## Emergence of collective motion

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### Outline

Binding and excitation energies with JISP16

Electromagnetic observables with JISP16

Emergence of rotational bands

Detailed example: 9Be

Intrinsic rotational band paramereters

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#### Ground state energies of *p*-shell nuclei with JISP16

Maris and Vary, IJMPE22, 1330016 (2013)



# Extrapolating to complete basis

Challenge: achieve numerical convergence for No-Core CI calculations using a finite amount of CPU time on current HPC systems

- Perform a series of calculations with increasing N<sub>max</sub> truncation
- $\blacktriangleright$  Extrapolate to infinite model space  $\longrightarrow$  exact results
  - Empirical: binding energy exponential in N<sub>max</sub>

$$E_{\text{binding}}^{N} = E_{\text{binding}}^{\infty} + a_1 \exp(-a_2 N_{\text{max}})$$

- use 3 or 4 consecutive  $N_{\text{max}}$  values to determine  $E_{\text{binding}}^{\infty}$
- use ħω and N<sub>max</sub> dependence to estimate numerical error bars
   Maris, Shirokov, Vary, PRC79, 014308 (2009)
- ► Recent studies of IR and UV behavior based on S.P. asymptotics: exponentials in √ħω/N and √ħωN Coon et al, PRC86, 054002 (2012);

Furnstahl, Hagen, Papenbrock, PRC86, 031301(R) (2012); More, Ekstrom, Furnstahl, Hagen, Papenbrock, PRC87, 044326 (2013); Wendt, Forssén, Papenbrock and Sääf, PRC91, 061301 (2015); More, PhD thesis OSU, arXiv:1608.01385 [nucl-th];

### Extrapolating to complete basis - in practice

Perform a series of calculations with increasing N<sub>max</sub> truncation

 H.O. basis up to N<sub>max</sub> = 16 and exponential extrapilation E<sub>b</sub> = -31.49(3) MeV

Cockrell, Maris, Vary, PRC86, 034325 (2012)



• Hyperspherical harmonics up to  $K_{max} = 14$ :  $E_b = -31.46(5)$  MeV

Vaintraub, Barnea, Gazit, PRC79, 065501 (2009)

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Example: <sup>9</sup>Be

Band paramereters

# Spectrum of <sup>6</sup>Li



- Excitation energies narrow states reasonably well converged
- No need for extrapolations

#### Energies of excited states of A = 6 to A = 9 nuclei

Maris and Vary, IJMPE22, 1330016 (2013)



#### Local one-body density

One-body density in single-particle coordinates

$$\rho(\vec{r}) = \int |\Psi(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_A)|^2 d^3 r_2 \dots d^3 r_A$$

• Lab-frame density  $\rho^{\omega}(\vec{r})$  includes Center-of-Mass motion

$$ho^{\omega}(\vec{r}) = \int 
ho_{
m rel}(\vec{r}-\vec{R}) 
ho^{\omega}_{
m CM}(\vec{R}) d^{3}\vec{R}.$$

depends on basis ħω, even in complete (infinitely large) basis
 Deconvolution of relative density and Center-of-Mass motion

$$\rho_{\rm rel}(\vec{r}) = F^{-1} \left[ \frac{F[\rho^{\omega}(\vec{r})]}{F[\rho^{\omega}_{\rm CM}(\vec{r})]} \right]$$

Multipole expansion (used to facilitate deconvolution)

$$\rho(\vec{r}) = \sum_{K=0}^{2J} \frac{\langle JMK0 | JM \rangle}{\sqrt{2J+1}} Y_K^0(\theta) \rho^{(K)}(r)$$

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# Density of <sup>6</sup>Li

Cockrell, Maris, Vary, PRC86, 034325 (2012)



- Slow convergence of asymptotic tail of wavefunction in particular for larger ħω values
- Hence, slow convergence of RMS radii, quadrupole moments, etc.

