

The No-Core Shell Model and High-Performance Computing

Pieter Maris

Dept. of Physics and Astronomy
Iowa State University
Ames, IA 50011
pmaris@iastate.edu

28th Indian Summer School
Ab Initio methods in Nuclear Physics
Aug. 29 – Sept. 2, 2016, Prague, Czech Republic

Outline

Introduction

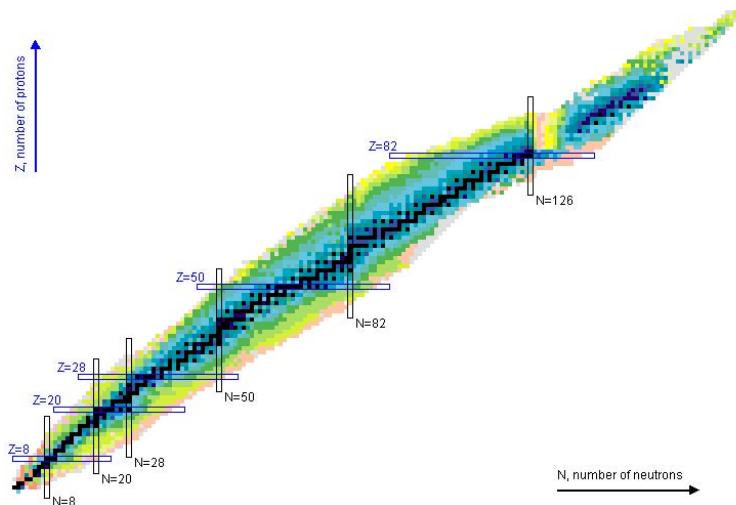
No-Core Configuration Interaction approach

High-Performance Computing

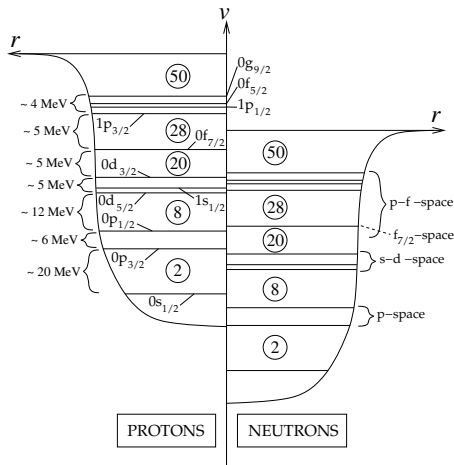
CI calculations on HPC platforms

For further reading

Nuclear Landscape: Valley of stability & Magic numbers

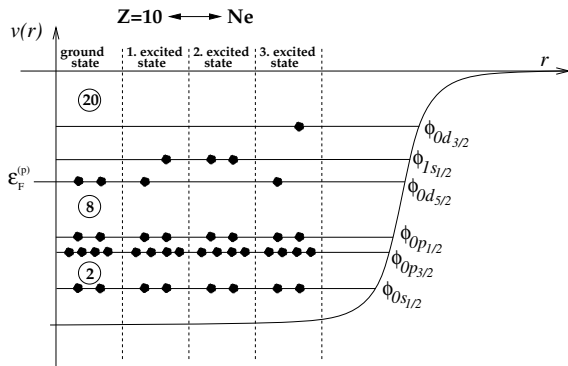


The nuclear mean field Shell Model



- Effective (phenomenological) potential for protons and neutrons

The nuclear mean field Shell Model



- ▶ Closed shell core
- ▶ Effective potential for non-interacting valence nucleons:
Single-Particle energies ϵ
- ▶ Interactions between valence nucleons:
Configuration mixing

The nuclear No-Core Shell Model

Given a Hamiltonian operator (in relative coordinates)

$$\hat{\mathbf{H}}_{\text{rel}} = \sum_{i < j} \frac{(\vec{p}_i - \vec{p}_j)^2}{2 m A} + \sum_{i < j} V_{ij} + \sum_{i < j < k} V_{ijk} + \dots$$

solve the eigenvalue problem for the bound state wavefunction of A nucleons

$$\hat{\mathbf{H}}_{\text{rel}} \Psi(r_1, \dots, r_A) = \lambda \Psi(r_1, \dots, r_A)$$

by expanding the $\Psi(r_1, \dots, r_A)$ in products of harmonic oscillator (H.O.) single-particle states

