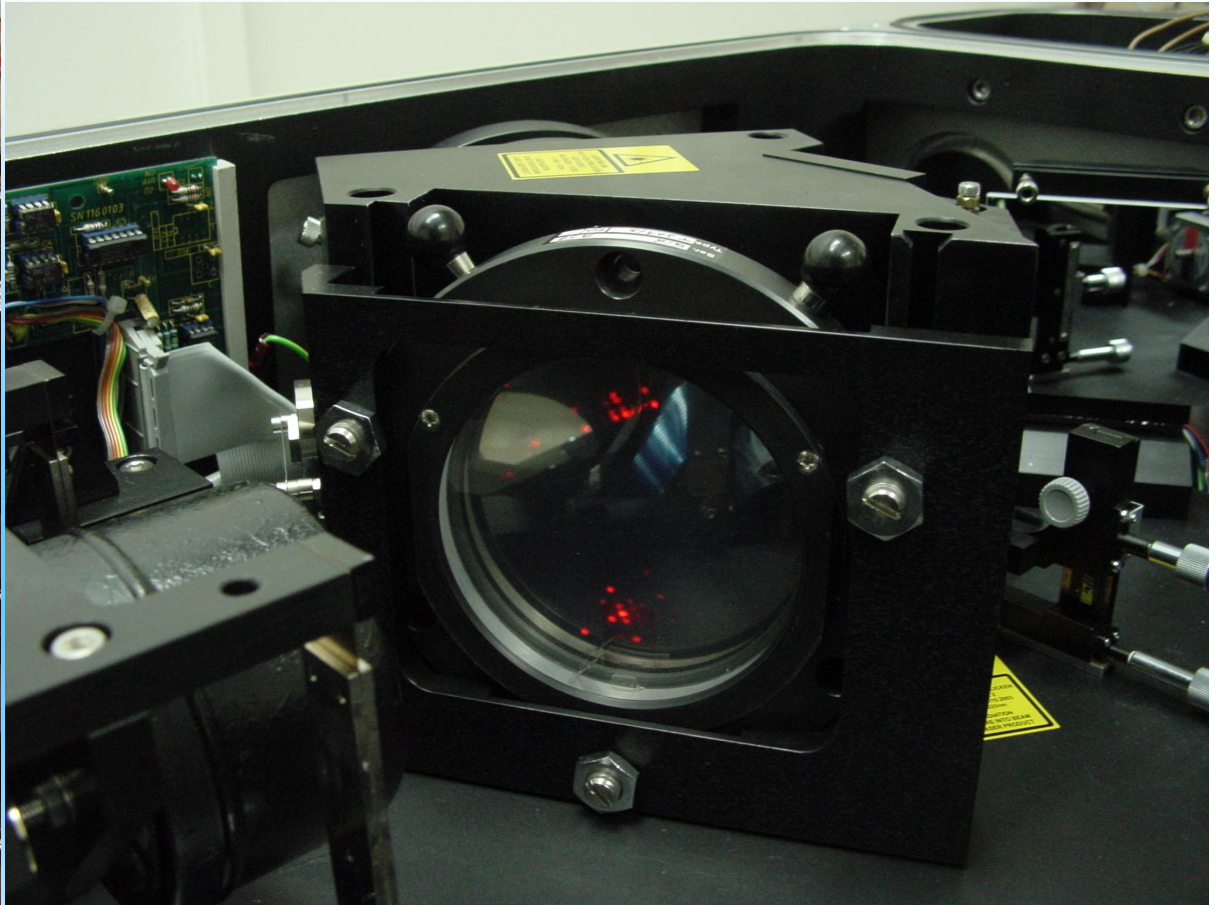
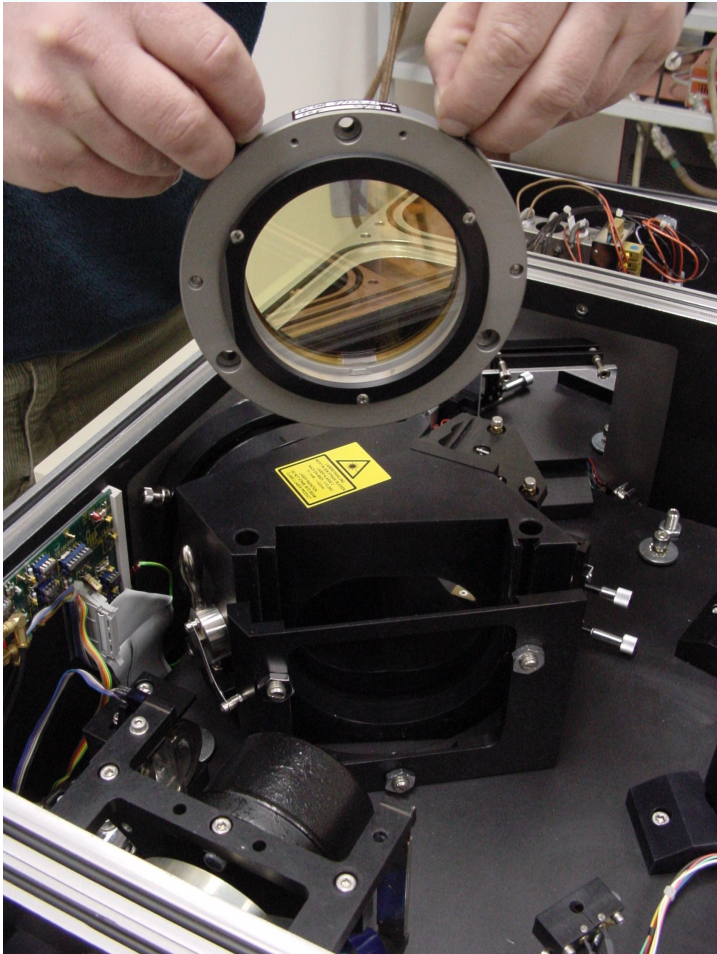


