

NN and 3N (effective) interactions for ab initio nuclear structure calculations

Pieter Maris

Dept. of Physics and Astronomy
Iowa State University
Ames, IA 50011
pmaris@iastate.edu

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Outline

Phenomenological NN and 3N interactions

Nuclear interactions from Chiral Effective Field Theory

Similarity Renormalization Group

Okubo–Lee–Suzuki renormalization

Appendix: From (effective) interactions to H.O. matrix elements in single-particle coordinates

For further reading

Phenomenological nuclear interactions

$$\hat{\mathbf{H}}_{\text{rel}} = \hat{\mathbf{T}}_{\text{rel}} + \sum_{i < j} V_{ij} + \sum_{i < j < k} V_{ijk} + \dots$$

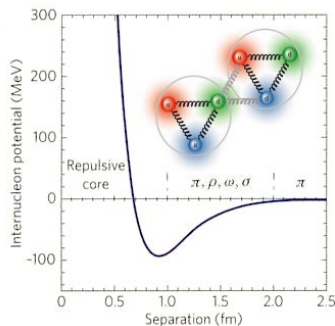
Nuclear interaction not well-determined

- ▶ In principle calculable from QCD
- ▶ Constrained by experimental (scattering) data

Alphabet of realistic NN potentials

- ▶ Argonne potentials
- ▶ Bonn potentials
- ▶ Chiral interactions
- ▶ Daejeon16
- ▶ ...

Most need 3N forces as well



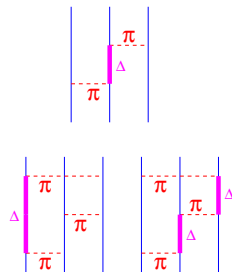
NN potential and scattering data

- ▶ Experimental cross-section data for pp and pn scattering, but not for nn scattering
 - ▶ analysis in terms of isoscalar $T = 0$ and isovector $T = 1$ channels
- ▶ Typically, cross-section data converted to phase shifts
- ▶ NN potentials fitted to phase shifts
 - ▶ propagation of experimental uncertainties?
 - ▶ fitted up to what energy?
- ▶ NN scattering data constrain only the on-shell NN potential, but not the off-shell behavior
 - ▶ many NN potentials describe NN scattering data equally well
- ▶ Additional (physics) input
 - ▶ meson exchange currents
 - ▶ chiral effective field theory
 - ▶ select observables from light nuclei (which?)
 - ▶ more or less suitable for intended computational framework (e.g. local vs. nonlocal interactions)

Argonne V18

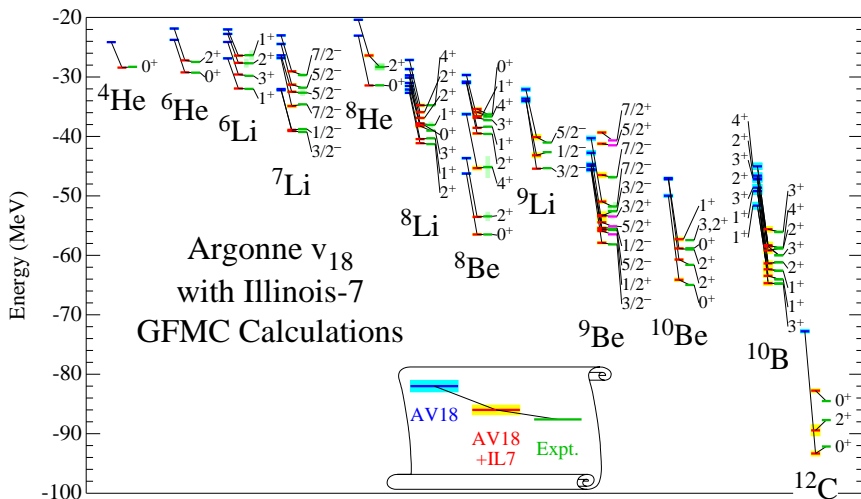
Wiringa, Stoks and Schiavilla, PRC 51, 38 (1995)

- ▶ Accurate local NN potential
 - ▶ 14 charge-independent operators (AV14)
 - ▶ 3 charge-dependent, 1 charge asymmetric operator
 - ▶ 40 parameters fitted to pp and np scattering data up to 350 MeV
- ▶ Extensively used in VMC and GFMC calculations of light nuclei
- ▶ Not used much in NCSM calculations
 - ▶ slow convergence ('hard' interaction)
 - ▶ would need significant renormalization
- ▶ Additional 3-nucleon forces
 - ▶ Urbana IX – based on 2π exchange diagrams (in particular Fujita–Miyazawa term) plus short-range phenomenological terms
 - ▶ Illinois 7 – most important additional term: 3π ring diagram with one or two Δ 's



