

A self-consistent equation of motion multiphonon method for even and odd mass nuclei

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Semiclassical

Microscopic

$$\{\alpha_\lambda, \pi_\lambda\} \rightarrow \{O_\lambda, O_\lambda^\dagger\}$$

EofM

$$[H, O_\lambda^\dagger] = \hbar\omega_\lambda O_\lambda^\dagger$$

TDA mapping

$$O_\lambda^\dagger = \sum_{ph} c_{ph}(\lambda) a_p^\dagger a_h$$

EoM

$$[H, O_\lambda^\dagger] |> = \hbar\omega_\lambda O_\lambda^\dagger |>$$

Collective modes

RPA mapping

$$O_\lambda^\dagger = \sum_{ph} [X_{ph}(\lambda) a_p^\dagger a_h - Y_{ph}(\lambda) a_h^\dagger a_p]$$

EoM

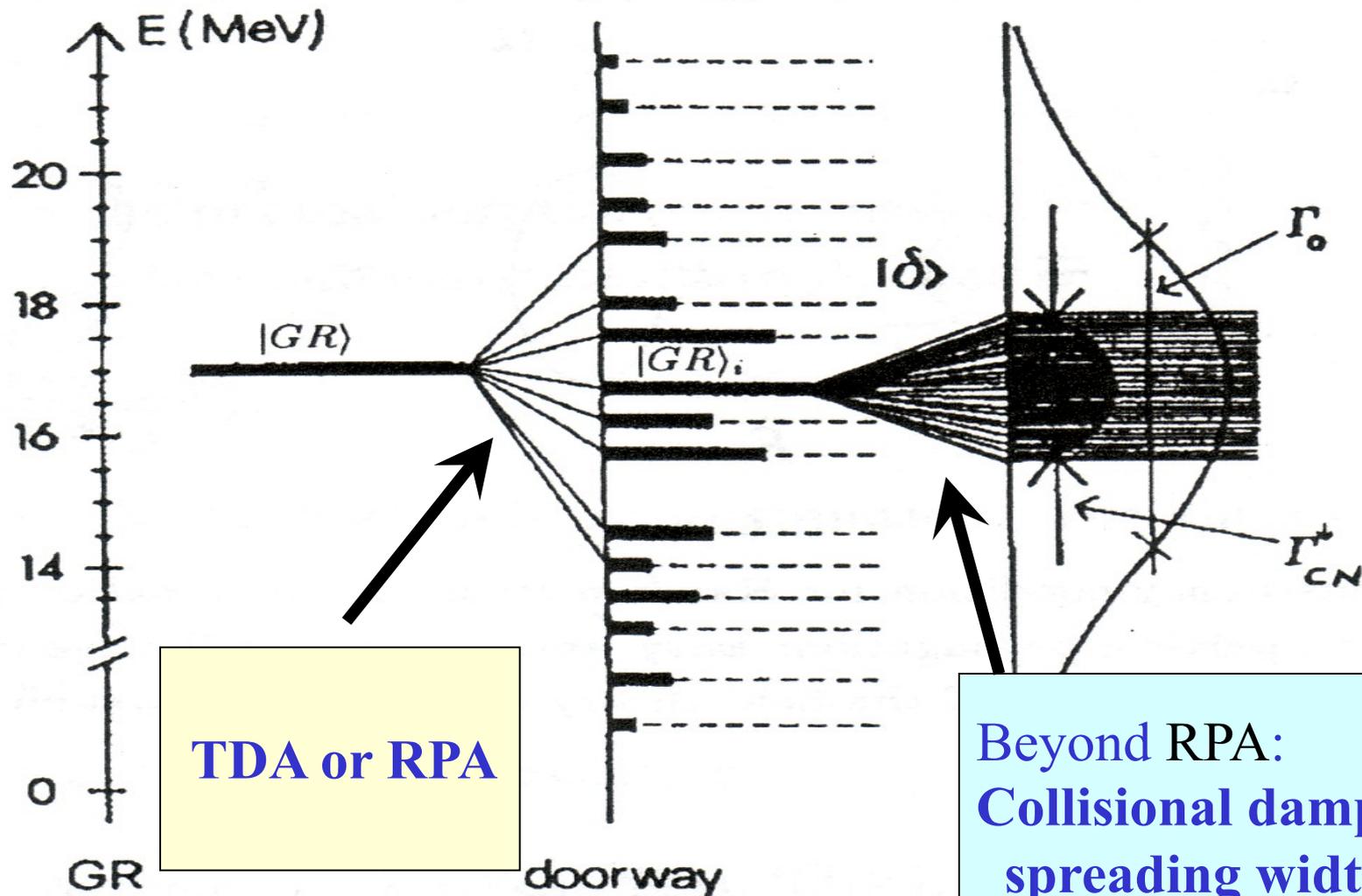
$$[H, O_\lambda^\dagger] |0> = \hbar\omega_\lambda O_\lambda^\dagger |0>$$

$|0> \equiv$ correlated g.s

Underlying approximation:

$$|0> \rightarrow |HF> \text{ QBA}$$

Collective modes: anharmonic features



From P.F. Bortignon, A. Bracco, R.A. Broglia, Giant Resonances (hap, 1998)

Beyond mean field: *Adopted Methods*

Non relativistic

- qp-Phonon

P. F. Bortignon et al. (Milano group)

- 2nd RPA

R. Roth et al.

D. Gambacurta et al.

Phenomenological

- QPM (*Soloviev School (Dubna)*)

Relativistic: RTBA

E. Litvinova, P. Ring, D. Vretenar.....

Our proposal: EMPM

*D. Bianco, F. Knapp, N. Lo Iudice, F. Andreozzi, A. Porrino, Phys. Rev. C **85**, 014313 (2012).*

Equation of motion phonon method (EMPM)

Eigenvalue problem

$$\mathbf{H} |\Psi_v\rangle = E_v |\Psi_v\rangle$$

$$|\Psi_v\rangle \in \mathcal{H} = \sum_n \oplus \mathcal{H}_n \quad \mathcal{H}_n \in |n; \beta\rangle \equiv \text{n-phonon basis states}$$

An obvious, **but unmanageable**, multiphonon basis

$$|\lambda_1, \dots, \lambda_i, \dots, \lambda_n\rangle = \mathbf{O}_{\lambda_1}^\dagger \dots \mathbf{O}_{\lambda_i}^\dagger \dots \mathbf{O}_{\lambda_n}^\dagger |0\rangle$$

$$\mathbf{O}_{\lambda}^\dagger = \sum_{ph} c_{ph}^\lambda \mathbf{a}_p^\dagger \mathbf{a}_h$$

A viable route

$$|\alpha_n\rangle = \sum_{\lambda\alpha} C_{\lambda\alpha}^\beta \mathbf{O}_{\lambda}^\dagger |\alpha_{n-1}\rangle$$

Construction of $|\mathbf{n}; \beta\rangle$: EoM

Assuming $|\alpha_{n-1}\rangle$ known, we solve the **Eq. of Motion**

$$\langle \alpha_n | [\mathbf{H}, \mathbf{O}^\dagger_\lambda] | \alpha_{n-1} \rangle = (\mathbf{E}_\beta^{(n)} - \mathbf{E}_\alpha^{(n-1)}) \langle \alpha_n | \mathbf{O}^\dagger_\lambda | \alpha_{n-1} \rangle$$


$$\mathcal{A} \mathcal{X} = \mathcal{E} \mathcal{X}$$


$$\mathcal{X}$$

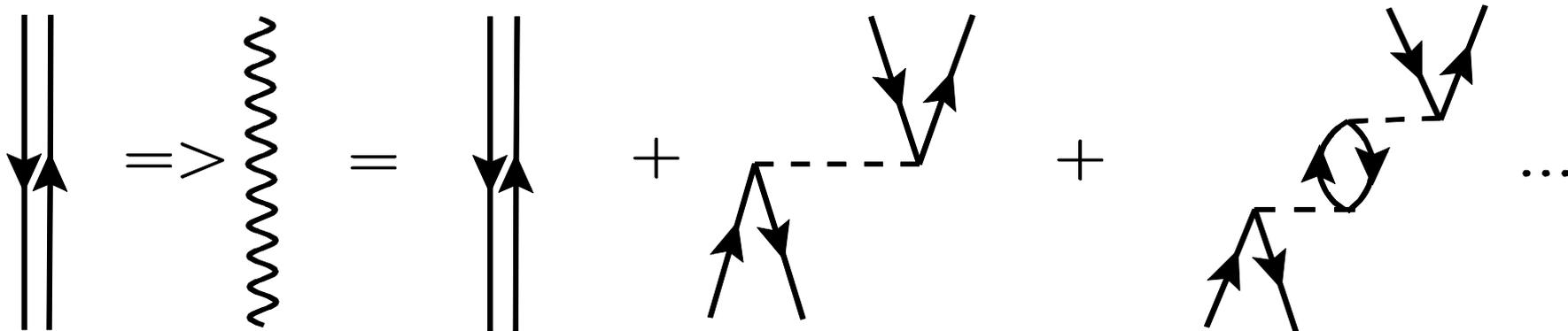
Where

$$\mathcal{A}_{(\mu\alpha)(\nu\gamma)} = (\mathbf{E}_\mu + \mathbf{E}_\alpha) \delta_{\alpha\gamma} \delta_{\mu\nu} + \mathcal{V}_{(\mu\alpha)(\nu\gamma)}$$