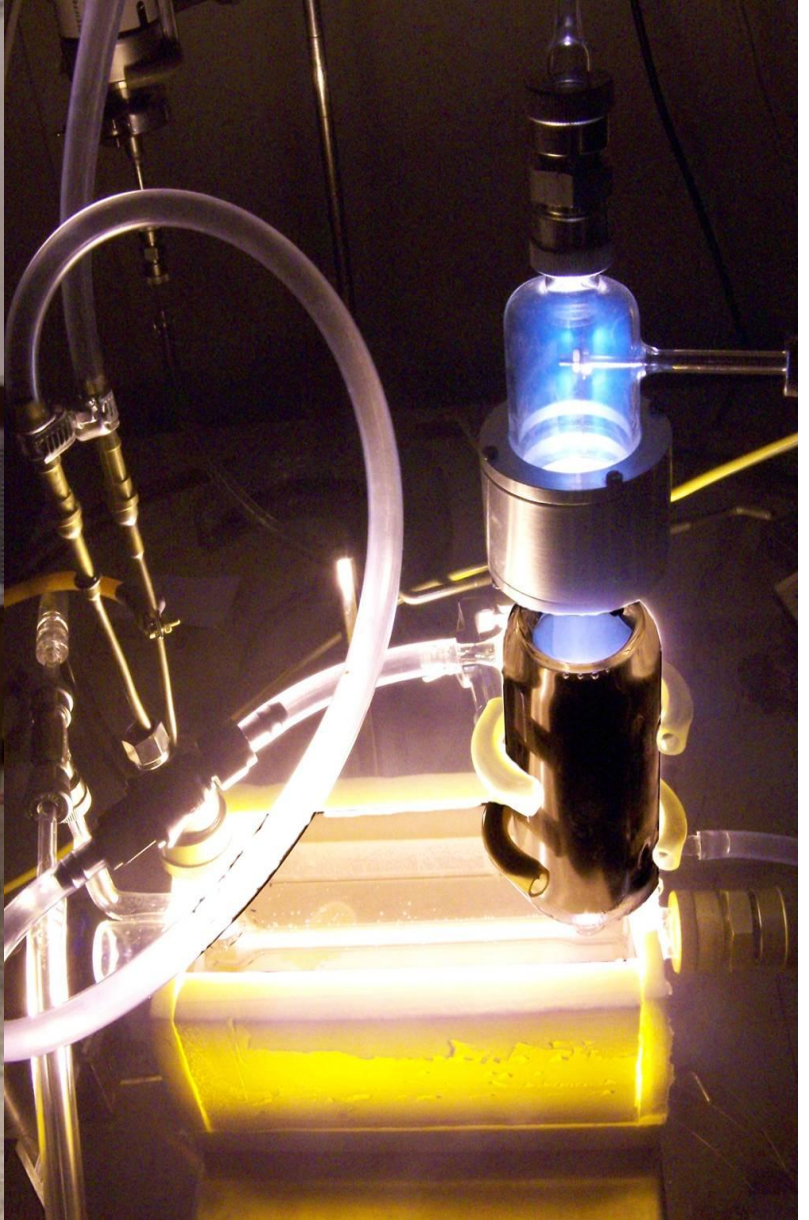
Svatopluk Civi–

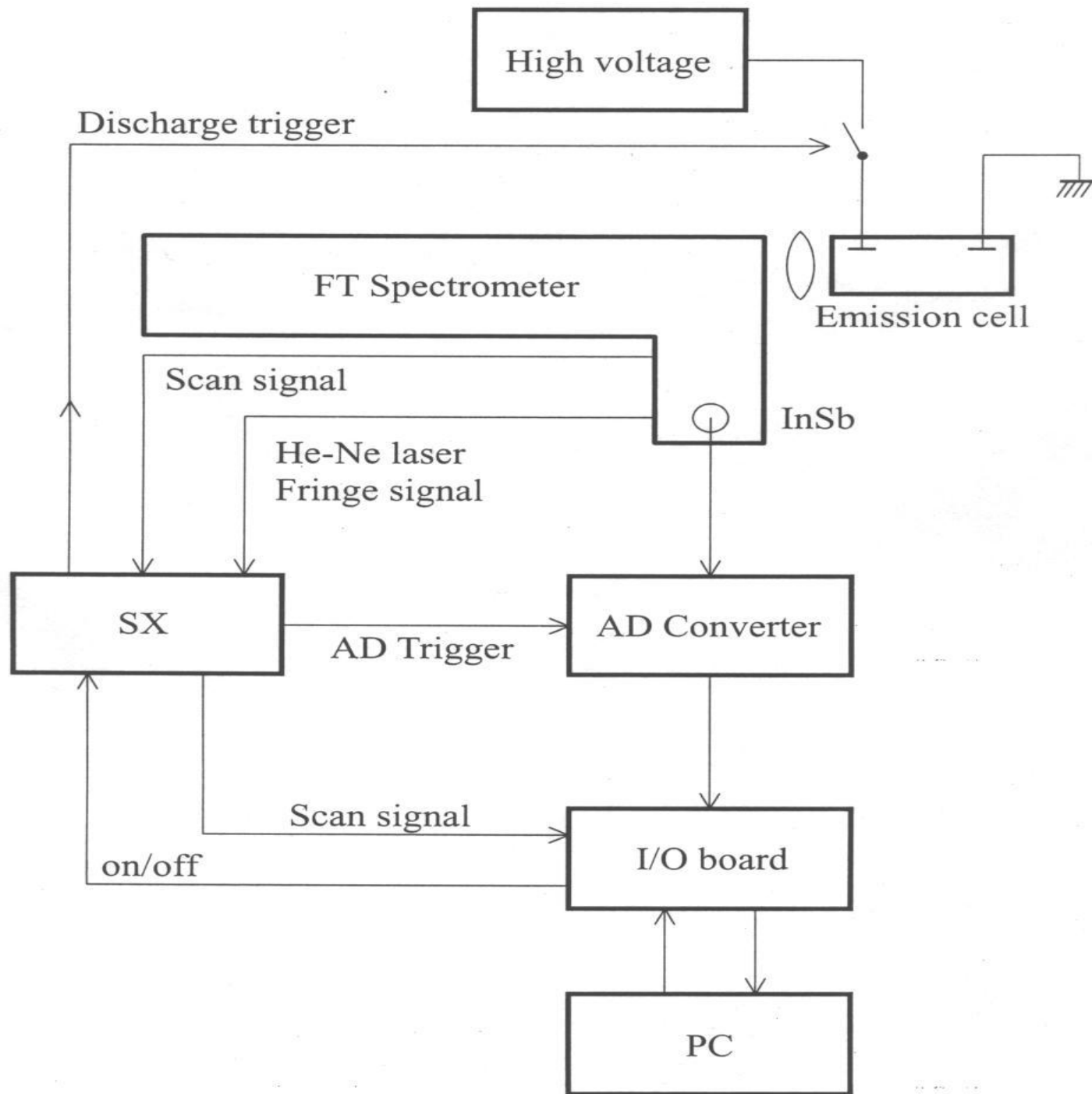