Energies with AV18 + IL7

Carlson, Gandolfi, Pederiva, Pieper, Schiavilla, Schmidt and Wiringa, RMP 87, 1067 (2015)



Carlson, Gandolfi, Pederiva, Pieper, Schiavilla, Schmidt and Wiringa, RMP 87, 1067 (2015)



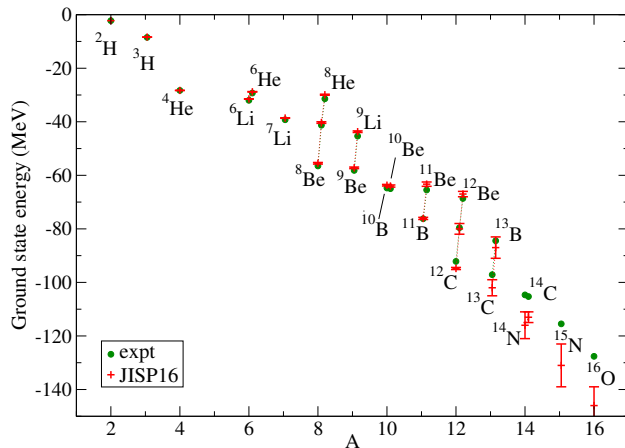
JISP16

Shirokov, Vary, Mazur and Weber, PLB 644, 33 (2007)

J-matrix Inverse Scattering Potential

- ▶ Nonlocal finite rank separable NN potential in H.O. representation, constructed to reproduce np scattering data
 - ▶ $2n + l \leq 8$ for even partial waves, limited to $J \leq 4$
 - ▶ $2n + l \leq 9$ for odd partial waves, limited to $J \leq 4$
 - ▶ $\hbar\omega = 40$ MeV
 - ▶ χ^2/datum of 1.05 for the 1999 np data base (3058 data)
- ▶ Use Phase-Equivalent Transformations (PET) to tune off-shell interaction to
 - ▶ deuteron quadrupole moment
 - ▶ binding energy of ^3H and ^4He
 - ▶ low-lying states of ^6Li (JISP6, precursor to JISP16)
 - ▶ binding energy of ^{16}O
- ▶ Unfortunately, convergence insufficiently understood and basis space was limited when tuning to ^{16}O

Ground state energies of p -shell nuclei with JISP16














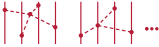



Maris and Vary,
IJMPE22, 1330016 (2013)

- ▶ ^{10}B – most likely JISP16 produces correct 3^+ ground state, but extrapolation of 1^+ states not reliable due to mixing of two 1^+ states
- ▶ ^{11}Be – expt. observed parity inversion within error estimates of extrapolation
- ▶ ^{12}B and ^{12}N – unclear whether gs is 1^+ or 2^+ (expt. at $E_x \approx 1$ MeV) with JISP16

Nuclear interactions from Chiral Effective Field Theory

- ▶ Controlled power series expansion in Q/Λ
- ▶ Hierarchy for many-body forces $V_{NN} \gg V_{NNN} \gg V_{NNNN}$

Chiral expansion of nuclear forces

	Two-nucleon force	Three-nucleon force	Four-nucleon force
LO (Q^0)			
NLO (Q^2)			
N ² LO (Q^3)			
N ³ LO (Q^4)			
N ⁴ LO (Q^5)			

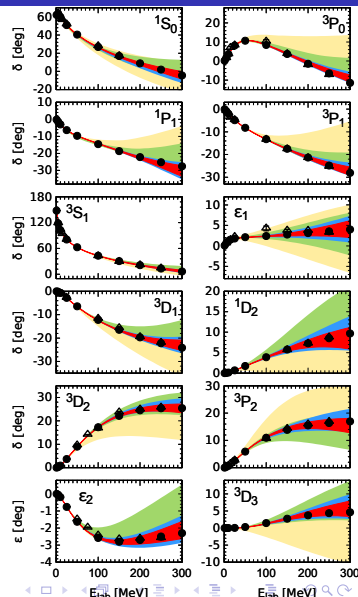
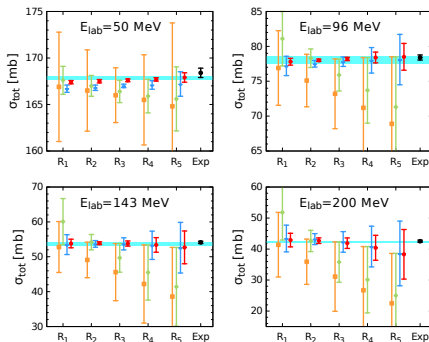
NN potential from χ EFT up to N⁴LO

