

# Hyperfragments from the lightest $p$ -shell hypernuclei

## Recent progress and the next steps

L. Majling,<sup>1</sup> O. Majlingova<sup>2</sup> and P. Bydžovsky<sup>1</sup>

<sup>1</sup> Nuclear Physics Institute, Řež near Prague

<sup>2</sup> Czech Technical University, Prague

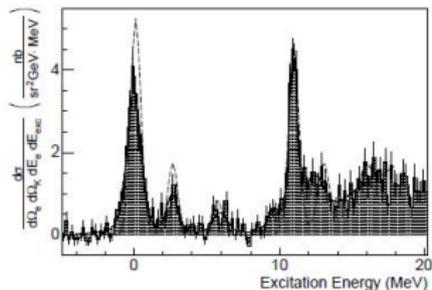
SPHERE Meeting  
Prague 10. IX. 2014

## AGENDA

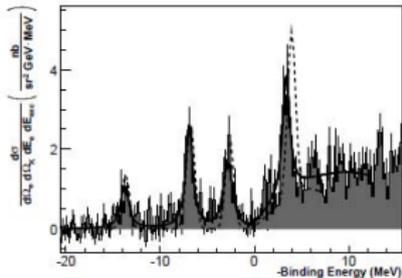
1. Motivation      Experiment **MAMI**, JLab
2. Approach      Shell Model
3. Extension       $N_{\min} + 2$
4. Suggestions

# Example : closed shell nuclei

$^{12}_{\Lambda}B$  : M. Iodice *et al.*, PRL **99**, 052501 (2007)

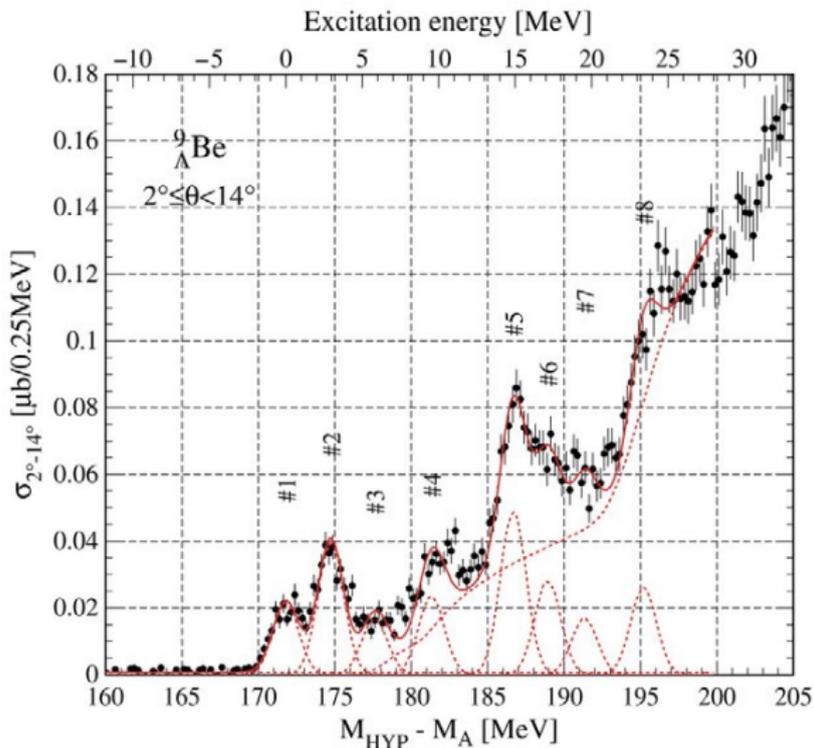


$^{16}_{\Lambda}N$  : F. Cusanno *et al.*, PRL **103**, 202501 (2009)