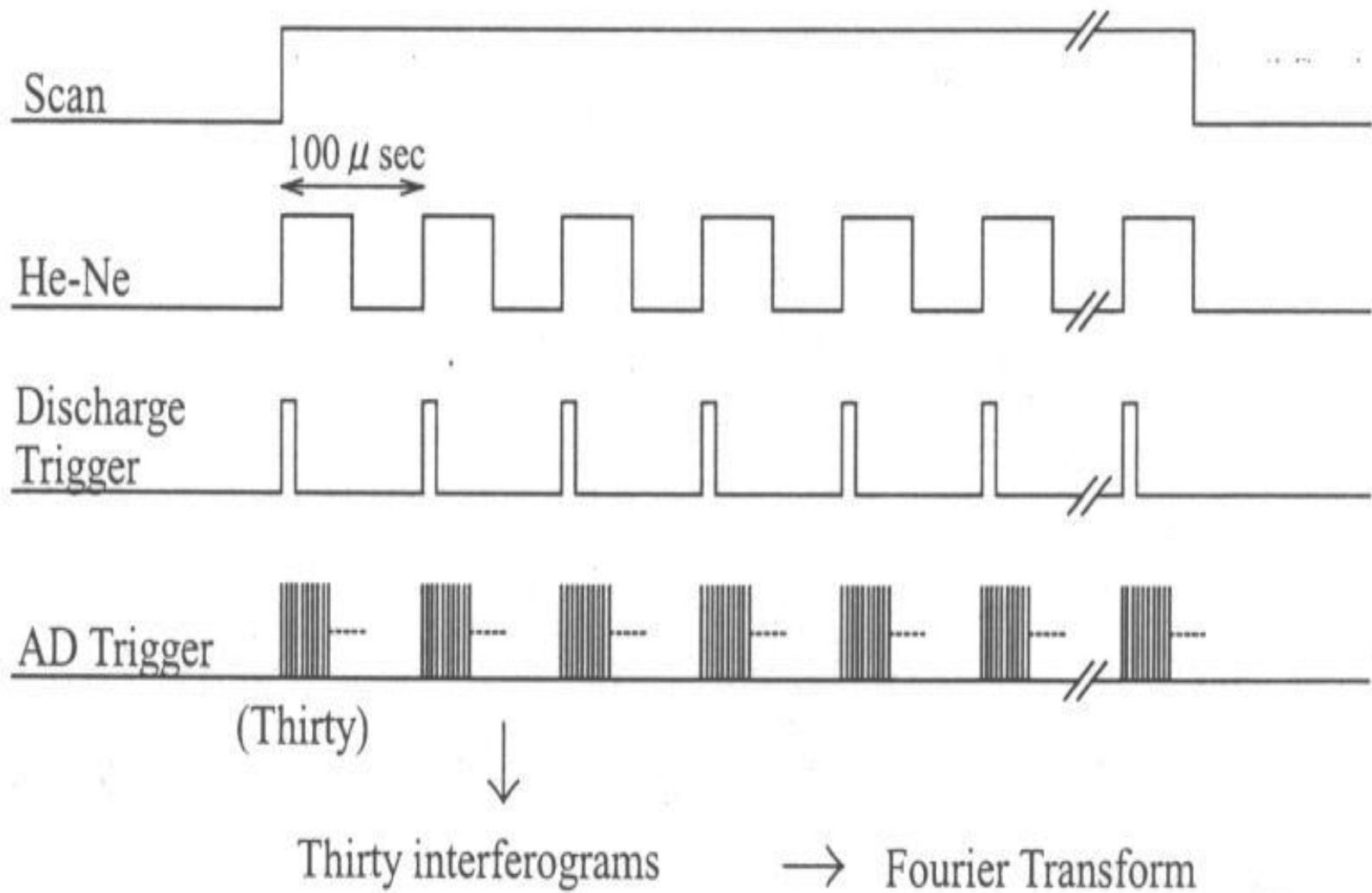
**š asov rozli-ená infra ervená emisní