TDA matrix

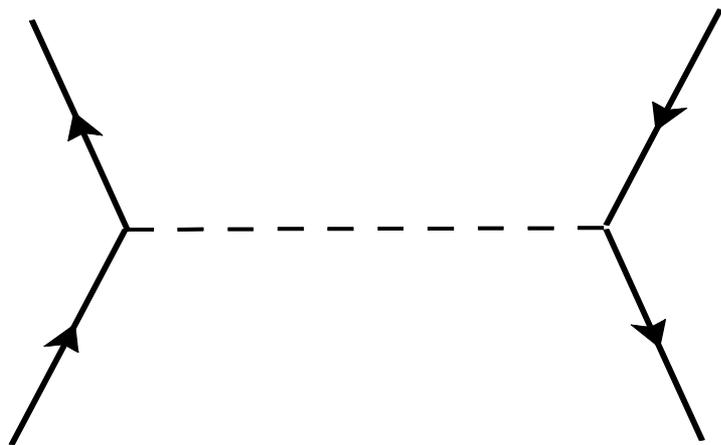
$$\mathbf{A}_{(ph)(p'h')} = (\epsilon_p - \epsilon_h) \delta_{pp'} \delta_{hh'} + \mathbf{V}_{ph'hp'}$$

From **p-h** to **TDA** to **EMPM**

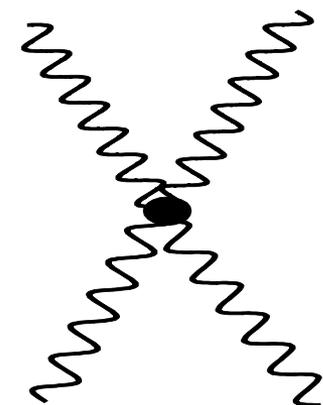


$$\varepsilon_p - \varepsilon_h \Rightarrow E_\mu$$

$$|\mathbf{ph}\rangle \Rightarrow |\boldsymbol{\mu}\rangle = \sum_{\mathbf{ph}} \mathbf{c}_{\mathbf{ph}}^\mu \mathbf{a}_p^\dagger \mathbf{a}_h | \rangle$$



$$\mathbf{V}_{\mathbf{ph}'\mathbf{hp}'}$$



$$\mathbf{V}_{(\mu\nu)(\mu'\nu')}$$

Construction of $|\mathbf{n}; \beta\rangle$: EoM

Problem

$$\mathcal{A} \mathcal{X} = \mathcal{E} \mathcal{X}$$

is **not** a true Eigenvalue Eq.!

$\{\mathbf{O}^\dagger_\lambda |\mathbf{n}-1, \alpha\rangle\}$ form a **non-orthogonal redundant** basis



$\mathcal{X} = \langle \mathbf{n}, \beta | \mathbf{O}^\dagger_\lambda |\mathbf{n}-1, \alpha\rangle$ is **not** a **true** expansion coefficient

Recipe for solving the problem

1° step

$$|\alpha_n\rangle = \sum_{\lambda\alpha} C_{\lambda\alpha}^\beta \mathbf{O}_{\lambda}^\dagger |\alpha_{n-1}\rangle \}$$

$$|\lambda\alpha\rangle = \mathbf{O}_{\lambda}^\dagger |\alpha_n\rangle$$

$$(X = \langle \alpha_n | \lambda \alpha_{n-1} \rangle)$$

$$\mathcal{D} \equiv \{ \langle \lambda' \alpha_n' | \lambda \alpha_n \rangle \}$$

$$\mathcal{A} X = \mathcal{E} X$$

$$X = \mathcal{D} C$$

$$[\mathcal{H} - \mathcal{E}\mathcal{D}] C = \mathbf{0}$$

where

$$\mathcal{H} = \mathcal{A}\mathcal{D}$$

But \mathcal{D} is **singular** !

2° conclusive step: Choleski

$$[\mathcal{H} - \mathcal{E}\mathcal{D}] \mathbf{C} = \mathbf{0}$$

Cholesky



$$\mathcal{D} \rightarrow \mathbf{D}$$

$$[\mathbf{H} - \mathbf{E}] \mathbf{C} = \mathbf{0}$$



$$|n, \beta\rangle = \sum_{\lambda\alpha} C_{\lambda\alpha}^{\beta} \mathbf{O}_{\lambda}^{\dagger} |\alpha\rangle$$

$$\langle n', \alpha | n; \beta \rangle = \delta_{nn'} \delta_{\alpha\beta}$$

$$\mathcal{H} = \mathcal{A}\mathcal{D}$$

$$\mathcal{D} \equiv \{ \langle \lambda' \alpha' | \lambda \alpha \rangle \}$$

$$\mathbf{H} = \mathbf{D}^{-1} \mathcal{A}\mathcal{D}$$

Iterative Generation of n-phonon states

$$\mathbf{A} \mathbf{c} = \hbar \boldsymbol{\omega} \mathbf{c}$$



$$|\lambda\rangle = \sum_{\text{ph}} \mathbf{c}_{\text{ph}}^{\lambda} \mathbf{a}_{\text{p}}^{\dagger} \mathbf{a}_{\text{h}} | \rangle$$



$$\mathbf{H} \mathbf{C} = \mathbf{E} \mathbf{C}$$



$$|\alpha_n\rangle = \sum_{\lambda\alpha(n)} \mathbf{C}_{\lambda\alpha(n)}^{\alpha(n)} \mathbf{O}_{\lambda}^{\dagger} | \alpha_{n-1} \rangle$$

- **No approximations except for truncation!**
- **Pauli** principle fully accounted for:
- **No redundant** states!
- $|\alpha_n\rangle$ form an **orthonormal** basis

$$n = 2, 3, \dots$$

Eigenvalue problem in Multiphonon basis $\{|\alpha_n\rangle\} \equiv \{|\alpha_0\rangle, |\alpha_1\rangle, \dots\}$

$$\sum_{n', \beta(n')} [(\mathbf{E}_{\alpha(n)} - \epsilon_v) \delta_{nn'} \delta_{\alpha(n) \beta(n')} + \mathbf{V}_{\alpha(n) \beta(n')}] \mathbf{C}_{\beta(n')}^v = \mathbf{0}$$



$$|\Psi_v\rangle = \sum_{\alpha(n)} \mathbf{C}_{\alpha(n)}^v |\alpha_n\rangle$$

where

$$|\alpha_n\rangle = \sum_{\lambda \alpha(n-1)} \mathbf{C}_{\lambda \alpha(n-1)}^\beta \mathbf{O}_{\lambda}^\dagger |\alpha_{n-1}\rangle$$

EMPM in q-p scheme

One proceeds as in p-h scheme

$$\langle \alpha_n | [\mathbf{H}, \mathbf{O}_\lambda^\dagger] | \alpha_{n-1} \rangle = (E_\beta^{(n)} - E_\alpha^{(n-1)}) \langle \alpha_n | \mathbf{O}_\lambda^\dagger | \alpha_{n-1} \rangle$$

namely with the replacement

$$\mathbf{O}_\lambda^\dagger = \sum_{ph} c_{ph}(\lambda) \mathbf{a}_p^\dagger \mathbf{a}_h \quad \Rightarrow \quad \mathbf{O}_\lambda^\dagger = \sum_{r \leq s} c_{rs}(\lambda) \alpha_r^\dagger \alpha_s$$

Implementation

1° STEP : Intrinsic Hamiltonian

$$H = T_{\text{int}} + V_{\text{NN}}$$

where

$$T_{\text{int}} = \frac{1}{2m} \sum_i p_i^2 - \frac{P^2}{2M_{\text{cm}}} \quad V_{\text{NN}} = V_{\chi} = \text{NNLO}_{\text{opt}}$$

2° STEP : HF(B) Self-consistent basis

3° STEP : Construction of TDA phonons

(free of spurious admixtures induced by CM and particle number violation)

$$O_{\lambda}^{\dagger} = \sum_{\text{ph}} c_{\text{ph}}(\lambda) a_p^{\dagger} a_h$$

$$O_{\lambda}^{\dagger} = \sum_{\underline{r} \leq \underline{s}} c_{rs}(\lambda) \alpha_r^{\dagger} \alpha_s^{\dagger}$$

Implementation

4° STEP : Generation of Multiphonon basis

$$\{|\alpha_n\rangle\} \equiv \{|\alpha_0\rangle, |\alpha_1\rangle, |\alpha_2\rangle, \dots\}$$