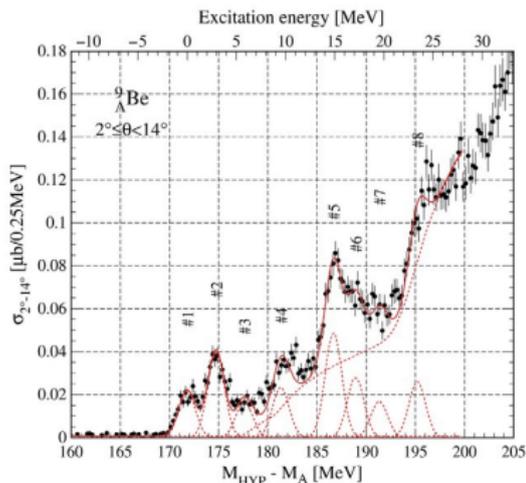


# Example : Spectrum of ${}^9_{\Lambda}\text{Be}$

Spectrum of  ${}^9_{\Lambda}\text{Be}$  measured using the  $(\pi^+, K^+)$  reaction.  
From O. Hashimoto and H. Tamura



# Deciphering spectrum of ${}^9_{\Lambda}\text{Be}$



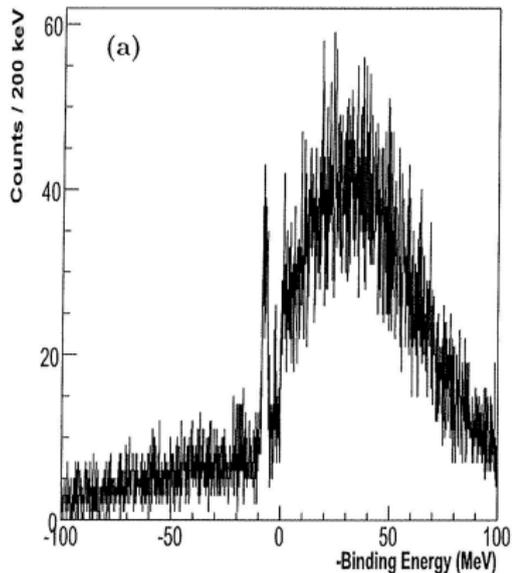
$p^{-1} s_{\Lambda}$      $p_{\Lambda}$   
 $p^{-1} s_{\Lambda}$      $p_{\Lambda}$   
 $s^{-1} s_{\Lambda}$      $s_{\Lambda}$

${}^8\text{Be}$  :

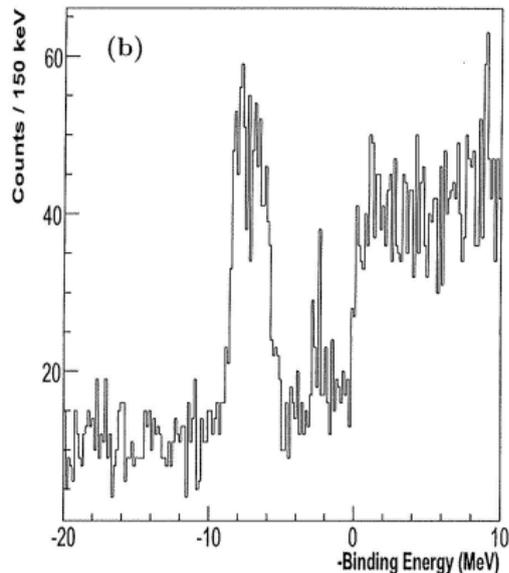
$J^{\pi}; T$	$E_{\text{ex}}$	shell model	"hole"
$0^+; 0$	0.0	$s^4 p^4$ [44] S	$p^{-1}$
$2^+; 0$	3.0	$s^4 p^4$ [44] D	$p^{-1}$
...			
$2^+; 0+1$	16.7	$s^4 p^4$ [431]D	$p^{-1}$
...			
$2^-; 0$	18.9	$s^3 p^5$ [44] P	$s^{-1}$
...			
$(1,2)^-; 1$	24.0	$s^3 p^5$ [431]P	$s^{-1}$

