-\frac{1}{2}$	0 -1	$-\frac{1}{2}$
ν_R	0	0	0	$-\frac{1}{2}$
e_R	0	-1	-1	

Standard model matter fields + 3 RH neutrinos

	T_L^3	Y	Q	$(B - L)/2$	Another T ³ -like generator?
$\begin{pmatrix} u \\ d \end{pmatrix}_L$	$+\frac{1}{2}$ $-\frac{1}{2}$	$+\frac{1}{6}$	$+\frac{2}{3}$ $-\frac{1}{3}$	$+\frac{1}{6}$	0
u_R	0	$+\frac{2}{3}$	$+\frac{2}{3}$	$+\frac{1}{6}$	$+\frac{1}{2}$
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$\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$	$+\frac{1}{2}$ $-\frac{1}{2}$	$-\frac{1}{2}$	0 -1	$-\frac{1}{2}$	0
ν_R	0	0	0	$-\frac{1}{2}$	$+\frac{1}{2}$
e_R	0	-1	-1	$-\frac{1}{2}$	$-\frac{1}{2}$

Left-right models

$$SU(3)_c \otimes SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$$

recent review: G. Senjanovic, Riv. Nuovo Cim. 034, 2011

$$\begin{pmatrix} u \\ d \end{pmatrix}_L$$

$$u_R$$

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$$\begin{pmatrix} u \\ d \end{pmatrix}_L$$

$$\begin{matrix} u_R \\ d_R \end{matrix} \longrightarrow \begin{pmatrix} u \\ d \end{pmatrix}_R$$

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$$\begin{matrix} \nu_R \\ e_R \end{matrix} \longrightarrow \begin{pmatrix} \nu \\ e \end{pmatrix}_R$$

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- high-scale parity restoration

- Z', W' gauge bosons

- Yukawa “unification”

$$V_{CKM} \approx 1$$

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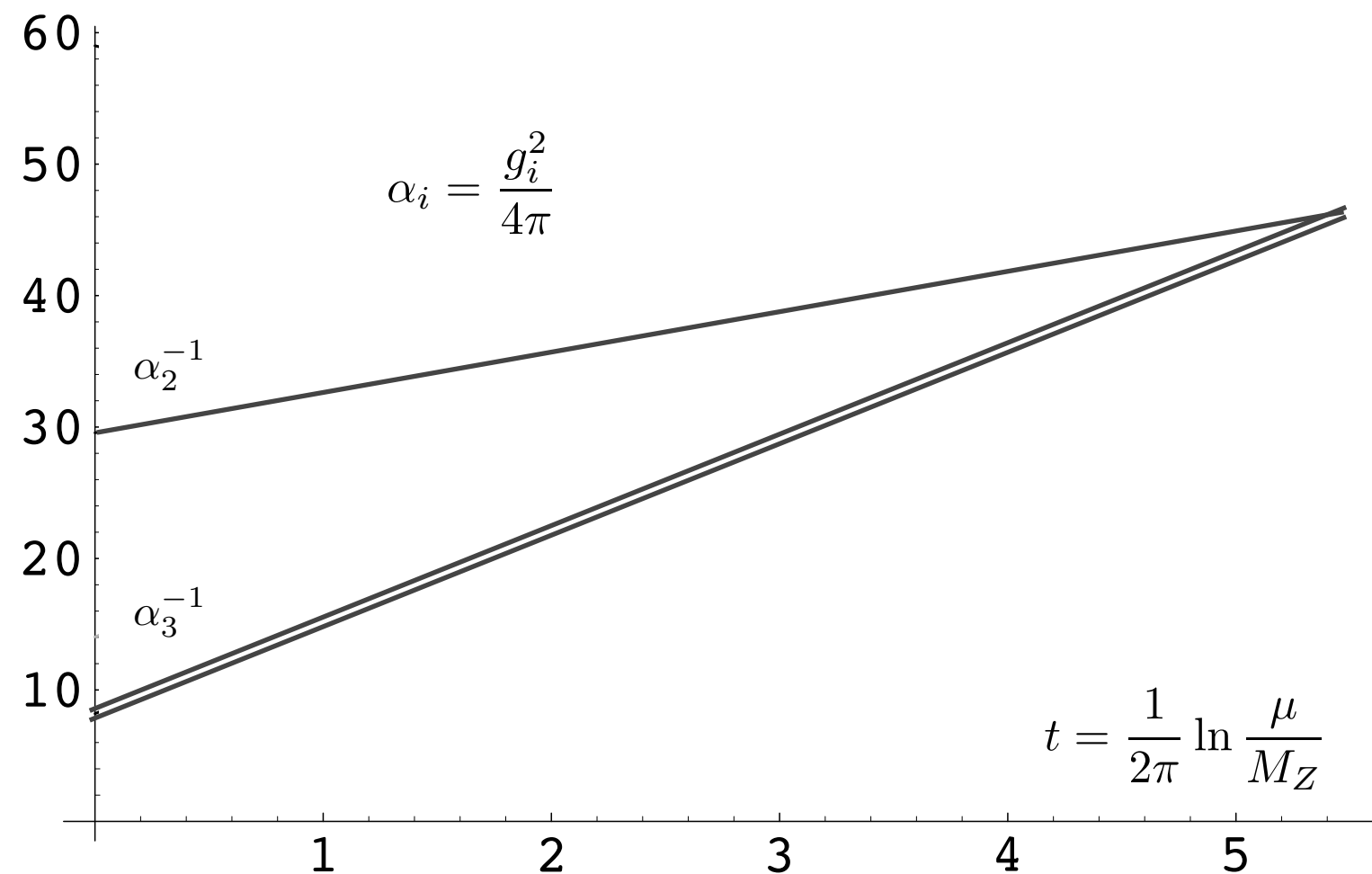
- Yukawa “unification”