Epelbaum, Krebs, Meißner, PRL 115 (2015); EPJ A51 (2015)

- Local regulator long-range terms

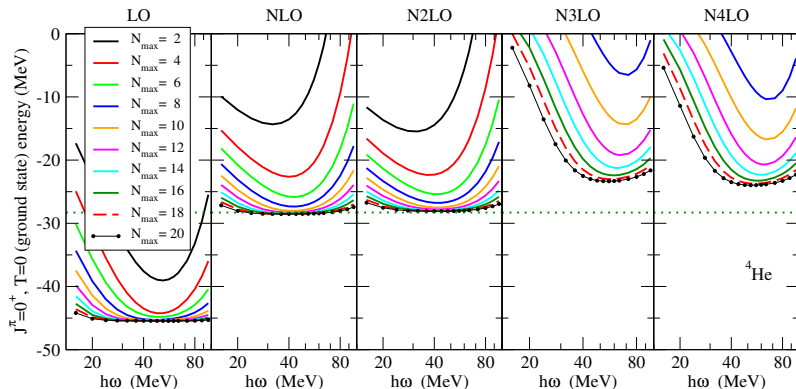
$$V(r) \rightarrow V(r) \left[1 - \exp(-r^2/R^2)\right]^6$$

- Regulators $R_1 = 0.8$ to $R_5 = 1.2$ fm
- Reduced finite-cutoff artefacts



Ground state energy of ^4He

LENPIC collaboration, PRC 93, 044002 (2016)

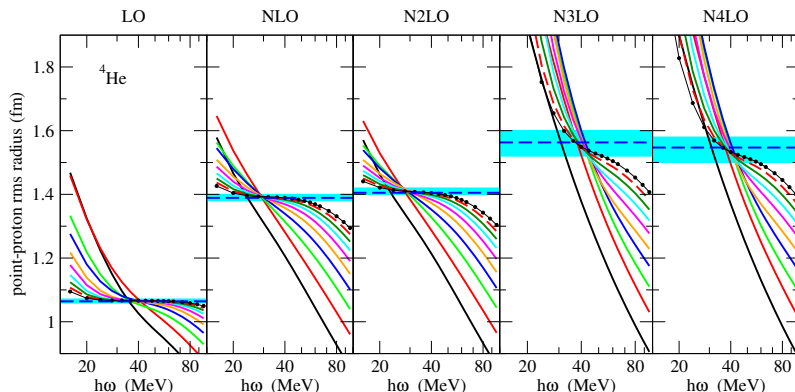


Chiral NN interaction with regulator $R = 1.0$ fm ($\Lambda_b = 600$ MeV)

- ▶ Many-body calc'ns converge rapidly at LO, NLO, and $N^2\text{LO}$
- ▶ Convergence significantly slower at $N^3\text{LO}$ and $N^4\text{LO}$
- ▶ No 3NFs included (yet) – should be present at $N^2\text{LO}$ and up

Point-proton rms radius of ^4He

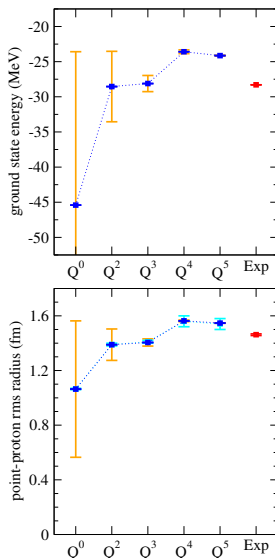
LENPIC collaboration, PRC 93, 044002 (2016)



- ▶ Convergence of many-body calculation for RMS radius slower than convergence for (ground state) energy
 - ▶ Long-range operator
 - ▶ H.O. basis function fall off like gaussians, instead of exponential
 - ▶ Nevertheless, agree with Faddeev–Yakubovsky calc'ns