No	$l_n^{-1} l_{\Lambda}$
1,2	$p^{-1} s_{\Lambda}$
3,4	$p^{-1} p_{\Lambda}$
5	$p^{-1} s_{\Lambda}$
6	$s^{-1} s_{\Lambda}$
7	$p^{-1} p_{\Lambda}$
8	$s^{-1} s_{\Lambda}$

# "Our" hypernucleus ${}^9_{\Lambda}\text{Li}$



the whole energy range



the region of interest

Excitation spectrum of  ${}^9_{\Lambda}\text{Li}$  measured using the  $(e, e'K^+)$  reaction  
Jefferson Lab Hall A Collaboration

# Hyperfragments from ${}^9_{\Lambda}\text{Li}$

*Pochodzalla, Acta Phys.Polon. (2011):*

	${}^6_{\Lambda}\text{Li}$ $3n$ 19.0	${}^7_{\Lambda}\text{Li}$ $2n$ 12.2	${}^8_{\Lambda}\text{Li}$ $n$ 3.7	${}^9_{\Lambda}\text{Li}$ $\Lambda$ 8.5
${}^4_{\Lambda}\text{He}$ $tnn$ 31.5	${}^5_{\Lambda}\text{He}$ $tn$ 9.9	${}^6_{\Lambda}\text{He}$ $t$ 9.7	${}^7_{\Lambda}\text{He}$ $d$ 13.0	${}^8_{\Lambda}\text{He}$ $p$ 13.8
${}^3_{\Lambda}\text{H}$ ${}^6\text{He}$ 18.2	${}^4_{\Lambda}\text{H}$ ${}^5\text{He}$ 11.8	$({}^5_{\Lambda}\text{H})$ ${}^4\text{He}$	${}^6_{\Lambda}\text{H}$ ${}^3\text{He}$ 38.5	$({}^7_{\Lambda}\text{H})$ $2p$

Seminal papers

Gal, Soper, Dalitz, AP **63** (1971), **72** (1972) **113** (1978)

$\Psi_{\Lambda}^{(A+1)Z} = \Psi^{(AZ)} \cdot \psi^{\Lambda}$  "weak coupling"

for  $p$ -shell hypernuclei:

$$\underbrace{|s^4 p^{A-4} : J_A^{\pi} T_A}_{\text{Cohen Kurath}} \otimes s_{\Lambda} : J >$$

*Cohen Kurath,*

Discrete part of spectra ( $0\hbar\omega$  excitations) :

Millener, Lecture Notes in Physics, vol. 724, 2007, Springer

EXCITED nuclear states:

LS coupling

$$|0s^{k_s}[f_s] 1p^{k_p}[f_p] \ell^{k_\ell}[f_\ell] : \underline{[f]} \quad \underline{(\lambda\mu)} : J \rangle$$

*Wigner(1939)*

$$\begin{array}{cc} TS & L \\ SU(4) & SU(3) \end{array}$$

*Elliott(1968)*

$$k_s + k_p + k_\ell = A$$

$$\ell = 2d, 2s; \quad 3f, 3p; \dots$$

# Excited states of $p$ -shell nuclei

$0\hbar\omega_N (k) :$

$$|0s^4[4] 1p^k[f_p] \quad : [f_A] (\lambda\mu) > \quad = \quad \Phi_k^{(A)}[f_A](\lambda\mu) \cdot \Psi_0(R_A) \\ [f_A] = [4f_p]$$

$1\hbar\omega_N (k+1) :$

$$\begin{array}{l} |0s^4[4] 1p^{k-1}[f_p] \quad 2d [1] \quad : [f_A] (\lambda\mu) > \\ |0s^3[3] 1p^{k+1}[f'_p] \quad : [f_A] (\lambda\mu) > \end{array} \quad \left. \vphantom{\begin{array}{l} |0s^4[4] 1p^{k-1}[f_p] \quad 2d [1] \\ |0s^3[3] 1p^{k+1}[f'_p] \quad : [f_A] (\lambda\mu) > \end{array}} \right\} \begin{array}{l} \Phi_k^{(A)}[f_A](\lambda\mu) \cdot \Psi_1(R_A) \\ \Phi_{k+1}^{(A)}[f_A](\lambda\mu) \cdot \Psi_0(R_A) \end{array} \\ [f_A] = [4f_p], [3f'_p]$$