The No-Core Shell Model: Plan for four lectures

1. NCSM and HPC
 - ▶ Remainder of this lecture
2. NN and 3N (effective) interactions
 - ▶ This afternoon
3. Emergence of collective motion
(plus additional topics, depending on interests?)
 - ▶ Tomorrow
4. Beyond the H.O. basis
(plus additional topics, depending on interests?)
 - ▶ Wednesday

Reminder: 3D quantum Harmonic Oscillator

- ▶ Hamiltonian $\hat{\mathbf{H}}^{\text{H.O.}} = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 r^2$
- ▶ Solutions of Schrödinger equation $\hat{\mathbf{H}}^{\text{H.O.}}\phi(r) = E_N\phi(r)$

$$\phi(r) = \mathcal{N}_{nl} \left(\frac{r}{b}\right)^l e^{-r^2/(2b^2)} L_n^{(l+\frac{1}{2})}(r^2/b^2) Y_{lm}(\theta, \phi)$$

with H.O. parameter $b^2 = \hbar/(m\omega)$

- ▶ Eigenvalues $E_N = (\frac{3}{2} + N)\hbar\omega$ with $N = 2n + l$
- ▶ Nucleons: spin- $\frac{1}{2}$ particles, spin up or down
- ▶ Single-particle quantum number in LS-scheme: $|nlmss_z\rangle$
- ▶ Couple orbital motion and intrinsic spin (and suppress $s = \frac{1}{2}$) to J-scheme: $\hat{\mathbf{J}} = \hat{\mathbf{L}} + \hat{\mathbf{S}}$: $|nljm_j\rangle$
- ▶ Degeneracy of major (H.O. energy) shells $(N+1)(N+2)$
 - ▶ magic numbers 2, 8, 20, 40, ...
 - ▶ s-shell, p-shell, sd-shell, pf-shell, ...

No-Core Configuration Interaction approach

- ▶ Expand wavefunction in basis states $|\Psi\rangle = \sum a_i |\Phi_i\rangle$
- ▶ Express Hamiltonian in basis $\langle \Phi_j | \hat{\mathbf{H}} | \Phi_i \rangle = H_{ij}$
- ▶ Diagonalize Hamiltonian matrix H_{ij}
- ▶ No-Core: all A nucleons are treated the same
- ▶ Complete basis \longrightarrow exact result
 - ▶ caveat: complete basis is infinite dimensional
- ▶ In practice
 - ▶ truncate basis
 - ▶ study behavior of observables as function of truncation
- ▶ Computational challenge
 - ▶ construct large ($10^{10} \times 10^{10}$) sparse symmetric matrix H_{ij}
 - ▶ obtain lowest eigenvalues & -vectors corresponding to low-lying spectrum and eigenstates

NCCI approach – Basis expansion

- ▶ Expand A -body wave function in basis functions

$$\Psi(r_1, \dots, r_A) = \sum a_i \Phi_i(r_1, \dots, r_A)$$

- ▶ Use basis of single Slater Determinants of single-particle states

$$\Phi_i(r_1, \dots, r_A) = \frac{1}{\sqrt{(A!)}} \begin{vmatrix} \phi_{i1}(r_1) & \phi_{i2}(r_1) & \dots & \phi_{iA}(r_1) \\ \phi_{i1}(r_2) & \phi_{i2}(r_2) & \dots & \phi_{iA}(r_2) \\ \vdots & \vdots & & \vdots \\ \phi_{i1}(r_A) & \phi_{i2}(r_A) & \dots & \phi_{iA}(r_A) \end{vmatrix}$$

which takes care of anti-symmetrization (Fermion statistics)

- ▶ Each many-body basis state $\Phi(r_1, \dots, r_A)$ eigenstate of $\hat{\mathbf{J}}_z$, with fixed eigenvalue M , but not of $\hat{\mathbf{J}}^2$ (M -scheme)
 - ▶ almost trivial to implement

$$\hat{\mathbf{J}}_z |\Phi\rangle = M |\Phi\rangle = \sum_{k=1}^A m_{j_k} |\Phi\rangle$$