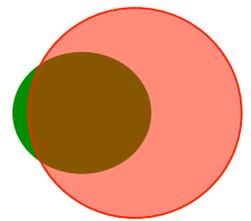
5° STEP : Solution of the eigenvalue problem in Multiphonon basis

$$|\Psi_\nu\rangle = \sum_{n\alpha} C_\alpha^{(\nu)} |n; \alpha\rangle$$

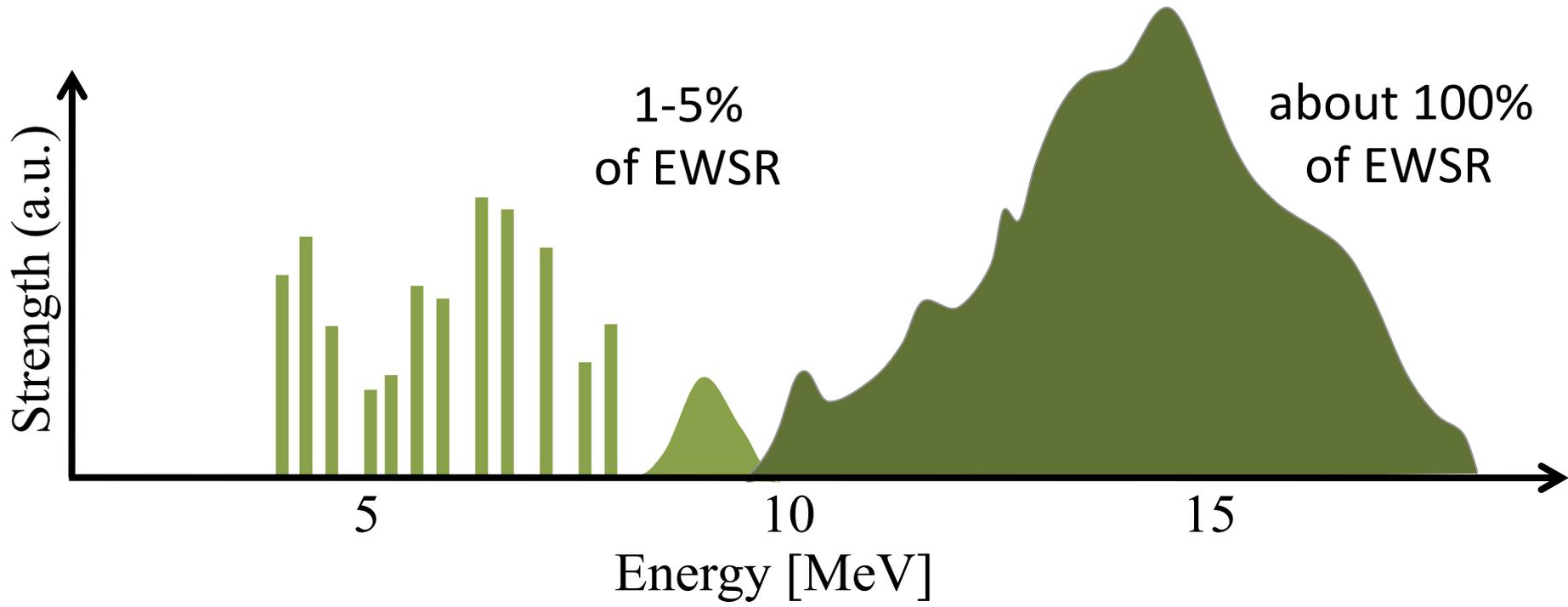
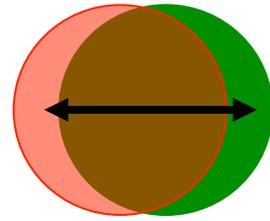
6° STEP : Transition Probabilities and cross section