Chiral EFT uncertainty estimates

LENPIC collaboration, PRC 93, 044002 (2016)



- ▶ Chiral expansion in $m_\pi/\Lambda_b \approx 0.23$

$$X^{(i)} = X^{(0)} + \Delta X^{(2)} + \dots + \Delta X^{(i)}$$

- ▶ Expected chiral corrections at order i

$$\Delta X^{(i)} \approx \mathcal{O}(Q^i X^{(0)})$$

- ▶ Chiral uncertainty estimates

$$\delta X^{(0)} \approx Q^2 |X^{(0)}|$$

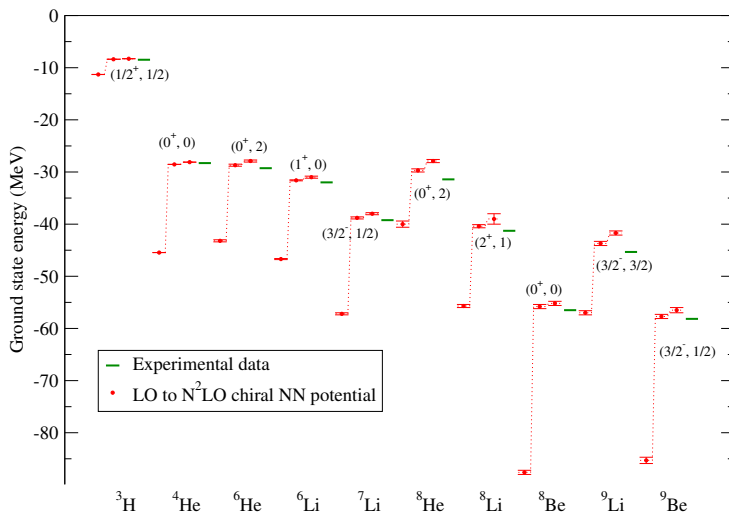
$$\delta X^{(2)} \approx \max(Q \delta X^{(0)}, Q |\Delta X^{(2)}|)$$

- ▶ 3NF at N²LO ($i = 3$) and up not yet included

$$\delta X^{(i \geq 3)} \approx \max(Q \delta X^{(i-1)}, Q^{i-2} |\Delta X^{(3)}|)$$

- ▶ Additional condition on LO uncertainty estimate $\delta X^{(0)} \geq \max(|X^{(i)} - X^{(j)}|)$

Ground state energies with χ EFT up to $A = 9$

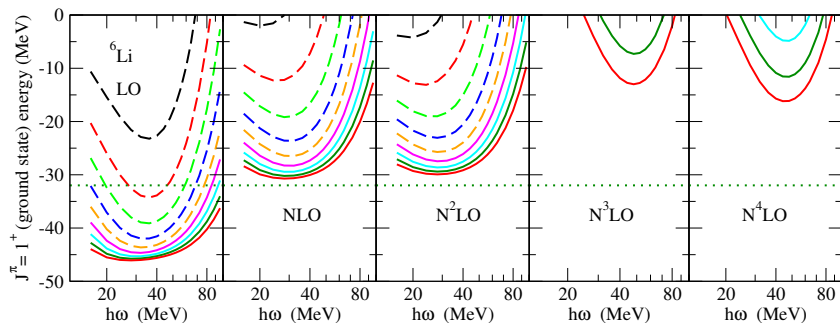


LENPIC collaboration

in preparation

Results for ${}^6\text{Li}$ with χEFT NN potential

LENPIC collaboration, PRC 93, 044002 (2016)



- ▶ Up to $N^2\text{LO}$ good numerical convergence
- ▶ At $N^3\text{LO}$ (and $N^4\text{LO}$) convergence significantly slower
- ▶ Need to use renormalization techniques to accelerate convergence

Renormalization

Challenge: achieve numerical convergence for No-Core CI calculations using a finite amount of CPU time on current HPC systems

- ▶ Use unitary transformations to renormalize interaction
 - ▶ can improve quality of results in small basis spaces
 - ▶ need to renormalize other operators as well
- ▶ Commonly used in NCSM calculations
 - ▶ Similarity Renormalization Group
 - ▶ Okubo–Lee–Suzuki
 - ▶ $V_{\text{low } k}$, V_{UCOM} , ...
- ▶ In principle, unitary transformations change the wavefunction, but should not change physical observables
- ▶ In practice, induced many-body effects are neglected ...
 - ▶ need to study effect of induced many-body forces