NCCI approach – Hamiltonian

- ▶ Expand Hamiltonian in basis

$$\hat{\mathbf{H}}_{\text{rel}} = \sum_{i < j} \frac{(\vec{p}_i - \vec{p}_j)^2}{2 m A} + \sum_{i < j} V_{ij} + \sum_{i < j < k} V_{ijk} + \dots$$

- ▶ A -body problem with a -body interaction: nonzero matrix elements iff at least $(A - a)$ particles are in identical single-particle states
 - ▶ many-body basis states are single Slater Determinants
 - ▶ A -body problem with 2-body interactions

$$H_{ij}^{(A)} = (-1)^{\text{permutations}} \delta_{i_1, j_1} \dots \delta_{i_{(A-2)}, j_{(A-2)}} \langle ab | \hat{\mathbf{H}} | cd \rangle$$

- ▶ Sparse symmetric eigenvalue problem

NCCI approach – Truncations

Complete space – exact (for bound states), but infinite-dimensional

- ▶ N_{\max} truncation:

Each many-body basis state $\Phi(r_1, \dots, r_A)$ satisfies

$$\sum_{k=1}^A (2 n_k + l_k) \leq N_0 + N_{\max}$$

- ▶ exact factorization of Center-of-Mass motion
- ▶ Alternatives:
 - ▶ FCI (commonly used in nuclear shell model, quantum chemistry, ...) truncation on single-particle basis states only
 - ▶ Importance Truncation
 - ▶ No-Core Monte-Carlo Shell Model
 - ▶ SU(3) Truncation
 - ▶ ...

Roth, PRC79, 064324 (2009)

Abe *et al*, PRC86, 054301 (2012)

Dytrych *et al*, PRL111, 252501 (2013)

Intermezzo: Center-of-Mass motion

- ▶ Use single-particle coordinates, not relative (Jacobi) coordinates
 - ▶ straightforward to extend to many particles
 - ▶ have to separate Center-of-Mass motion from relative motion
- ▶ Center-of-Mass wavefunction factorizes for H.O. basis functions in combination with N_{\max} truncation

$$\begin{aligned} |\Psi_{\text{total}}\rangle &= |\phi_1\rangle \otimes \dots \otimes |\phi_A\rangle \\ &= |\Phi_{\text{Center-of-Mass}}\rangle \otimes |\Psi_{\text{rel}}\rangle \end{aligned}$$

where

$$\hat{\mathbf{H}}_{\text{rel}} |\psi_{i, \text{rel}}\rangle = E_i |\psi_{i, \text{rel}}\rangle$$

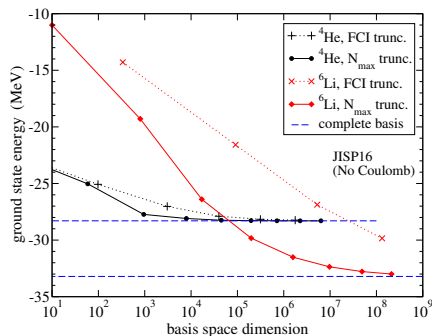
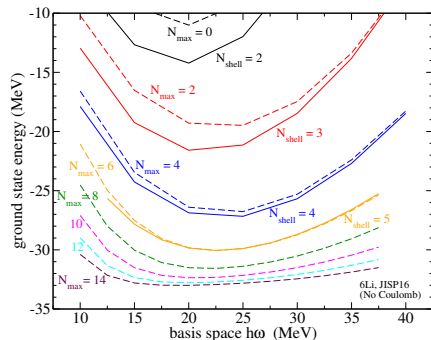
- ▶ Add Lagrange multiplier to Hamiltonian (Lawson term)

$$\hat{\mathbf{H}}_{\text{rel}} \longrightarrow \hat{\mathbf{H}}_{\text{rel}} + \Lambda_{\text{CM}} \left(\hat{\mathbf{H}}_{\text{CM}}^{\text{H.O.}} - \frac{3}{2} \hbar \omega \right)$$

with $\hat{\mathbf{H}}_{\text{rel}} = \hat{\mathbf{T}}_{\text{rel}} + \hat{\mathbf{V}}_{\text{rel}}$ the relative Hamiltonian