$2\hbar\omega_N (k+2) :$

$$\begin{array}{l} |0s^4[4] 1p^{k-1}[f_p] \quad 3f [1] \quad : [f_A] (\lambda\mu) > \\ |0s^4[4] 1p^{k-2}[f'_p] \quad 2d^2[f_d] \quad : [f_A] (\lambda\mu) > \\ |0s^3[3] 1p^k [f''_p] \quad 2d [1] \quad : [f_A] (\lambda\mu) > \\ |0s^2[2] 1p^{k+2}[f'''_p] \quad : [f_A] (\lambda\mu) > \end{array} \quad \begin{array}{l} \Phi_k^{(A)}[f_A](\lambda\mu) \cdot \Psi_2(R_A) \\ \Phi_k^{(A)}[f_A](\lambda\mu) \cdot \Psi_2(R_A) \\ \Phi_{k+1}^{(A)}[f_A](\lambda\mu) \cdot \Psi_1(R_A) \\ \Phi_{k+2}^{(A)}[f_A](\lambda\mu) \cdot \Psi_0(R_A) \end{array} \\ [f_A] = [4f_p], [3f'_p], [2f''_p]$$

Due to NON-CENTRAL forces,  
 the wave function for the  ${}^4\text{He } J^\pi = 0^+$  ground state  
 is a mixture of  ${}^1S_0$ ,  ${}^3P_0$ , and  ${}^5D_0$  states

Recent calculations:

ref.	interaction	${}^1S_0$	${}^3P_0$	${}^5D_0$
[1]	AV8'	85.7 %	0.4 %	13.9 %
[2]	AV18 + UIX	83.2 %	0.8 %	16.0 %

[1] PR C 64, 044001 ('01) K. Kamada *et al.*:  
*Benchmark test calculation of a 4N bound state*

[2] PR C 65 054003 ('02) A. Nogga *et al.*:  
*The  $\alpha$  particle based on modern forces*

$${}^{15}D_0: |0s^2 1p^2 : [22] \rangle$$

# Excited states of $p$ -shell HYPERNUCLEI:

$$\hbar\omega : \hbar\omega_k + \hbar\omega_\Lambda$$

0	0 + 0	$\Phi_k^{(A)} \cdot \varphi_0^\Lambda(R_A - r_\Lambda)$	$(\varphi_0^\Lambda \equiv 0s_{\frac{1}{2}}^\Lambda)$
1	0 + 1	$\Phi_k^{(A)} \cdot \varphi_1^\Lambda$	$(\varphi_1^\Lambda \equiv 1p_{\frac{3}{2}}^\Lambda, 1p_{\frac{1}{2}}^\Lambda)$
	1 + 0	$\Phi_{k+1}^{(A)} \cdot \varphi_0^\Lambda$	
2	0 + 2	$\Phi_k^{(A)} \cdot \varphi_2^\Lambda$	$(\varphi_2^\Lambda \equiv 2d_{\frac{5}{2}}^\Lambda, 2d_{\frac{3}{2}}^\Lambda, 2s_{\frac{1}{2}}^\Lambda)$
	1 + 1	$\Phi_{k+1}^{(A)} \cdot \varphi_1^\Lambda$	
	2 + 0	$\Phi_{k+2}^{(A)} \cdot \varphi_0^\Lambda$	