Similarity Renormalization Group

Glazek, Wilson, PRD 48, 5863 (1993); PRD 49, 4214 (1994); Wegner, Ann. Phys. 3, 77 (1994)

- ▶ Consider series of unitary transformations

$$H_\alpha = U_\alpha H U_\alpha^\dagger$$

- ▶ Define anti-hermitian operators η_α

$$\eta_\alpha = \frac{dU_\alpha}{d\alpha} U_\alpha^\dagger = -\eta_\alpha^\dagger$$

such that H_α evolves according to

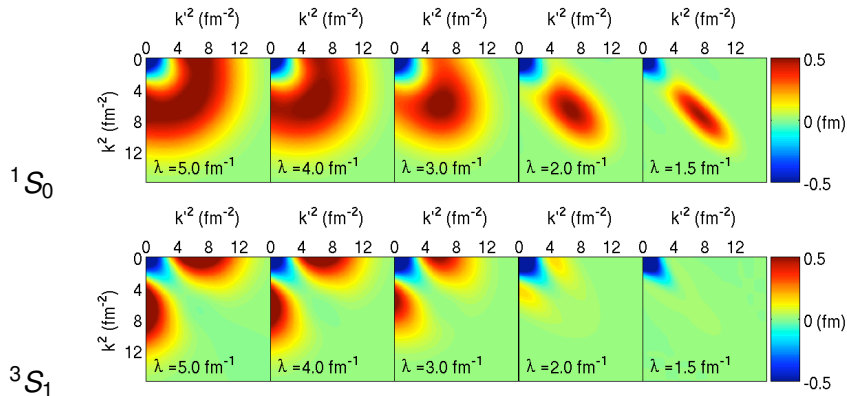
$$\frac{dH_\alpha}{d\alpha} = [\eta_\alpha, H_\alpha]$$

- ▶ Common choice for 'generator' η_α

$$\eta_\alpha = (2\mu)^2 [T_{\text{rel}}, H_\alpha]$$

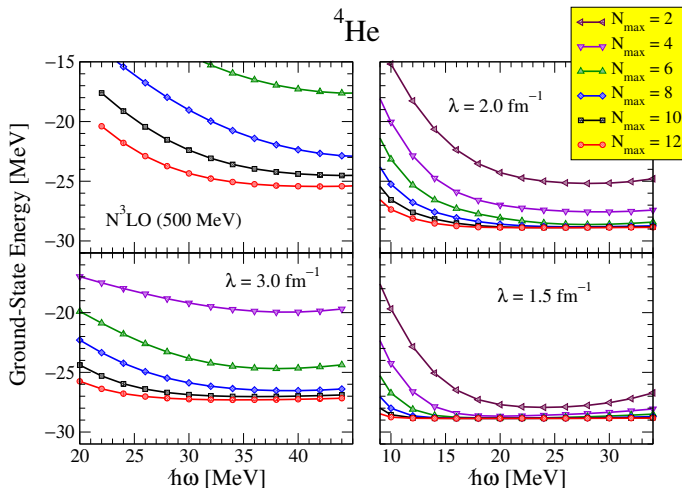
SRG evolution of Idaho χ EFT at N^3 LO(500 MeV)

Bogner, Furnstahl, Perry, PRC 75, 061001 (2007)



- Drives interaction to diagonal
- Decouples low and high momenta
- Note: SRG parameter $\lambda = \alpha^{-\frac{1}{4}}$

^4He with SRG evolved χEFT

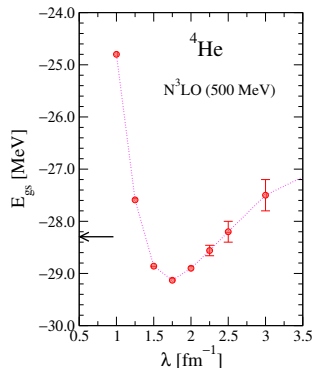
Bogner *et al*, NPA 801, 21 (2008)

- ▶ SRG renormalization improves convergence significantly
- ▶ SRG parameter dependence indicates omitted induced 3NF