Applications: Selfconsistent study of Dipole response

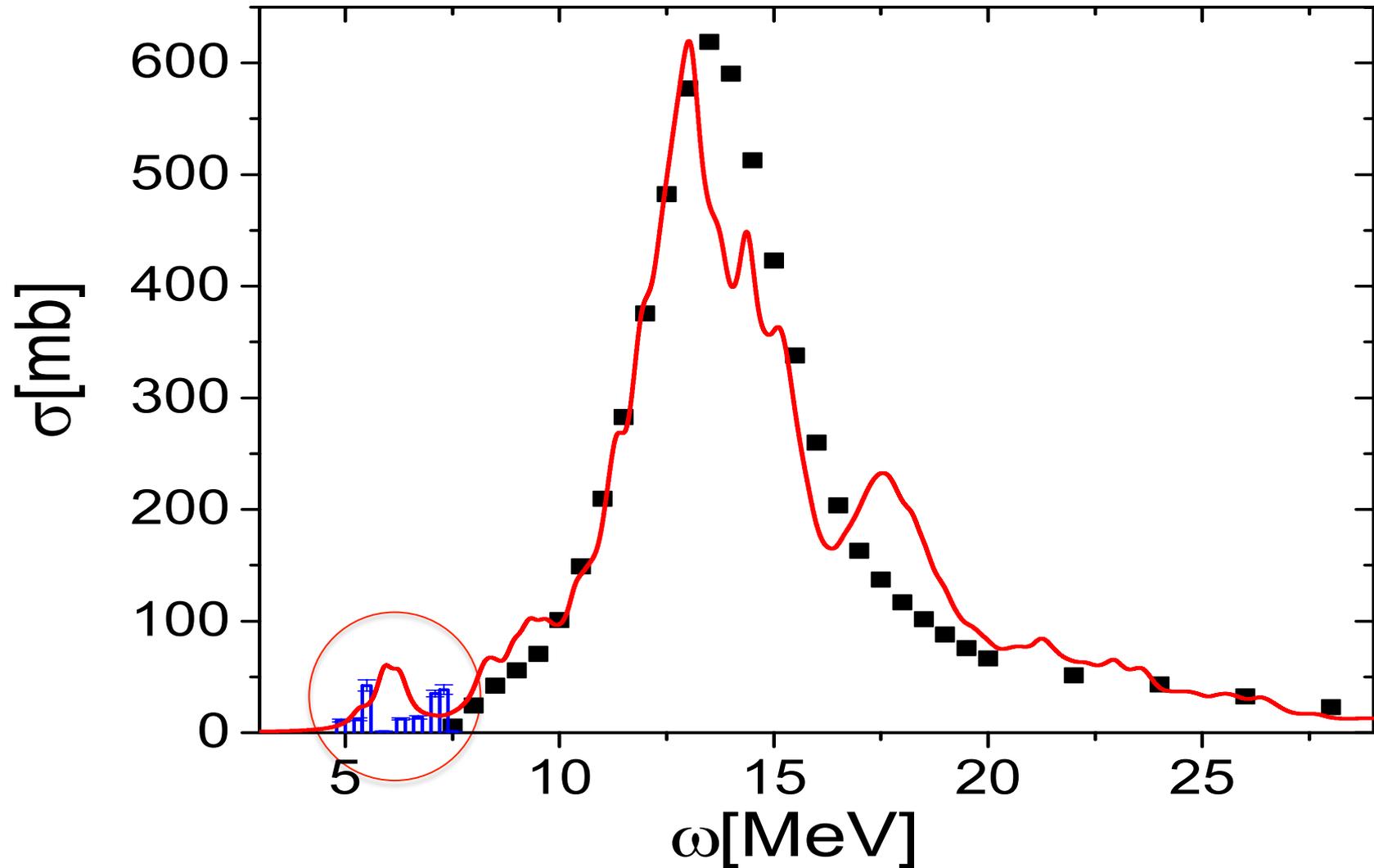
PDR



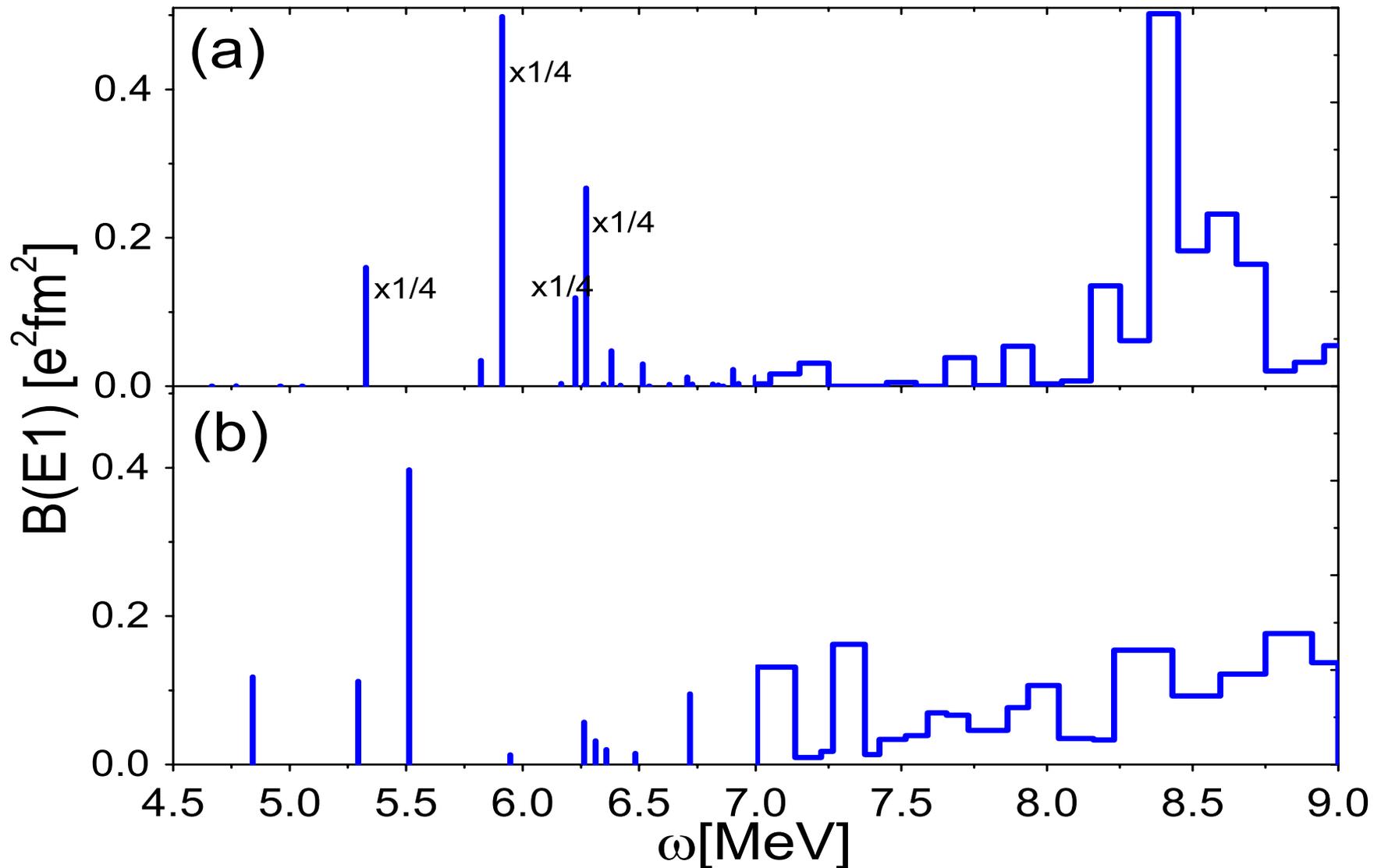
GDR



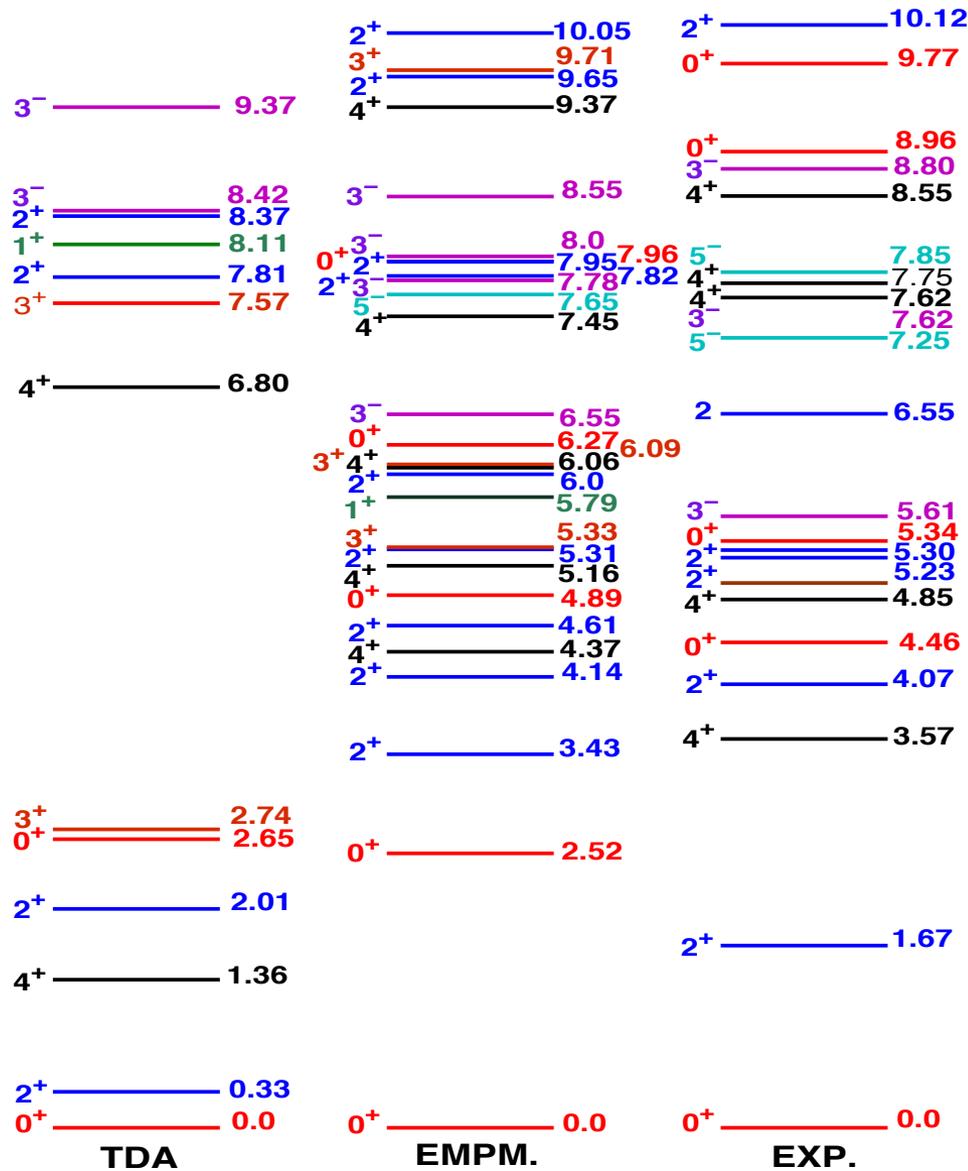
EMPM in p-h scheme: Application to neutron rich ^{208}Pb