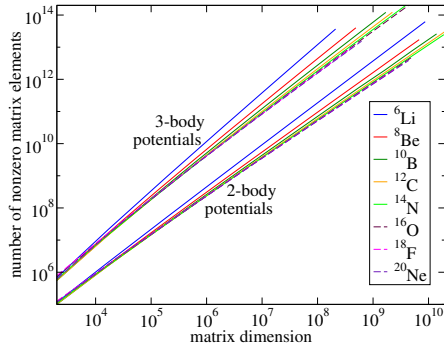
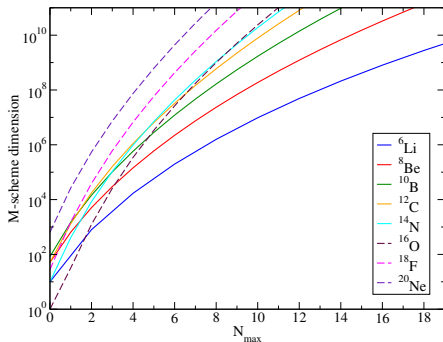
- ▶ separates states with CM excitations from states with 0s CM motion
 $|\Phi_{\text{CM}}\rangle = |\Phi_{0s}\rangle$

Intermezzo: FCI vs. N_{\max} truncation



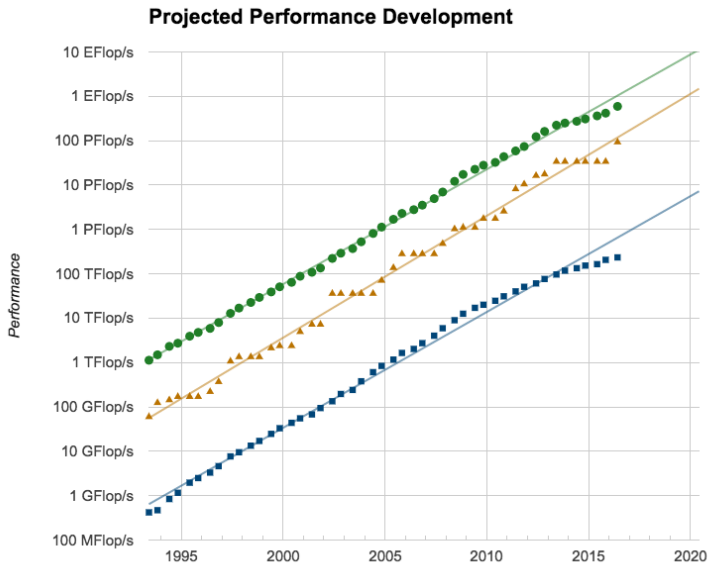
- ▶ N_{\max} truncation
 - ▶ exact factorization of Center-of-Mass motion
- ▶ Infinite basis space limit
 - ▶ both N_{\max} truncation and FCI converge to the same results
 - ▶ N_{\max} truncation does so much more rapidly

NCCI approach – Main Challenge



- ▶ Increase of basis space dimension with increasing A and N_{\max}
 - ▶ need calculations up to at least $N_{\max} = 8$ for meaningful extrapolation and numerical error estimates
- ▶ More relevant measure for computational needs
 - ▶ number of nonzero matrix elements
 - ▶ current limit 10^{13} to 10^{14} (Edison, Mira, Titan)

High-Performance Computing: Moore's law



High-Performance Computing

- ▶ Parallel computing
 - ▶ Initially: Shared memory **or** distributed memory parallel systems
 - ▶ Currently: Systems have shared **and** distributed memory
 - ▶ use OpenMP within a node **and** MPI between nodes
- ▶ Accelerators
 - ▶ GPU's (NVIDIA), Xeon Phi (Intel), ...
 - ▶ Initially: as co-processor
 - ▶ Soon: self-hosted
- ▶ Vectorization
 - ▶ Xeon Phi (KNL) has 512-bit vector units (8 double precision floats)
- ▶ Increasing performance gap between processor and memory
 - ▶ Available memory and memory bandwidth per PU decreases
 - ▶ Data locality and data placement is crucial