^4He with SRG evolved χEFT including (induced) 3NF

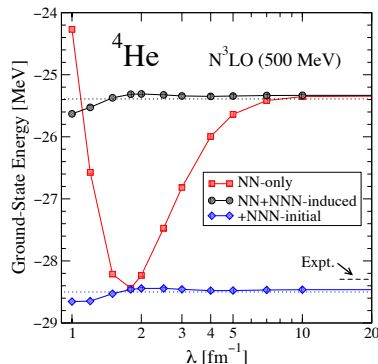
Bogner *et al*, NPA 801, 21 (2008)

(without Coulomb interaction)



Jurgenson, Navratil, Furnstahl, PRL 103, 082501 (2009)

(with Coulomb interaction)

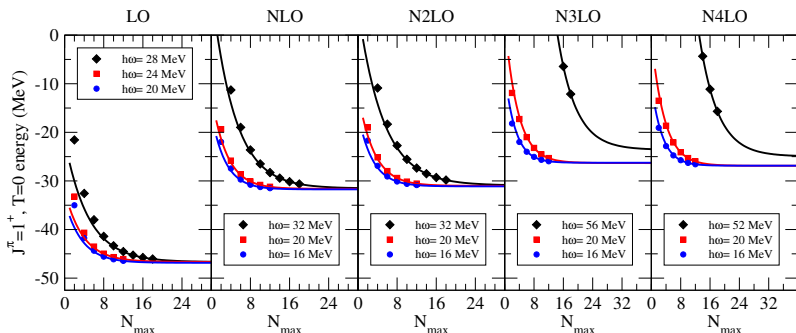


- ▶ Strong SRG parameter $\alpha = \lambda^{-4}$ dependence without induced 3NF
- ▶ Almost no SRG parameter $\alpha = \lambda^{-4}$ dependence with induced 3NF
- ▶ Explicit 3NF needed for agreement with experiment

${}^6\text{Li}$ with SRG evolved χEFT including induced 3NF

- Empirical extrapolation method (ground state) energies

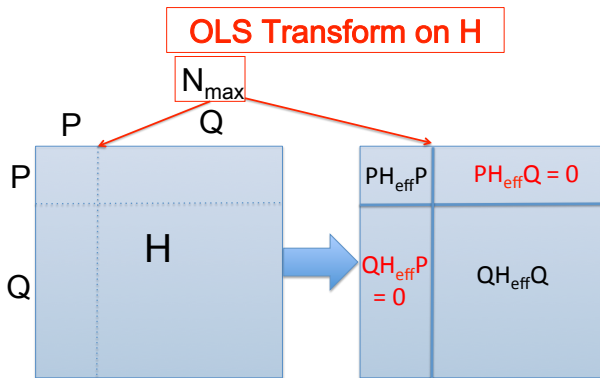
$$E(N_{\text{max}}) \approx E_{\infty} + a \exp(-bN_{\text{max}})$$



- Extrapolations at different SRG α and without SRG are consistent with each other to within estimated extrapolation uncertainty

Okubo–Lee–Suzuki renormalization

- ▶ Divide large (but finite) space into a 'P'-space and a 'Q'-space
- ▶ Construct unitary transformation that decouples 'P'- and 'Q'-space
- ▶ Effective Hamiltonian and other operators in 'P'-space
- ▶ Physical observables remain the same



Okubo–Lee–Suzuki renormalization

- ▶ Diagonalize original Hamiltonian $H = T + V$,
to obtain D eigenvalues E_i and matrix of eigenvectors U

$$U H U^\dagger = \text{diag}(E_i)$$

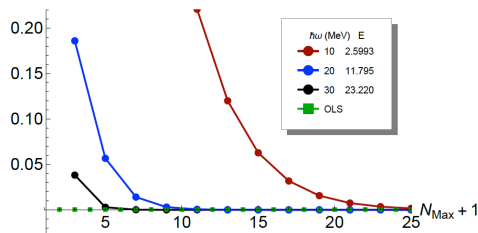
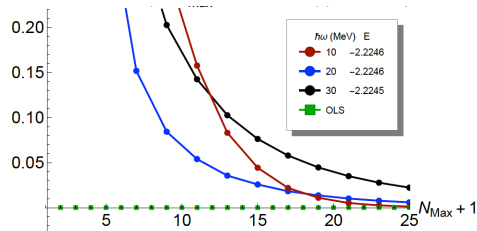
- ▶ Project U to the P-space: $U_P = P^T U P$
- ▶ Calculate the norm of $U_P^\dagger U_P$, that is,
the d eigenvalues N_i of $U_P^\dagger U_P$ and corresponding eigenvectors W
- ▶ Iff $\|U_P^\dagger U_P\| > 0$, that is,
if all eigenvalues are positive, we can proceed

$$\begin{aligned} H_{\text{OLS}} &= \left[W \text{diag}(1/\sqrt{N_i}) W^\dagger P^T \right] H \left[P W \text{diag}(1/\sqrt{N_i}) W^\dagger \right] \\ &= U_{\text{OLS}} H U_{\text{OLS}}^\dagger = P^T T P + V_{\text{OLS}} \end{aligned}$$