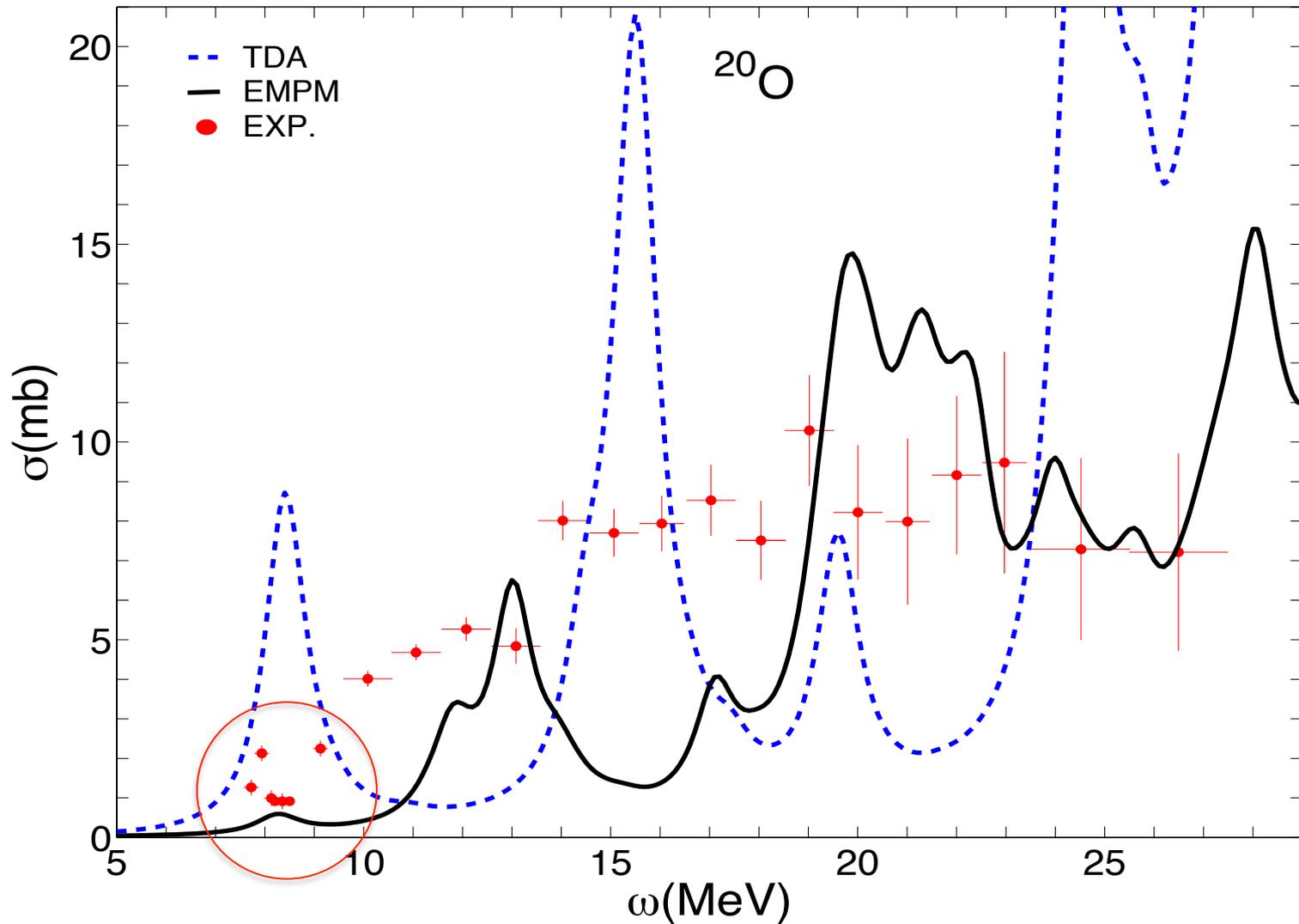
Low-lying Dipole strength in ^{208}Pb



EMPM in q-p scheme: Application to neutron rich ^{20}O



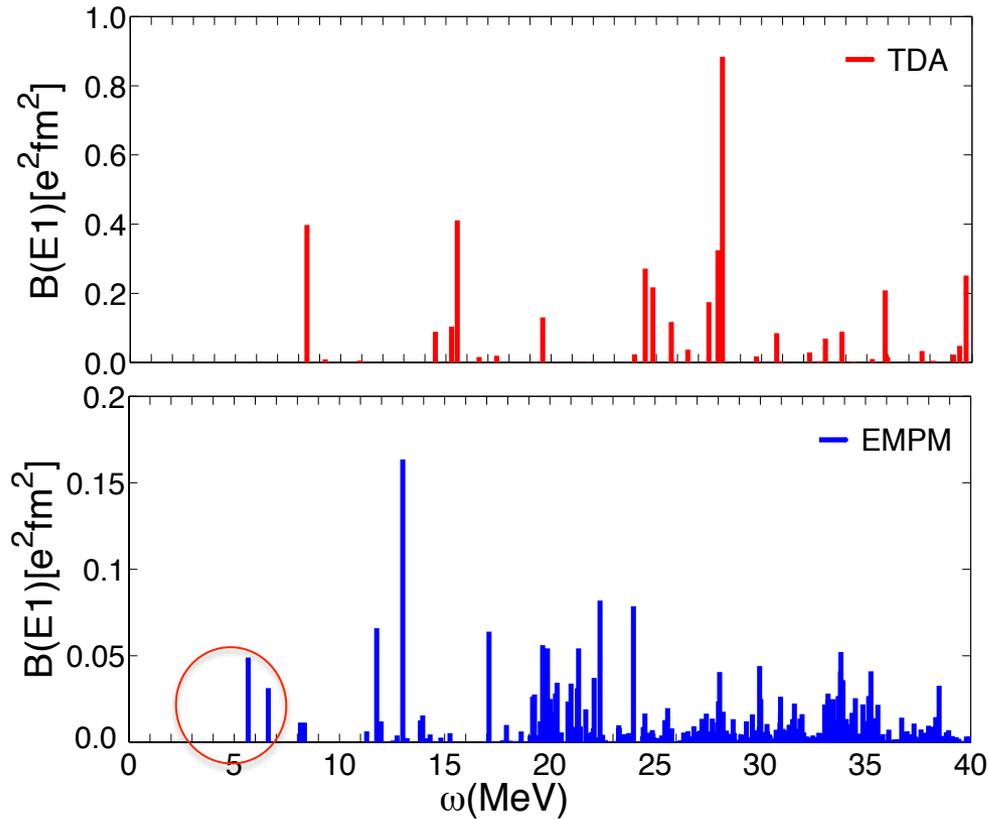
Cross section in ^{20}O



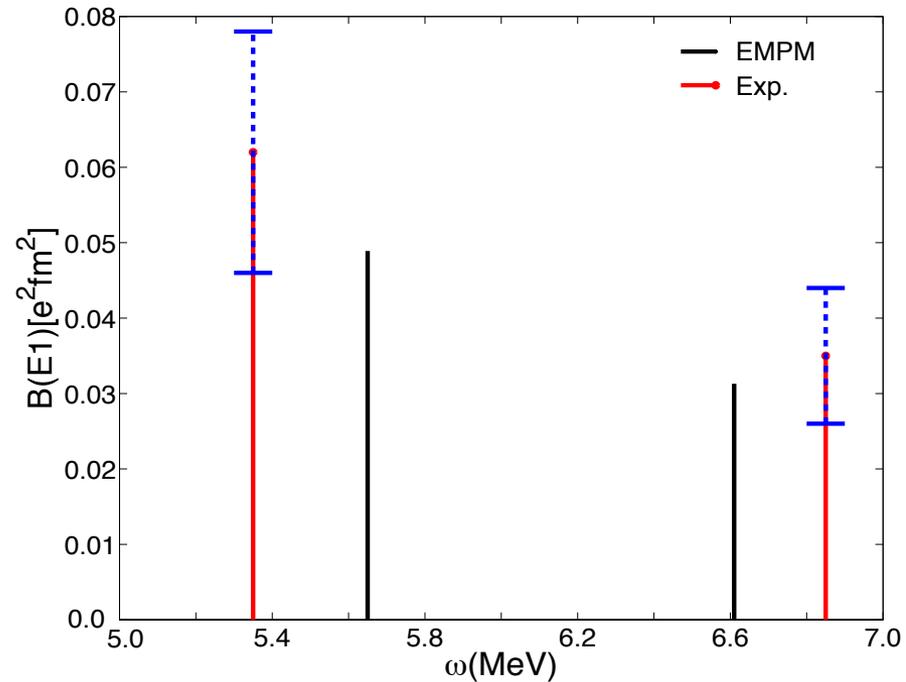
G. De Gregorio, F. Knapp, N. Lo Iudice, P. Vesely
Phys. Rev. C 044314 (2016)

E. Tryggestad, et al., Phys. Rev. C
67, 064309 (2003).

Dipole strength in ^{20}O



E. Tryggestad, et al., Phys Rev. C 67, 064309 (2003).



G. De Gregorio, F. Knapp, N. Lo Iudice, P. Vesely
Phys. Rev. C 044314 (2016)

Even Nuclei: Summary

EMPM

- Extends Mean Field without **approximations**, except for truncation, for any potential
- Includes **multiphonon states** **explicitly**

Results

2-phonon configurations are **needed** for

- describing the **fine** structure of the GDR and PDR
 - obtaining a **complete** low-energy spectrum

EMPM: Odd nuclei

$$|\lambda \alpha\rangle = O^\dagger_\lambda |\alpha\rangle$$



$$|p \alpha\rangle = a^\dagger_p |\alpha\rangle$$

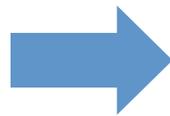
$$\langle \alpha_n | [H, O^\dagger_\lambda] | \alpha_{n-1} \rangle$$



$$\langle v_n | [H, a^\dagger_p] | \alpha_n \rangle$$

$$X = DC$$

$$AX = EX$$



$$[H - ED]C = 0$$

$$D \equiv \{ \langle p \alpha | p' \alpha' \rangle \}$$

Cholesky

$$D \rightarrow D$$

EMPM: Odd nuclei

$$[\mathbf{H} - \mathbf{E}] \mathbf{C} = \mathbf{0}$$

$$\mathbf{H} = \mathbf{D}^{-1} \mathcal{A} \mathcal{D}$$

$$|\mathbf{v}_n\rangle = \sum_{p\alpha} C_{p\alpha}^{\mathbf{v}} a_p^\dagger |\alpha\rangle$$