$$V_{CKM} \approx 1$$

Unfortunately, not strong enough to tell us much about neutrinos

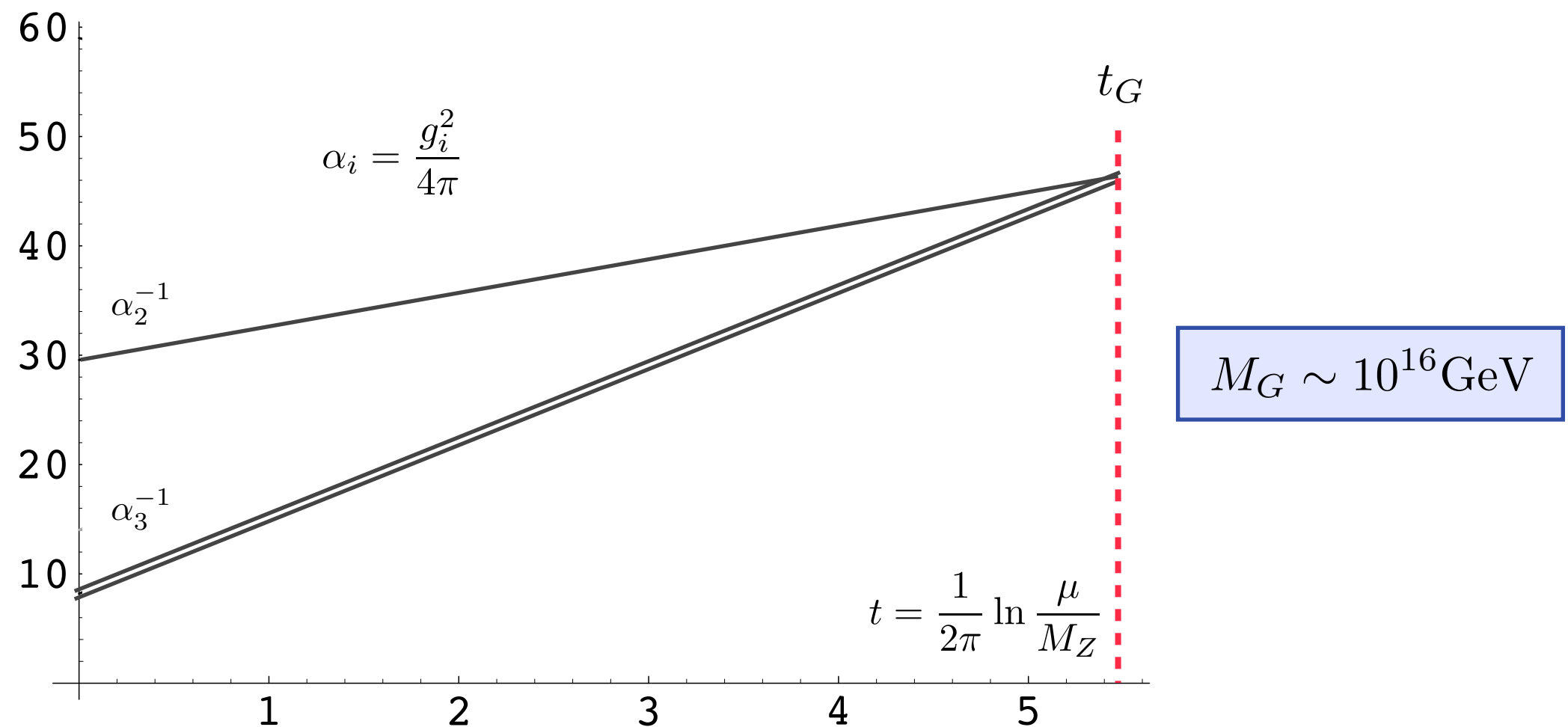
Even stronger gauge symmetries?