OLS renormalization for deuteron

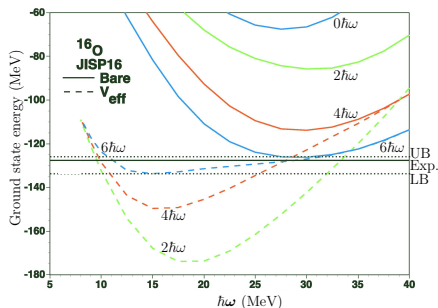
Toy model: coupled sd channel in relative H.O. basis with JISP16

- Fractional difference between converged ground state energy and g.s. energy in smaller bases with $H = T + V_{\text{JISP16}}$
- Same, but with $H = T + V_{\text{JISP16}} + V_{\hbar\omega}^{\text{H.O.}}$

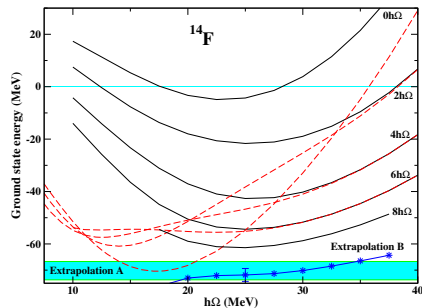


Energies for JISP16 with OLS renormalization

Shirokov, Vary, Mazur and Weber, PLB 644, 33 (2007)



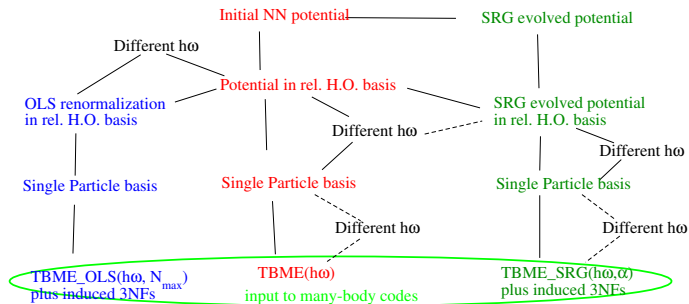
PM, Shirokov and Vary, PRC81, 021301(R) (2010)



OLS renormalized V_{OLS} without induced many-body forces

- ▶ About 10 years ago, was believed to be a lower bound
- ▶ Extrapolation to complete basis with bare potential is monotonic
- ▶ Results V_{OLS} approach bare results as N_{max} increases

Appendix: From (effective) interactions to H.O. matrix elements in single-particle coordinates



Essential H.O. transformations

- ▶ Relative \rightarrow single-particle (Talmi-Moshinsky brackets)
- ▶ Transformation $\hbar\omega \rightarrow \hbar\tilde{\omega}$ (in relative basis)

Renormalization: SRG evolution and/or OLS in relative basis

For further reading

- ▶ J. Carlson, S. Gandolfi, F. Pederiva, S.C. Pieper, R. Schiavilla, K.E. Schmidt and R.B. Wiringa, *Quantum Monte Carlo methods for nuclear physics*, Rev. Mod. Phys. **87**, 1067 (2015).
- ▶ E. Epelbaum, H.-W. Hammer and U.-G. Meißner, *Modern Theory of Nuclear Forces*, Rev. Mod. Phys. **81** 1773 (2009).
- ▶ R. Machleidt and D.R. Entem, *Chiral effective field theory and nuclear forces*, Phys. Rept. 503, 1 (2011).
- ▶ S.K. Bogner, R.J. Furnstahl and A. Schwenk, *From low-momentum interactions to nuclear structure*, Prog. Part. Nucl. Phys. **65**, 94 (2010).
- ▶ R. Roth, A. Calci, J. Langhammer and S. Binder, *Evolved Chiral NN+3N Hamiltonians for Ab Initio Nuclear Structure Calculations*, Phys. Rev. C90, 024325 (2014).