$$\langle \mathbf{v}' | \mathbf{v} \rangle = \delta_{\mathbf{v}\mathbf{v}'}$$

$$|\Psi_\mu\rangle = \sum_i C_{\mathbf{v}_i}^{(\mu)} |\mathbf{v}_i\rangle$$

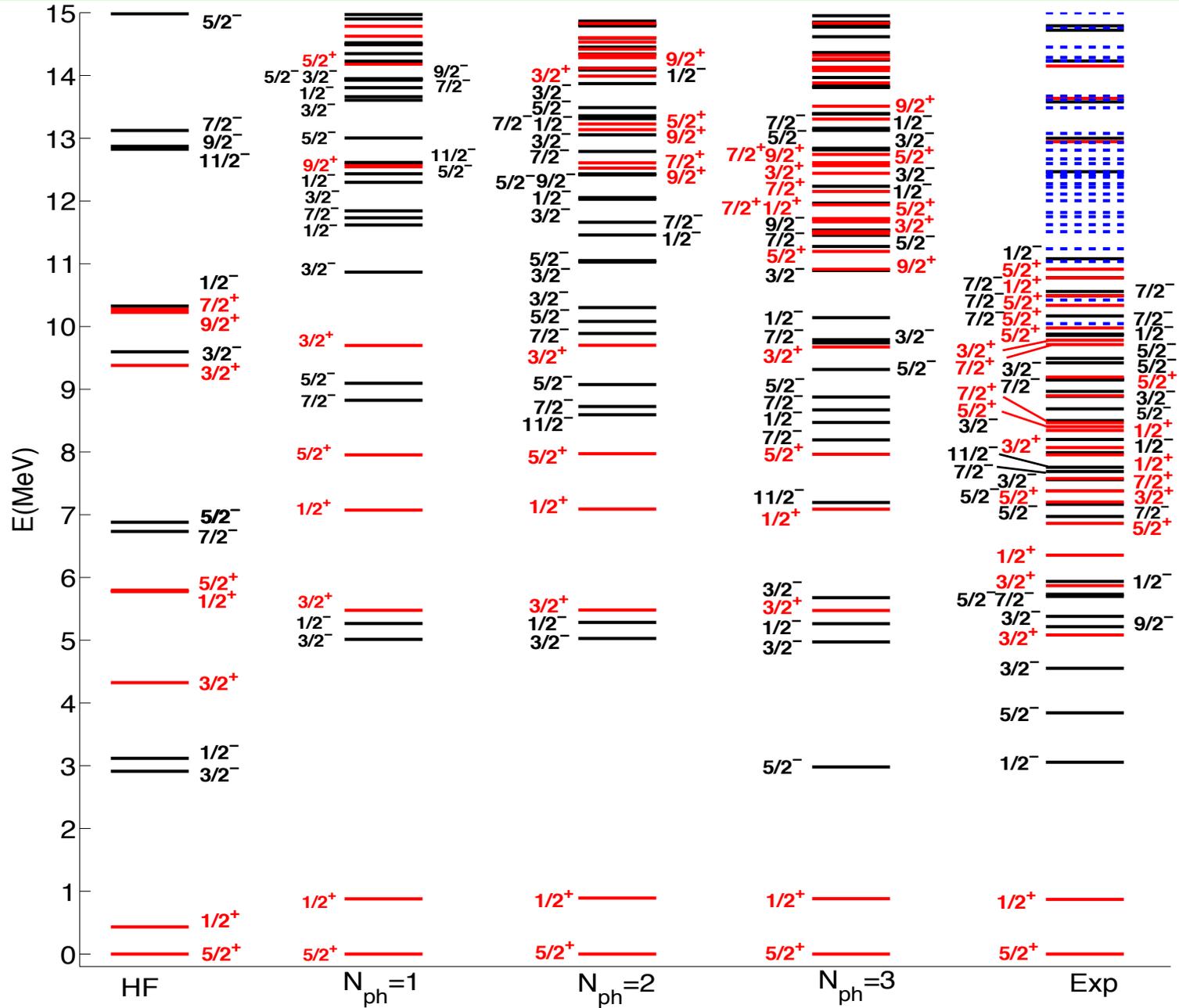
- **No approximations except for truncation !**
- **Pauli** principle fully accounted for:
- **No redundant** states!
- $|\mathbf{v}_n\rangle$ form an **orthonormal** basis

Implementation

Potential: NNLO_{opt}

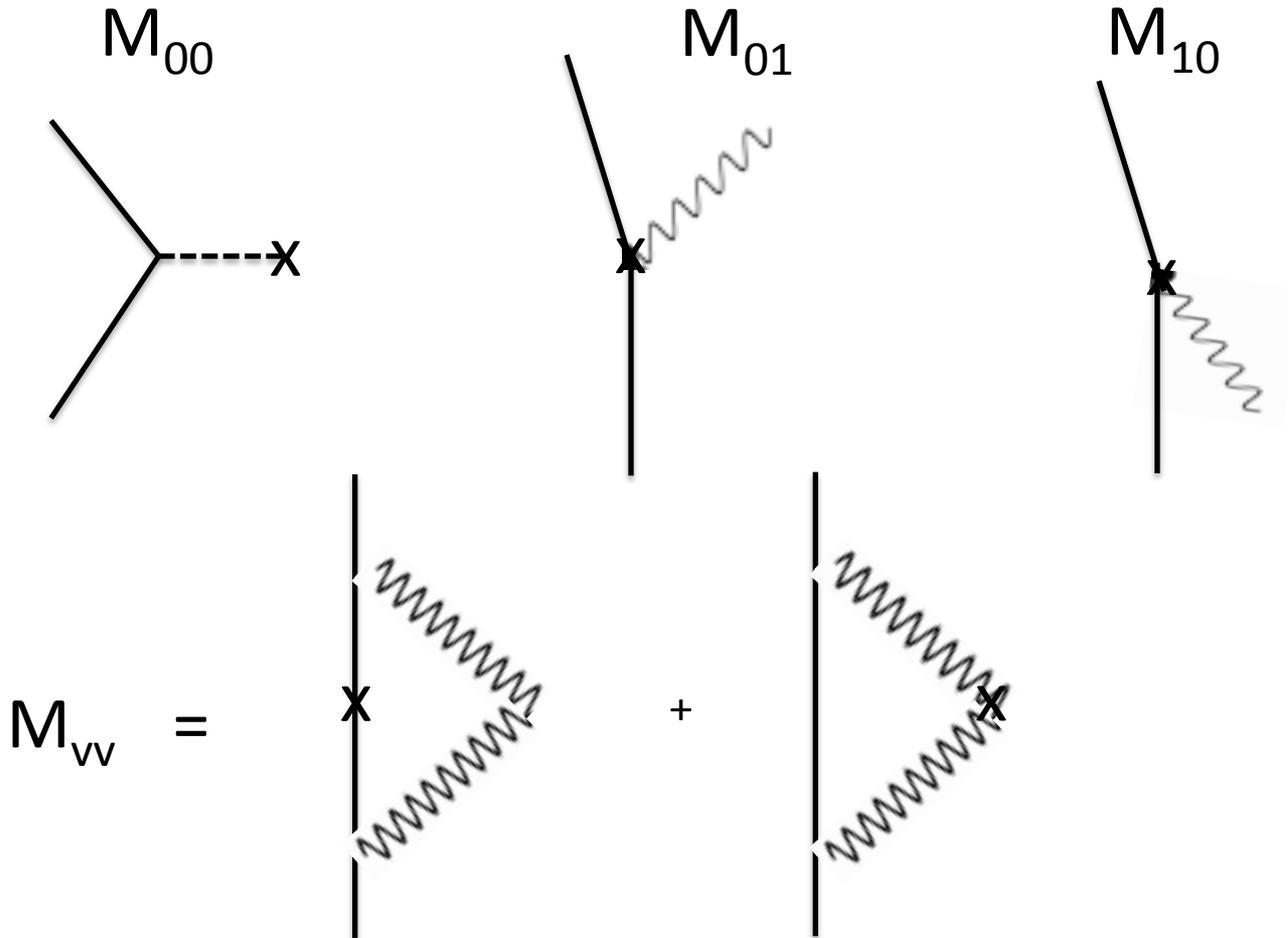
- Perform **HF**
- Construct **TDA** phonons (**free** of **CM** spurious admixtures)
- Generate the multiphonon basis $\{|\alpha_n\rangle\} \equiv \{|\alpha_0\rangle, |\alpha_1\rangle, |\alpha_2\rangle, |\alpha_3\rangle\}$
- Generate the **orthonormal** particle-phonon basis $\{|\mathbf{v}_n\rangle\}$
- Full diagonalization