the SM gauge couplings seem to converge at high energies



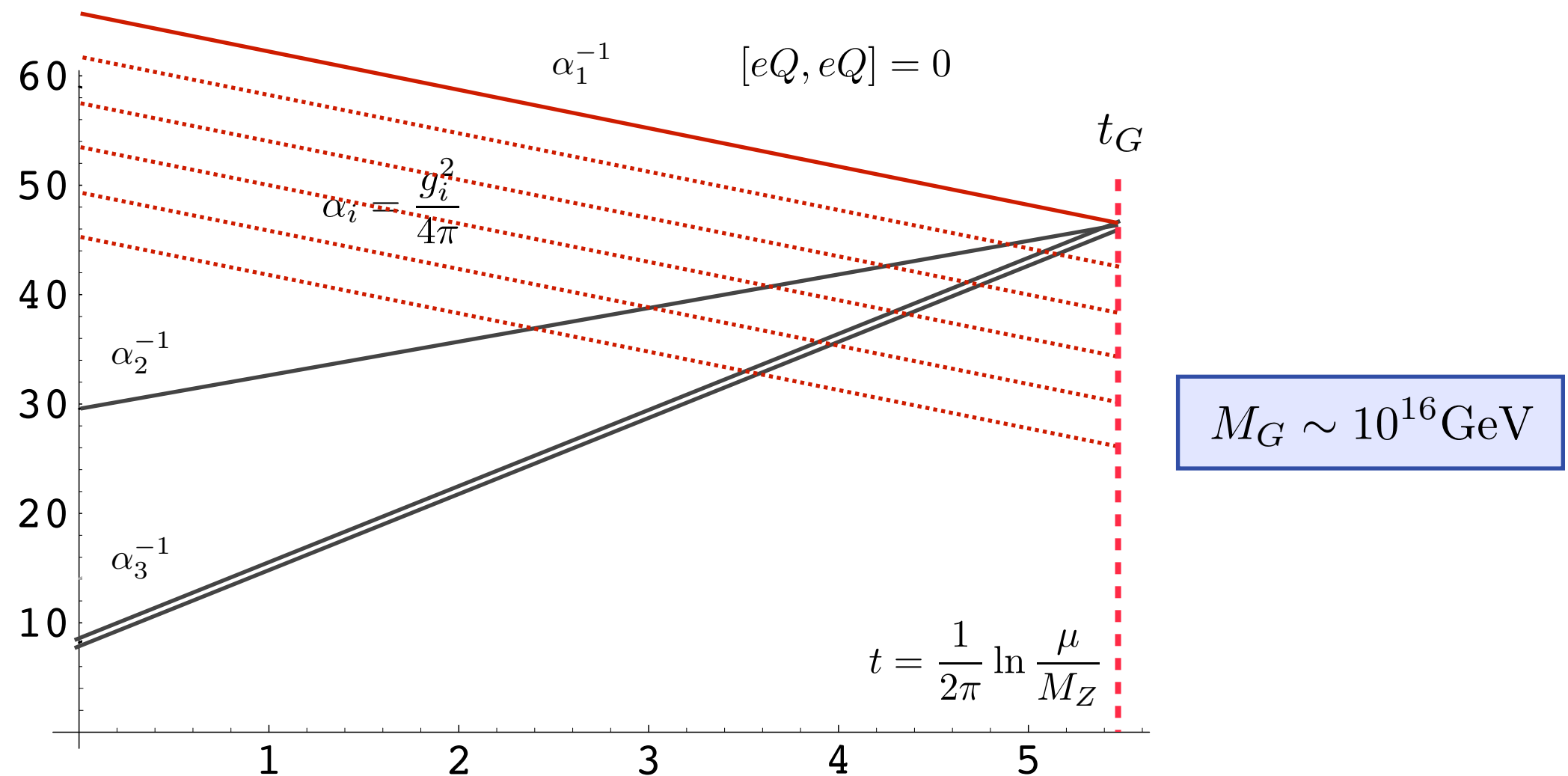
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GUT basics