Highly nontrivial to achieve good performance

- ▶ Need to collaborate with applied mathematicians and computer scientists

Many-Fermion Dynamics for nuclear structure

Configuration Interaction code for nuclear structure calculations

- ▶ Platform-independent, hybrid OpenMP/MPI, Fortran 90
- ▶ Generate many-body basis space
subject to user-defined single-particle and many-body truncation
- ▶ Construct of many-body matrix H_{ij}
 - ▶ determine which matrix elements can be nonzero
based on quantum numbers of underlying single-particle states
 - ▶ evaluate and store nonzero matrix elements
in compressed sparse block format
- ▶ Obtain lowest eigenpairs using Lanczos algorithm (or LOBPCG)
 - ▶ typical use: 8 to 12 lowest eigenvalues and eigenvectors
 - ▶ typically need ~ 300 to ~ 500 Lanczos iterations
 - ▶ some applications need hundreds of eigenvalues
- ▶ Write eigenvectors (wavefunctions) to disk
- ▶ Calculate selected set of observables

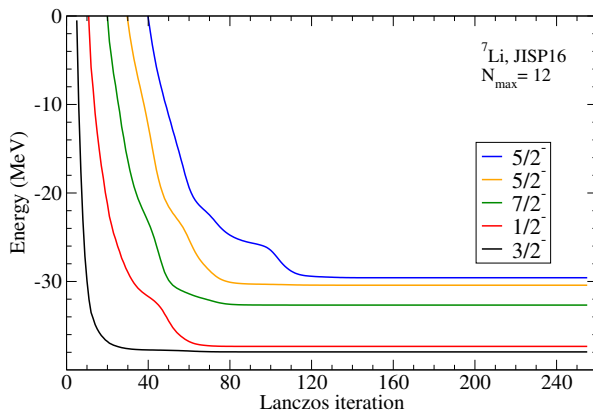
Lanczos algorithm

as implemented in MFDn

Let H be a symmetric matrix. Then H can be reduced to a symmetric tridiagonal matrix T via orthogonal unitary transformations, $H = Q_n T_n Q_n^T$

- ▶ For $i = 1$, set $\beta_1 = 0$ and initial vector q_1 with $\|q_1\| = 1$
- ▶ While (not converged) do
 1. compute $p = H q_i$ i.e. perform Sparse Matrix-Vector Multiplication
 2. compute $\alpha_i = q_i^T \cdot (H q_i)$ i.e. perform dot-product
 3. compute $k = p - \alpha_i q_i - \beta_i q_{i-1}$
 4. (orthogonalize k w.r.t. q_i for numerical stability) more dot-products
 5. compute $\beta_{i+1} = \|k\|$ and one more dot-products
 6. set $q_{i+1} = k / \|k\|$
 7. increment $i = i + 1$
 8. check (convergence) diagonalize small tridiagonal matrix
 - ▶ obtain eigen-values λ and -vectors v of T_n LAPACK
 - ▶ compute $\beta_i |\frac{v_i}{\lambda}|$ for each desired eigenvalue
- ▶ Compute approximate eigenvectors of H from T_n and Q_n

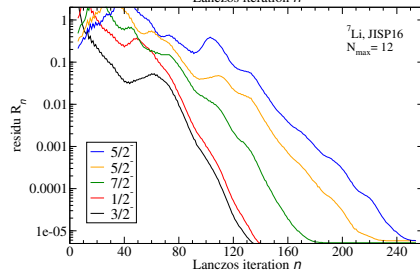
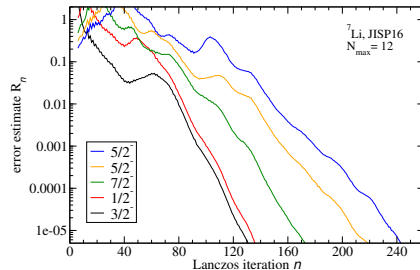
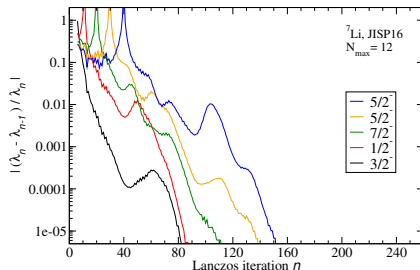
Lanczos algorithm



- ▶ dimension 252 million, with 400 billion nonzero matrix elements
- ▶ runs on 124 nodes Edison at NERSC using 496 MPI ranks with 6 OpenMP threads/MPI
- ▶ total runtime less than 10 minutes