# Example : structure $(\begin{smallmatrix} 7 \\ \Lambda \end{smallmatrix} \text{He})$

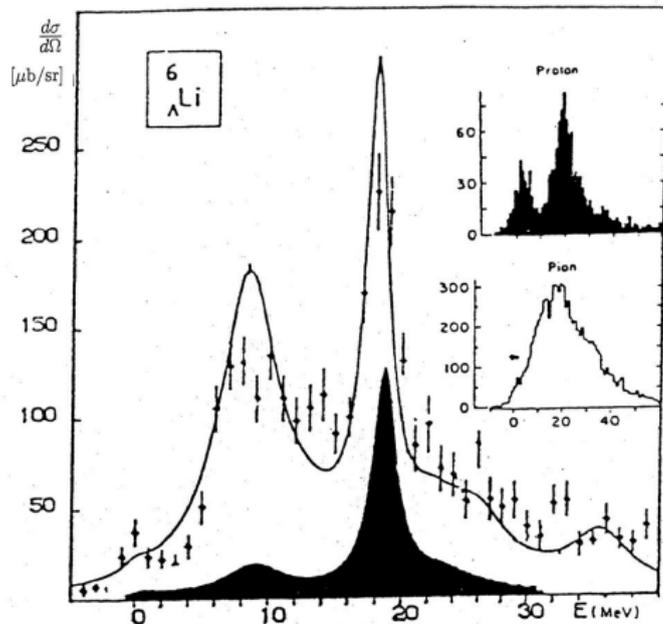
N	${}^6\text{He}$	$\Phi_N^{(6)}$	$[f_6]$	$\begin{smallmatrix} 7 \\ \Lambda \end{smallmatrix} \text{He}$
2	$s^4 p^2$	$\Phi_2^{(6)}$	$[42], \dots$	$s_\Lambda \quad p_\Lambda \quad d_\Lambda$
3	$s^4 p d$	$\Phi_3^{(6)}$	$[42], \dots$	$s_\Lambda \quad p_\Lambda$
3	$s^3 p^3$	$\Phi_3^{(6)}$	$[42], \dots \quad [ \underline{\mathbf{3}} \mathbf{3} ], \dots$	$s_\Lambda \quad p_\Lambda$
4	$s^4 p f$	$\Phi_4^{(6)}$	$[42], \dots$	$s_\Lambda$
4	$s^4 d^2$	$\Phi_4^{(6)}$	$[42], \dots$	$s_\Lambda$
4	$s^3 p^2 d$	$\Phi_4^{(6)}$	$[42], \dots \quad [33], \dots$	$s_\Lambda$
4	$s^2 p^4$	$\Phi_4^{(6)}$	$[42], \dots \quad [33], \dots \quad [ \underline{\mathbf{2}} \mathbf{22} ]$	$s_\Lambda$

# Spectrum of ${}^6_{\Lambda}\text{Li}$

deciphering:

$$p^{-1}s_{\Lambda} \quad p^{-1}p_{\Lambda} \quad s^{-1}s_{\Lambda}$$

$$[41] \quad [41] \quad [32]$$



The  $(K^-, \pi^-)$  spectrum of the  ${}^6_{\Lambda}\text{Li}$  production.

# Cluster wave function in Shell Model

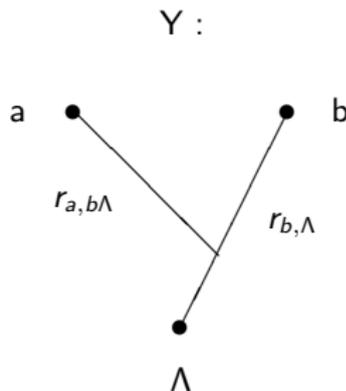
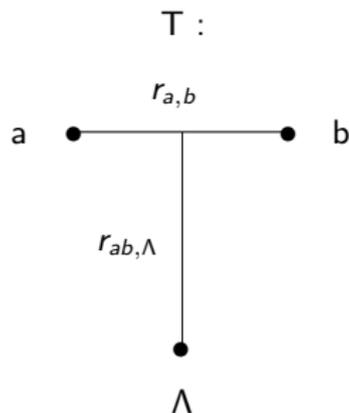
$$\Phi_{NA}^{(A)}[f_A](\lambda\mu)_A = \sum g^a g^b \phi_{na}^{(a)}[f_a](\lambda\mu)_a \phi_{nb}^{(b)}[f_b](\lambda\mu)_b \varphi_\nu(r_{a,b})$$