H.Georgi, S.Glashow, Phys.Rev.Lett. 30 (1974)

GUTs are spontaneously broken BSM gauge theories based on simple compact gauge groups

The physics case

- charge quantization
- monopoles
- baryon and lepton number violation
- partly flavour

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Georgi-Glashow model - a prototype GUT

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$$(1, 2, -\frac{1}{2}) \quad \begin{pmatrix} \nu_e \\ e \end{pmatrix} \quad \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}$$

$$(1, 1, +1) \quad e^c \quad \mu^c$$

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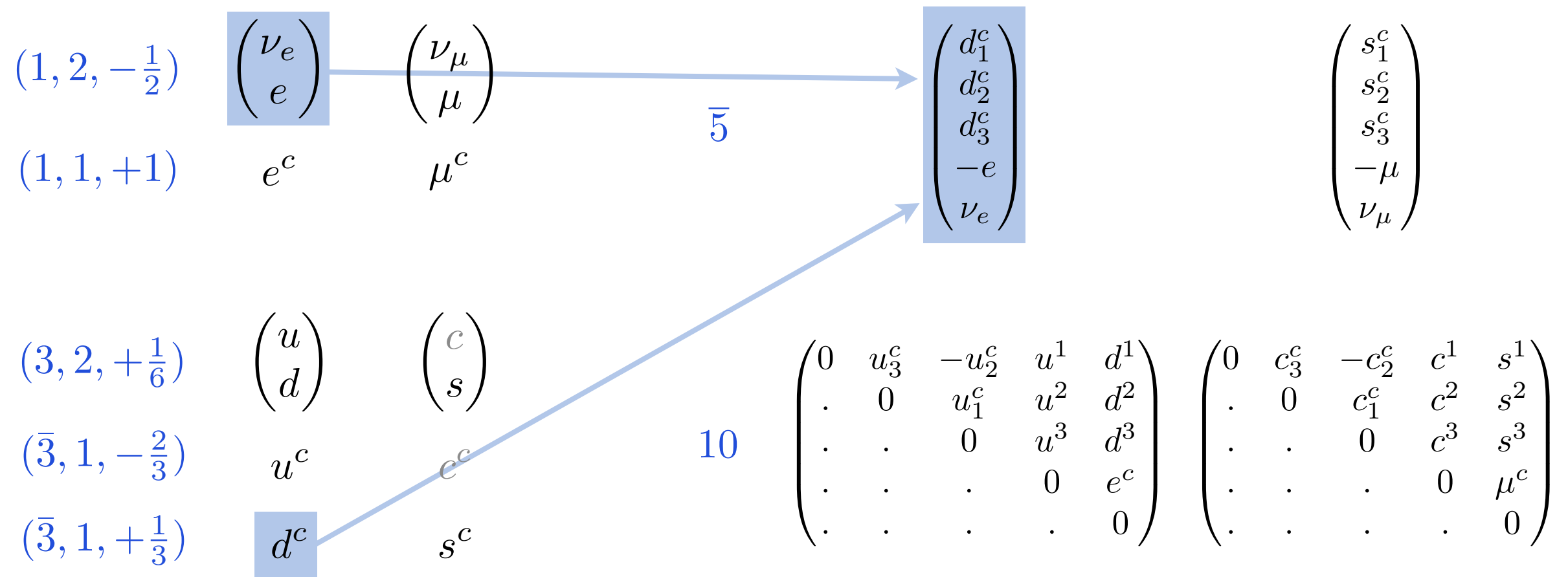
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Gauge sector:

$$\begin{array}{ll}
 (8, 1, 0) & G^\mu \\
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 \end{array}
 \left. \vphantom{\begin{array}{l} (8, 1, 0) \\ (1, 3, 0) \\ (1, 1, 0) \end{array}} \right\} W^\pm, Z, \gamma
 \qquad
 \begin{array}{ccccccc}
 24 = (8, 1, 0) \oplus (1, 3, 0) \oplus (1, 1, 0) \oplus (3, 2, -\frac{5}{6}) \oplus (\bar{3}, 2, +\frac{5}{6}) \\
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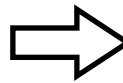
GUT-breaking Higgs: $SU(5) \rightarrow SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$ variety of extra Higgses

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Baryon and lepton number violation in GUTs