- ▶ Lowest 5 eigenvalues of T_n after n Lanczos iterations
- ▶ Note: in MFDn we use single-precision for H and $Q = \{q_i\}$ but double-precision for dot-products and T_n for numerical stability

Lanczos algorithm – convergence criterium



- Error estimate (cheap)

$$R_n \approx \beta_n \left| \frac{v_n}{\lambda} \right|$$

- Residu (expensive)

$$R^2 = \frac{\sum_i \left(\sum_j H_{ij} a_j - \lambda a_i \right)^2}{\lambda^2}$$

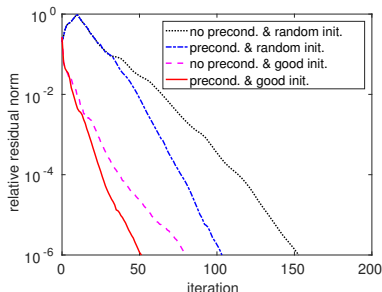
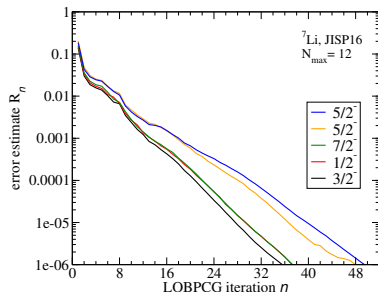
LOBPCG

in collaboration with applied mathematicians from Berkeley

- ▶ Set **initial guess** for $X^{(1)}$ consisting of k orthonormal vectors
 - ▶ ideally, consisting of approximate eigenvectors
 - ▶ e.g. smaller basis space, different H.O. parameter $\hbar\omega$, ...
- ▶ While (not converged) do
 1. apply **preconditioner T**
 - ▶ preconditioning is an art ...
 - ▶ kinetic energy is likely to be efficient, but too expensive
 - ▶ diagonal matrix element is cheap, but not efficient
 - ▶ compromise:
diagonal tiles of H , based on many-body (n, l, j) orbitals
 2. orthonormalize using Cholesky QR
 3. compute $H X^{(i)}$ Sparse Matrix-Matrix Multiplication
 4. do LOBPCG magic ...
 5. check convergence
- ▶ $X^{(n)}$ consists of k orthonormal eigenvectors

LOBPCG

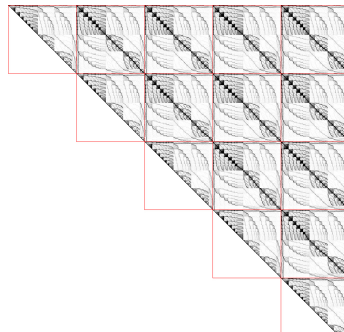
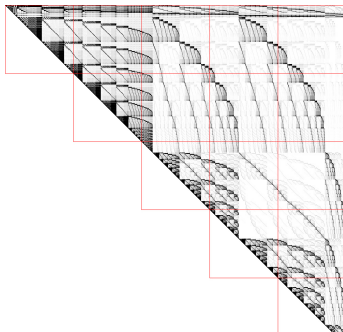
Shao, Aktulga, Yang, Ng, PM, Vary, to be submitted



- ▶ Blocks of 8 (12) vectors, targeting lowest 5 (8) eigenstates
- ▶ $N_{\max} = 8$: 114 iterations in 6.5 seconds (using random initial vectors)
- ▶ $N_{\max} = 10$: 67 iterations in 19.8 seconds
- ▶ $N_{\max} = 12$: 50 iterations in 109.4 seconds
- ▶ Despite doing approximately 1.6 times more work in SpMV/SpMM, **LOBPCG factor of 2 faster than Lanczos**

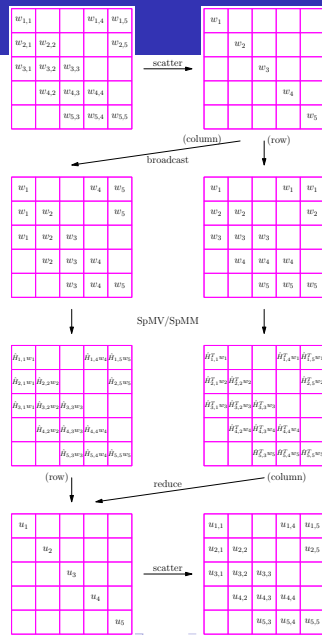
Efficient distributed SpMV

- ▶ Symmetric matrix
 - ▶ store only half the matrix (upper or lower triangle)
 - ▶ have to do SpMV and SpMV^T with same data structure
- ▶ Load-balancing
 - ▶ load determined by number of nonzero matrix elements
 - ▶ 2-dimensional distribution of matrix
 - ▶ round-robin distribution of (groups of) many-body states