# Radius of <sup>7</sup>Be



• Calculation one-body observables  $\langle i | \mathcal{O} | j \rangle \sim \int \mathcal{O}(r) r^2 \rho_{ij}(r) dr$ 

- RMS radius:  $\mathcal{O}(r) = r^2$
- Slow convergence of RMS radius due to slow build up of asymptotic tail
- Ground state RMS radius in agreement with data

### Multipole operators

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Electric quadrupole (E2) operator

$$\mathbf{Q}_{2} = \sum_{i=1}^{N} e_{i} r_{i}^{2} Y_{2u}(\mathbf{r}_{i}) = e_{p} \mathbf{Q}_{p} + e_{n} \mathbf{Q}_{n} \qquad e_{p} = e \quad e_{n} = 0$$
$$\mathbf{Q}_{p} \sim \sum_{i=1}^{Z} r_{p,i}^{2} Y_{2u}(\mathbf{r}_{p,i}) \qquad \mathbf{Q}_{n} \sim \sum_{i=1}^{N} r_{n,i}^{2} Y_{2u}(\mathbf{r}_{n,i}) \qquad Proton \ \mathcal{E} \ neutron \ tensors$$

Magnetic dipole (M1) operator

$$\mathbf{M}_{1} = \sqrt{\frac{3}{4\pi}} \mu_{N} \sum_{i=1}^{A} (g_{\ell}^{(i)} \boldsymbol{\ell}_{i} + g_{s}^{(i)} s_{i}) \qquad \begin{array}{c} g_{\ell,p} = 1 & g_{\ell,n} = 0 \\ g_{s,p} \approx 5.585 & g_{s,n} \approx -3.826 \end{array}$$
$$= g_{\ell,p} \mathbf{D}_{\ell,p} + g_{\ell,n} \mathbf{D}_{\ell,n} + g_{s,p} \mathbf{D}_{s,p} + g_{s,n} \mathbf{D}_{s,n},$$
$$\mathbf{D}_{\ell,p} \sim \sum_{i=1}^{Z} \boldsymbol{\ell}_{p,i} \quad \mathbf{D}_{\ell,n} \sim \sum_{i=1}^{N} \boldsymbol{\ell}_{n,i} \quad \mathbf{D}_{s,p} \sim \sum_{i=1}^{Z} s_{p,i} \quad \mathbf{D}_{s,n} \sim \sum_{i=1}^{N} s_{n,i} \quad Dipole \ terms$$

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# Dipole terms of <sup>7</sup>Li

Maris, Vary, IJMPE22, 1330016 (2013)

$$J = \frac{1}{J+1} \left( \langle \vec{J} \cdot \vec{L}_{p} \rangle + \langle \vec{J} \cdot \vec{L}_{n} \rangle + \langle \vec{J} \cdot \vec{S}_{p} \rangle + \langle \vec{J} \cdot \vec{S}_{n} \rangle \right)$$



• Converged with  $N_{\text{max}}$ , persistent weak  $\hbar\omega$  dependence  $\frac{5}{2}^{-}$  states

• Two  $\frac{5}{2}^{-}$  states have very different structure

#### Magnetic moments of *p*-shell nuclei with JISP16



Magnetic moments reasonably well converged

Deviations from experiment: missing meson exchange currents

#### Quadrupole moment and E2 transition strengths 7Li

Cockrell, Maris, Vary, PRC86 034325 (2012)



- E2 observables not converged, due to gaussian fall-off of HO wavefunction
- Nevertheless, qualitative agreement of Q and B(E2) with data

Example: <sup>9</sup>Be

Band paramereters

### Spectrum of A = 6 to 9 nuclei with JISP16

Maris and Vary, IJMPE22, 1330016 (2013)



Rotational band?

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Example: <sup>9</sup>Be Band parame

#### Rotational model predictions

Intrinsic state  $|\phi_K\rangle$  & rotation in Euler angles  $\vartheta$  (J = K, K + 1, ...) $|\psi_{JKM}\rangle \propto \int d\vartheta \Big[ \mathscr{D}^{J}_{MK}(\vartheta) |\phi_{K};\vartheta\rangle + (-)^{J+K} \mathscr{D}^{J}_{M-K}(\vartheta) |\phi_{\bar{K}};\vartheta\rangle \Big]$ Rotational energy  $E(J) = E_0 + A[J(J+1) + a(-)^{J+1/2}(J+\frac{1}{2})] \qquad A \equiv \frac{\hbar^2}{2\pi}$ Coriolis (K = 1/2) Rotational relations on electromagnetic transitions (E2, M1, ...)Ē - 01  $E_{c}$ 5/2 7/2 1/23/29/2J

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### Rotational band: Quadrupole matrix elements

$$\begin{aligned} \langle \psi_{J_f \mathcal{K}} || \mathcal{E}_2 || \psi_{J_f \mathcal{K}} \rangle &= \frac{(2J_i + 1)^{1/2}}{1 + \delta_{\mathcal{K}0}} \Big( (J_i, \mathcal{K}, 2, 0 | J_f, \mathcal{K}) \langle \phi_{\mathcal{K}} || \mathcal{E}_{2,0} || \phi_{\mathcal{K}} \rangle \\ &+ (-)^{J_i + \mathcal{K}} (J_i, -\mathcal{K}, 2, 2\mathcal{K} | J_f, \mathcal{K}) \langle \phi_{\mathcal{K}} || \mathcal{E}_{2,2\mathcal{K}} || \phi_{\bar{\mathcal{K}}} \rangle \Big) \end{aligned}$$