- Quarks and leptons share common GUT multiplets
 - gauge bosons coupled to a universal charge
 - Yukawas do not care about who is who either

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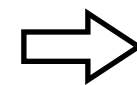
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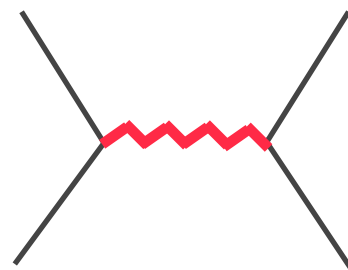
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- Proton decay

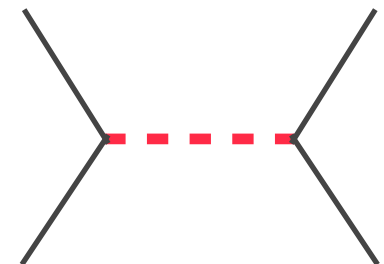
d=6

gauge-induced



$$\frac{f_1}{M_G^2} \bar{Q} u^c \bar{Q} e^c, \quad \frac{f_2}{M_G^2} u^c \bar{Q} d^c \bar{L}$$

Higgs-induced



$$\frac{f_3}{M_G^2} Q Q Q L, \quad \frac{f_4}{M_G^2} u^c u^c d^c e^c$$

SO(10) GUT completion of the L-R model

- all matter in a single IRREP of SO(10)

Fritzsch & Minkowski 1975

$$SO(10) \supset \supset SU(3)_c \otimes SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$$

Q_L

Q_R

L_L

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$$16_F = (3, 2, 1, +1/3) \oplus (3, 1, 2, -1/3) \oplus (1, 2, 1, -1) \oplus (1, 1, 2, +1)$$

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- very simple Yukawa sector

Φ

$$10_H = (3, 1, 1, -2/3) \oplus (3, 1, 1, +2/3) \oplus (1, 2, 2, 0)$$

$$\mathcal{L} \ni 16_F Y_{10} 16_F 10_H$$

all Dirac masses correlated!

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Majorana and extra Dirac masses generated!

SO(10) “flavour miracle”

Assuming dominant type II seesaw:

B.Bajc, G.Senjanovic and F.Vissani
Phys.Rev.Lett. 90 (2003) 051802

$$M_u = Y_{10} v_u^{10} + Y_{126} v_u^{126}$$

$$M_d = Y_{10} v_d^{10} + Y_{126} v_d^{126}$$

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$$\frac{m_b}{m_\tau} \rightarrow 1$$

$$\theta_{23} \sim \frac{\pi}{4}$$

$$\theta_{13} \sim \mathcal{O}(\lambda)$$

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Daya-Bay: $\sin^2 2\theta_{13} = 0.092 \pm 0.016(\text{stat}) \pm 0.005(\text{syst})$