^{17}O spectrum

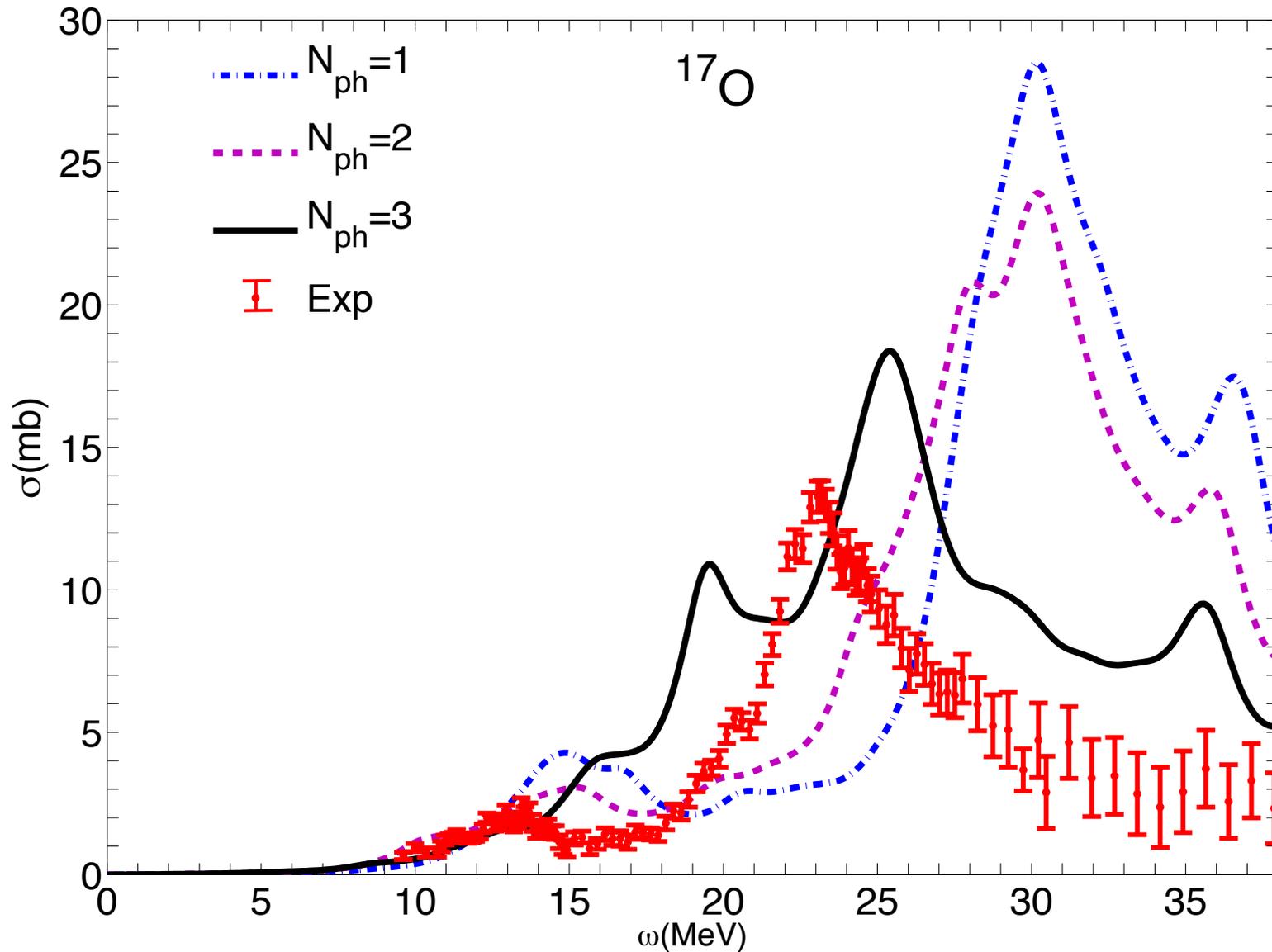


Transitions Amplitudes

$$\langle \Psi_{\Omega_f} \| M(\lambda) \| \Psi_{\Omega_i} \rangle = M_{00}(\lambda) + M_{01}(\lambda) + M_{10}(\lambda) + M_{11}(\lambda) + M_{12}(\lambda) + M_{21}(\lambda) + M_{22}(\lambda)$$

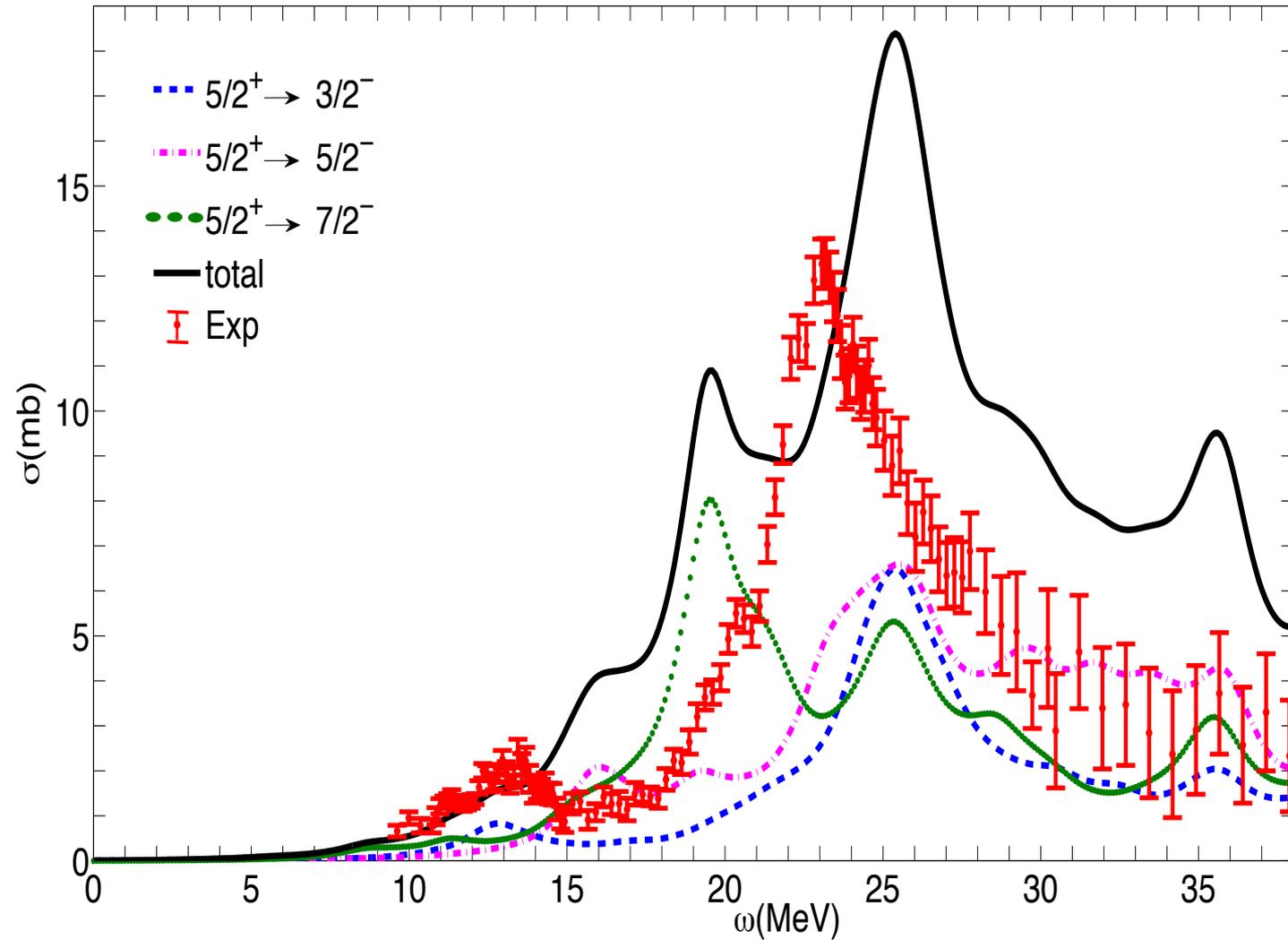


Cross section



J. W. Jury, et al., Phys Rev. C **21**, 503 (1980).

Cross section



Mean Values and transitions

^{17}O	HF	EMPM	Exp ⁽¹⁾
Q(barn)	0	-0.00841	-0.025
μ (nm)	-1.9130	-1.833	-1.893
$B(E2; 5/2^+_{1} \rightarrow 1/2^+_{1})$ ($e^2\text{fm}^4$)	0	0.17	2.18(16)