Consider both proton and neutron quadrupole tensors

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#### Rotational band: Dipole matrix elements

Magnetic moments

$$\mu(J) = a_0 J + a_1 \frac{K}{J+1} + a_2 \delta_{K,\frac{1}{2}} \frac{(-1)^{J-\frac{1}{2}}}{2\sqrt{2}} \frac{2J+1}{J+1}$$

Magnetic transition matrix elements

$$\langle \psi_{J-1,K} || M_1 || \psi_{J,K} \rangle = -\sqrt{\frac{3}{4\pi}} \sqrt{\frac{J^2 - K^2}{J}} \left( a_1 + a_2 \delta_{K,\frac{1}{2}} \frac{(-1)^{J-\frac{1}{2}}}{\sqrt{2}} \right)$$

Define dipole terms D<sub>l,p</sub>, D<sub>l,n</sub>, D<sub>s,p</sub>, and D<sub>s,n</sub> for both the magnetic moments and for the M<sub>1</sub> transitions

$$M_1 = g_{l,p} D_{l,p} + g_{l,n} D_{l,n} + g_{s,p} D_{s,p} + g_{s,n} D_{s,n}$$

with  $g_{l,p} = 1$ ,  $g_{l,n} = 0$ ,  $g_{s,p} = 5.586$ , and  $g_{s,n} = -3.826$ 

Example: <sup>9</sup>Be Band parar

### <sup>8</sup>Be ground state rotational band

Shell model: Valence space angular momentum  $J \le 4$ Cluster model: Molecular rotation of  $\alpha + \alpha$  dimer



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Example: <sup>9</sup>Be Band param

## <sup>8</sup>Be ground state rotational band

Shell model: Valence space angular momentum  $J \le 4$ Cluster model: Molecular rotation of  $\alpha + \alpha$  dimer



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Example: <sup>9</sup>Be I

Band paramereters

## <sup>8</sup>Be ground state rotational band

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### Candidate Rotational bands in Be isotopes



### E2 moments



### E2 moments and transitions



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Example: <sup>9</sup>Be Band par

### E2 and M1 moments and transitions



## Convergence with basis space

Absolute binding energy? NO!



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Example: <sup>9</sup>Be

### Convergence with basis space

Absolute binding energy? NO! Excitation within band? ~YES



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#### Convergence with basis space

Absolute *E*2? **NO**! Ratio of *E*2? ~**YES** Absolute *M*1? ~**YES** 



Example: <sup>9</sup>Be Ba

Band paramereters

## Inter- and Intra-band E2 transtions

Caprio, Maris, Vary, Smith, IJMPE 24, 1541002 (2015)



E2 transition strenght between natural (negative) parity states in <sup>9</sup>Be

Transitions within g.s. (K = 3/2) and (K = 1/2) bands significantly enhanced over typical E2 transition strenght



# One-body density of 9Be ground state $(\frac{3}{2}, \frac{1}{2})$





#### and their difference



nagnetic observables

Rotational banc

Example: <sup>9</sup>Be

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### Extraction of band parameters

Caprio, Maris, Vary, Smith, IJMPE 24, 1541002 (2015)



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#### Convergence of intrinsic quadrupole moments

Caprio, Maris, Vary, Smith, IJMPE 24, 1541002 (2015)

Quadrupole moments? NO! Ratios of p & n moments? ~ YES Enhanced relative to single-particle strength (Weisskopf)



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## Convergence of observables

- Binding energies
  - need extrapolations
  - alternatives for H.O. basis and/or different truncation scheme?
- Excitation energies
  - okay for narrow states (of similar structure as bound state)
  - H.O. basis not suited for broad resonances
  - alternatives for H.O. basis, specifically for resonances?
- Magnetic moments and transitions
  - converge generally rapidly
  - need to add meson-exchange currents
- Quadrupole moments and transitions
  - converge slowly
  - however, ratio's of E2 observables reasonably well converged
  - extrapolations? Odell, Papenbrock and Platter, PRC 93, 044331 (2016)
  - alternatives for H.O. basis and/or different truncation scheme?
- Emergence of rotational structure and clustering