$$U_{e3} \sim 0.15$$

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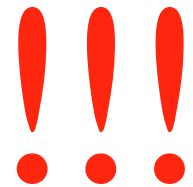
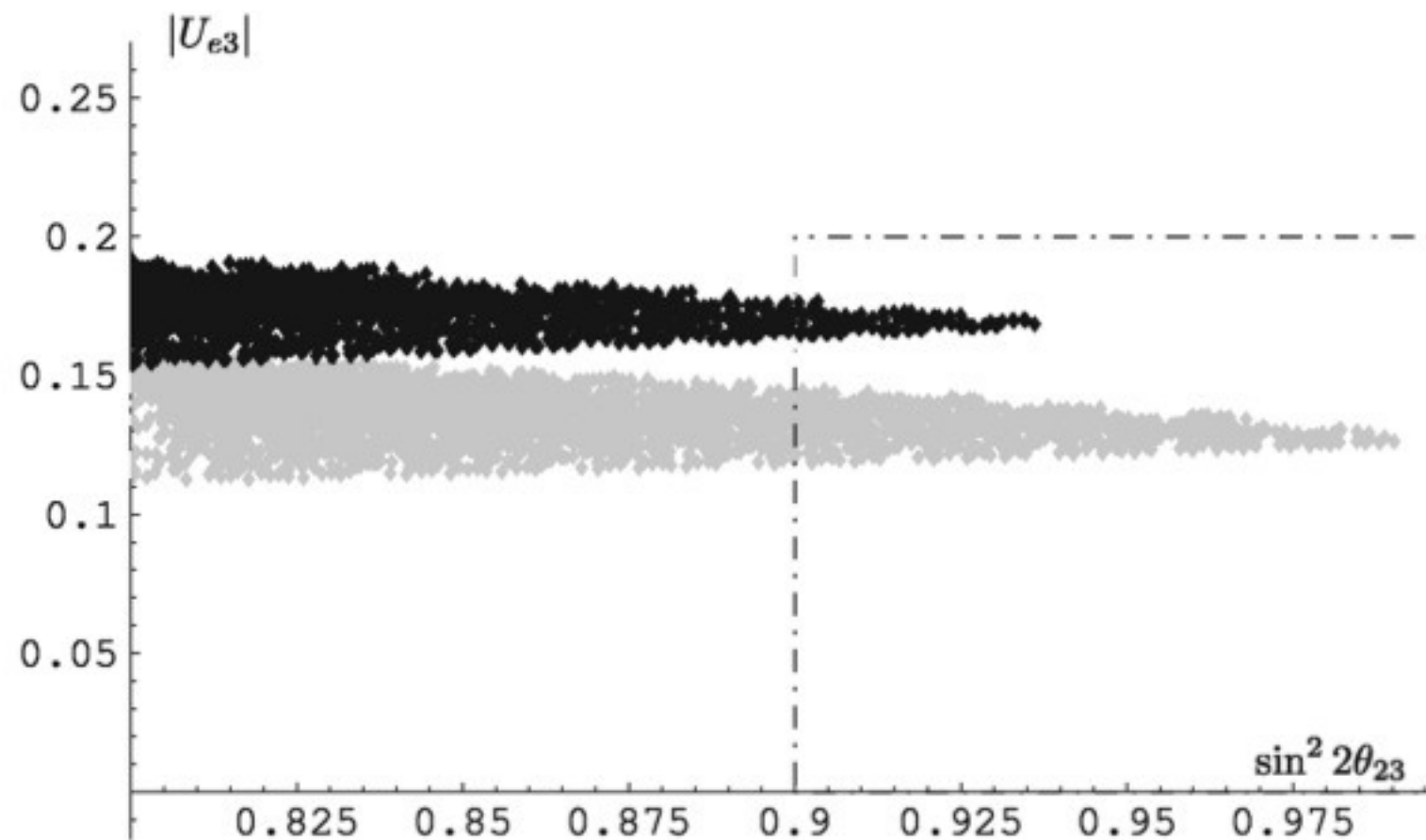


figure from: S. Bertolini, M. Frigerio, MM, Phys. Rev. D70, 2004



Thanks for your kind attention!

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VOLUME 32, NUMBER 8

PHYSICAL REVIEW LETTERS

25 FEBRUARY 1974

Unity of All Elementary-Particle Forces

Howard Georgi* and S. L. Glashow

Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138

(Received 10 January 1974)

Strong, electromagnetic, and weak forces are conjectured to arise from a single fundamental interaction based on the gauge group $SU(5)$.

We present a series of hypotheses and speculations leading inescapably to the conclusion that $SU(5)$ is the gauge group of the world—that all elementary particle forces (strong, weak, and electromagnetic) are different manifestations of the same fundamental interaction involving a single coupling strength, the fine-structure constant. Our hypotheses may be wrong and our speculations idle, but the uniqueness and simplicity of our scheme are reasons enough that it be taken seriously.

of the GIM mechanism with the notion of colored quarks⁴ keeps the successes of the quark model and gives an important bonus: Lepton and hadron anomalies cancel so that the theory of weak and electromagnetic interactions is renormalizable.⁵

The next step is to include strong interactions. We assume that *strong interactions are mediated by an octet of neutral vector gauge gluons associated with local color $SU(3)$ symmetry*, and that there are no fundamental strongly interacting scalar-meson fields.⁶ This insures that

- Actually, what Georgi and Glashow have shown was uniqueness of $SU(5)$ for rank=4 GUTs

Monopoles

No way to produce in lab, only cosmics + Callan-Rubakov effect

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- galactic magnetic field depletion
- pulsar stability
- proton stability

Freese, Turner

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Upper limits on the flux density around Earth

Theory: $\Phi_M(\text{Earth})_{\text{theory}} \lesssim 10^{-22} \sim 10^{-27} \text{ cm}^{-2} \text{ sr}^{-1} \text{ s}^{-1}$

Experiment: $\Phi_M(\text{Earth})_{\text{exp.}} \lesssim 10^{-16} \text{ cm}^{-2} \text{ sr}^{-1} \text{ s}^{-1}$ MACRO 2001 (Gran Sasso)

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N.B. early (fake) monopole-like events Price et al., 1975 PRL August 25