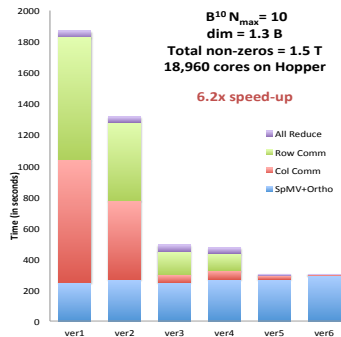
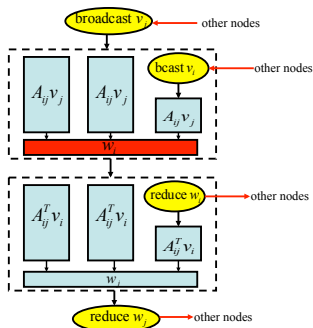
Efficient distributed SpMV

- Communication needs to be load-balanced as well
- Vectors distributed over all processors for orthogonalization



Efficient distributed SpMV – MPI communication

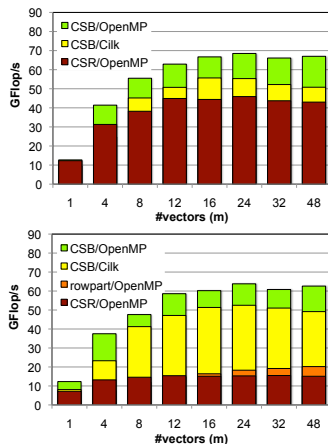
Aktulga, Yang, Ng, PM, Vary, *Concurr. Comput.* 26 (2014), doi:10.1002/cpe.3129



- ▶ Overlap communication with computation
- ▶ Optimize mapping onto network topology for non-overlapping communication

see also Oryspayev, PhD thesis 2016, ISU

Efficient distributed SpMV – single-node performance



- ▶ Compressed sparse block (CSB)
 - ▶ improves data locality and cache performance
 - ▶ allows for efficient OpenMP parallelization within nodes, avoiding race conditions for both SpMV and SpMV^T
- ▶ Block algorithm (LOBPCG)
 - ▶ SpMV on 'block of vectors' allows for vectorization
- ▶ Performance tuning for KNL in progress in collaboration with NERSC, Cray, and Intel

Aktulga, Afibuzzaman, Williams, Buluç, Shao, Yang, Ng, PM, Vary, submitted for publication

For further reading

- ▶ J. Suhonen,
From Nucleons to Nucleus,
Springer, 2007
- ▶ R.R. Whitehead, A. Watt, B.J. Cole and I. Morrison,
Computational Methods for Shell Model Calculations,
Adv. Nucl. Phys. **9**, 123 (1977).
- ▶ B.R. Barrett, P. Navrátil and J.P. Vary,
Ab initio no core shell model,
Prog. Part. Nucl. Phys. **69** (2013) 131.

For further reading – CI codes

- ▶ OXBASH/NuShellX

B.A. Brown, A. Etchegoyen and W.D.M. Rae, *Computer code OXBASH: the Oxford University-Buenos Aires-MSU shell model code*, MSU Cyclotron Laboratory Report No. 524, 1985;
B.A. Brown and W.D.M. Rae. Nucl. Data Sheets 120, 115 (2014).

- ▶ Antoine

F. Nowacki and E. Caurier, Acta Phys. Pol. **B 30**, 705 (1999);
E. Caurier *et al.*, Rev. Mod. Phys. **77**, 427 (2005).

- ▶ Bigstick

C.W. Johnson, W.E. Ormand and P.G. Krastev,
Comp. Phys. Comm. **184**, 2761 (2013) [arXiv:1303.0905 [nucl-th]];
H. Shan, S. Williams, C.W. Johnson, K. McElvain and
W.E. Ormand, SuperComputing 2015.

- ▶ MFDn

should be made public this fall/winter ...