$g^a(g^b)$ : coefficient of fractional parentage for a (b) particles

constrains:

$$\begin{aligned} na + nb + \nu &= N_A \\ [f_a] + [f_b] &= [f_A] \\ (\lambda\mu)_a + (\lambda\mu)_b + (\nu 0) &= (\lambda\mu)_A \\ L_a + L_b &= L_A \end{aligned}$$

# Hyper FRAGMENTS



Transformation of Jacobi coordinates:

$$\langle \varphi_\nu(r_{ab}) \varphi_0^\Lambda(r_{ab,\Lambda}) | \varphi_0(r_{b,\Lambda}) \varphi_\nu(r_{a,b\Lambda}) \rangle = \left( \frac{m_A+1}{m_A} \frac{m_b}{m_b+1} \right)^{\frac{\nu}{2}}$$

$E_{\text{th}}$		$[f_c]$
2.82	${}^6_{\Lambda}\text{He} + n$	[41]
3.08	${}^5_{\Lambda}\text{He} + 2n$	[4]
5.23	${}^6\text{He} + \Lambda$	[42]
6.95	${}^5\text{He} + \Lambda n$	[41]
15.49	${}^4_{\Lambda}\mathbf{H} + t$	[3]
21.41	${}^4_{\Lambda}\text{He} + 3n$	[3]
23.66	${}^3_{\Lambda}\mathbf{H} + tn$	[2]
23.81	${}^6_{\Lambda}\mathbf{H} + p$	[32]

# Hyperfragments from ${}^9_{\Lambda}\text{Li}$

$E_{\text{thr}}$	decay		$[f_i][f_k]$	$T_1 T_2$
3.7	$n$	$+ {}^8_{\Lambda}\text{Li}$	[43][1]	$\frac{1}{2} \frac{1}{2}$
8.5	${}^8\text{Li}$	$+ \Lambda$	[431]	1 0
9.7	$t$	$+ {}^6_{\Lambda}\text{He}$	[3][41]	$\frac{1}{2} \frac{1}{2}$
		$\downarrow$		
9.9	$tn$	$+ {}^5_{\Lambda}\text{He}$	[3][1][4]	$\frac{1}{2} \frac{1}{2} 0$
11.8	${}^5\text{He}$	$+ {}^4_{\Lambda}\text{H}$	[41][3]	$\frac{1}{2} \frac{1}{2}$
12.2	$2n$	$+ {}^7_{\Lambda}\text{Li}$	[2][42]	1 0
13.0	$d$	$+ {}^7_{\Lambda}\text{He}$	[2][42]	0 1
13.8	$p$	$+ {}^8_{\Lambda}\text{He}$	[1][421]	$\frac{1}{2} \frac{3}{2}$
18.2	${}^6\text{He}$	$+ {}^3_{\Lambda}\text{H}$	[42][2]	1 0
31.5	$tnn$	$+ {}^4_{\Lambda}\text{He}$	[3][1][1][3]	$\frac{1}{2} \frac{1}{2} \frac{1}{2} \frac{1}{2}$
38.5	${}^3\text{He}$	$+ {}^6_{\Lambda}\text{H}$	<b>[3][32]</b>	$\frac{1}{2} \frac{3}{2}$

Some SUGGESTIONS for next steps:

- see for  ${}^4_{\Lambda}\text{H}$  in other  $p$ -shell targets;
- see for  ${}^3_{\Lambda}\text{H}$  in  $p$ -shell targets;
- ${}^7\text{Li}$  is a source of  ${}^6_{\Lambda}\text{H}$  hyperfragment.