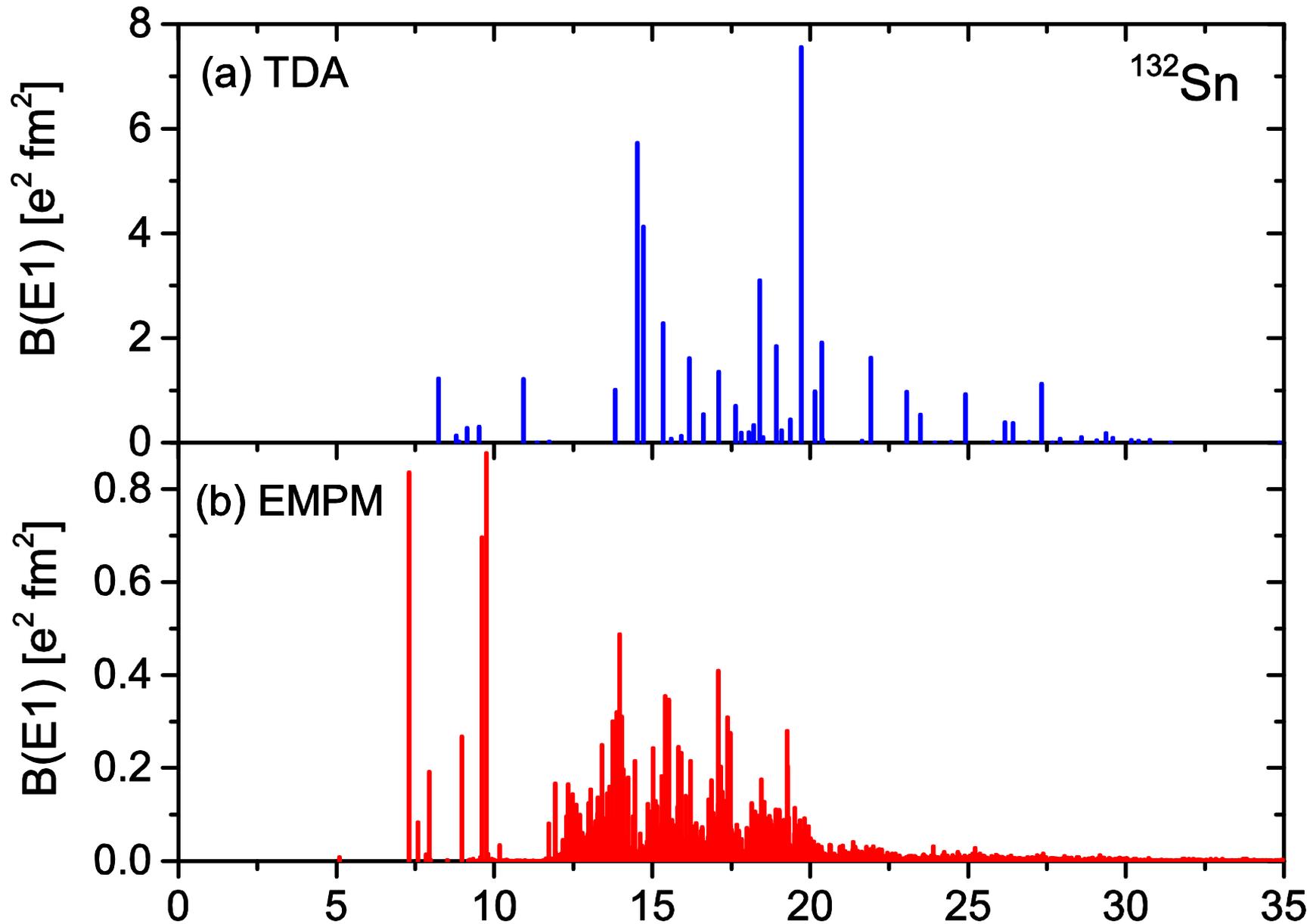
¹ N.P. Stone, Atomic Data and Nuclear Data Tables **90** (2005) 75–176

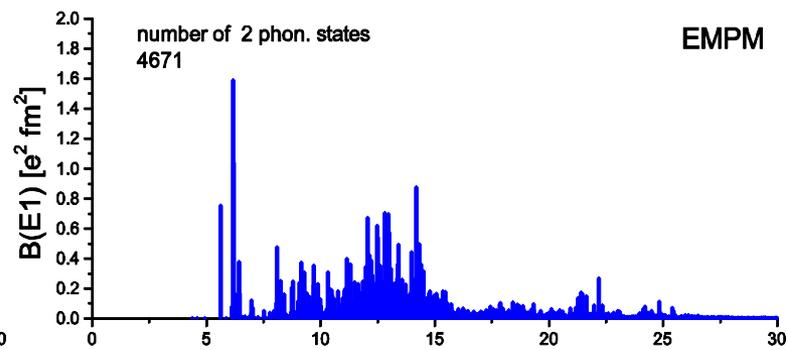
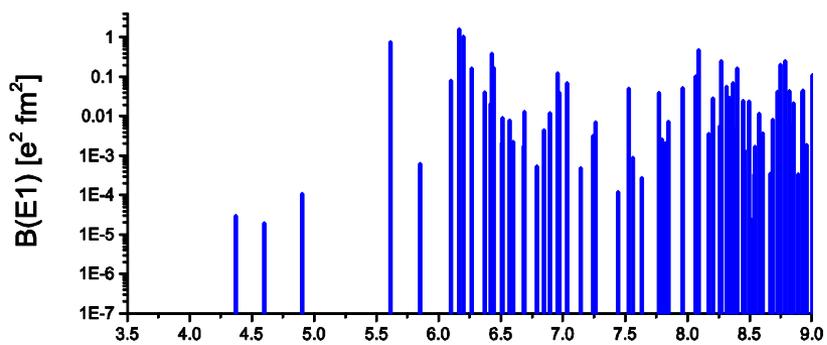
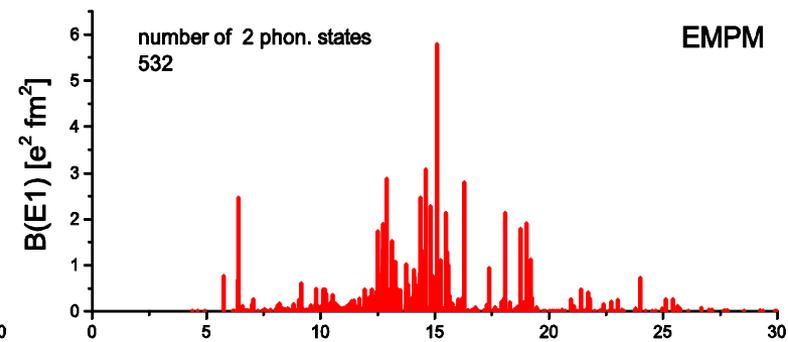
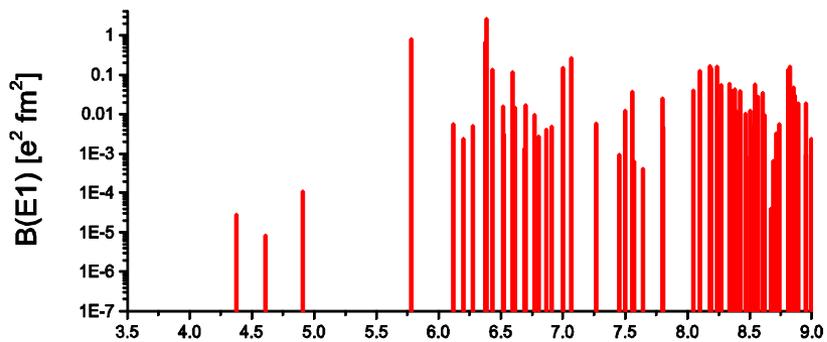
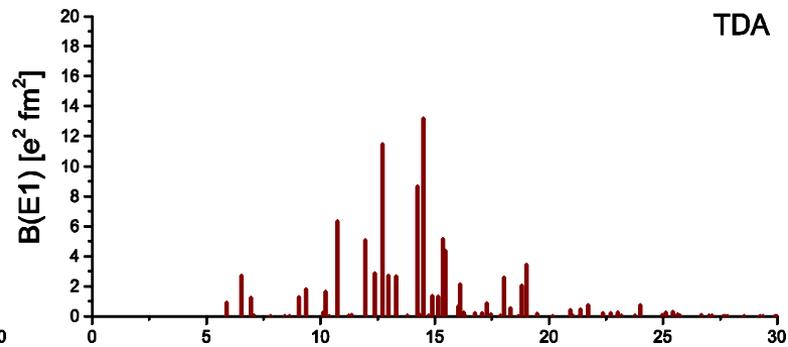
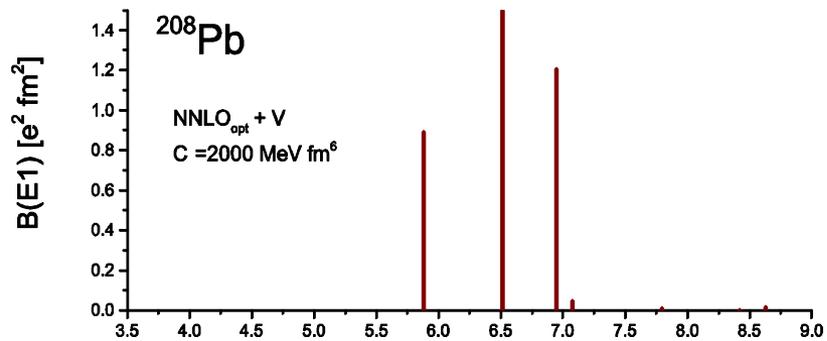
Odd nuclei : Concluding remarks

- One- and two-phonon states enhance greatly the density of levels consistently with experiments.
- Three-phonon states play an important role, their (strong) coupling to particle-phonon states $|v_1\rangle$
 - a) shift downwards the energies of $|v_1\rangle$
 - b) improve the cross section
 - c) increase the density of levels in low-energy part of the spectrum

Thank you

Dipole strength and cross section in ^{132}Sn





[MeV]

[MeV]