spektroskopie
s Fourierovou transformací a její aplikace v pulzních výbojích a
laserové ablaci kov ů.**



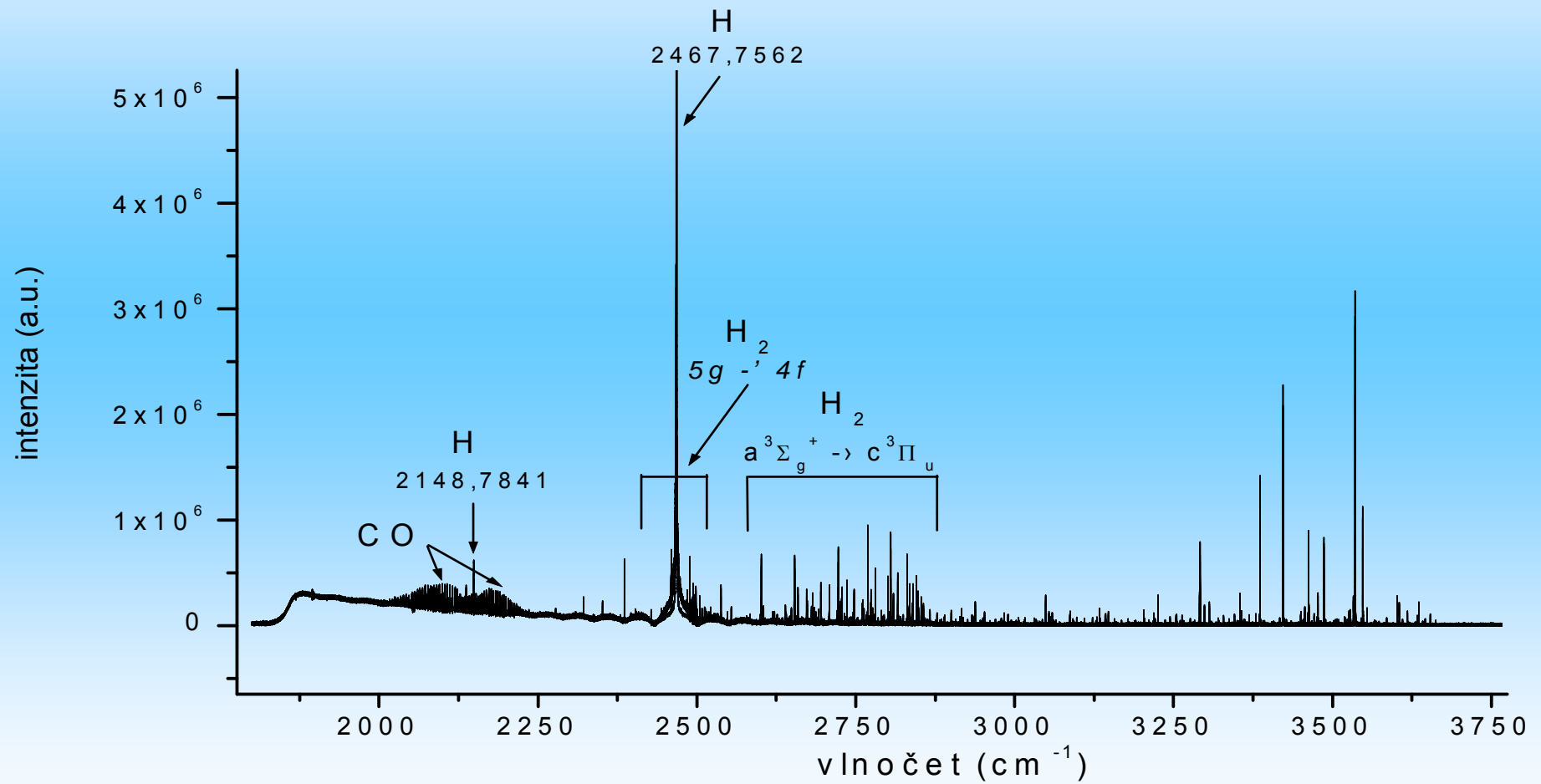




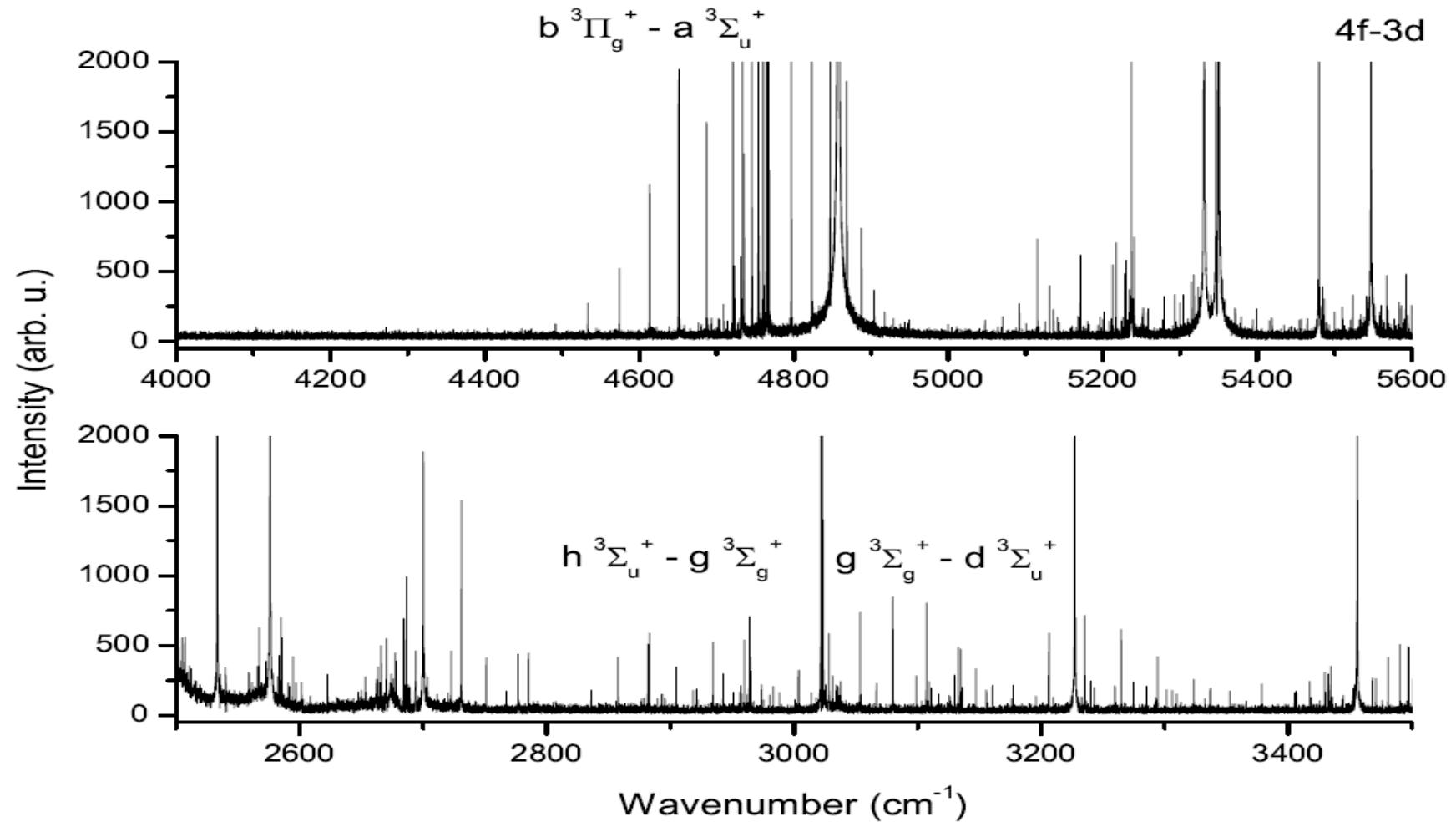




Výboj v čistém vodíku



Helium discharge



He discharge plasma

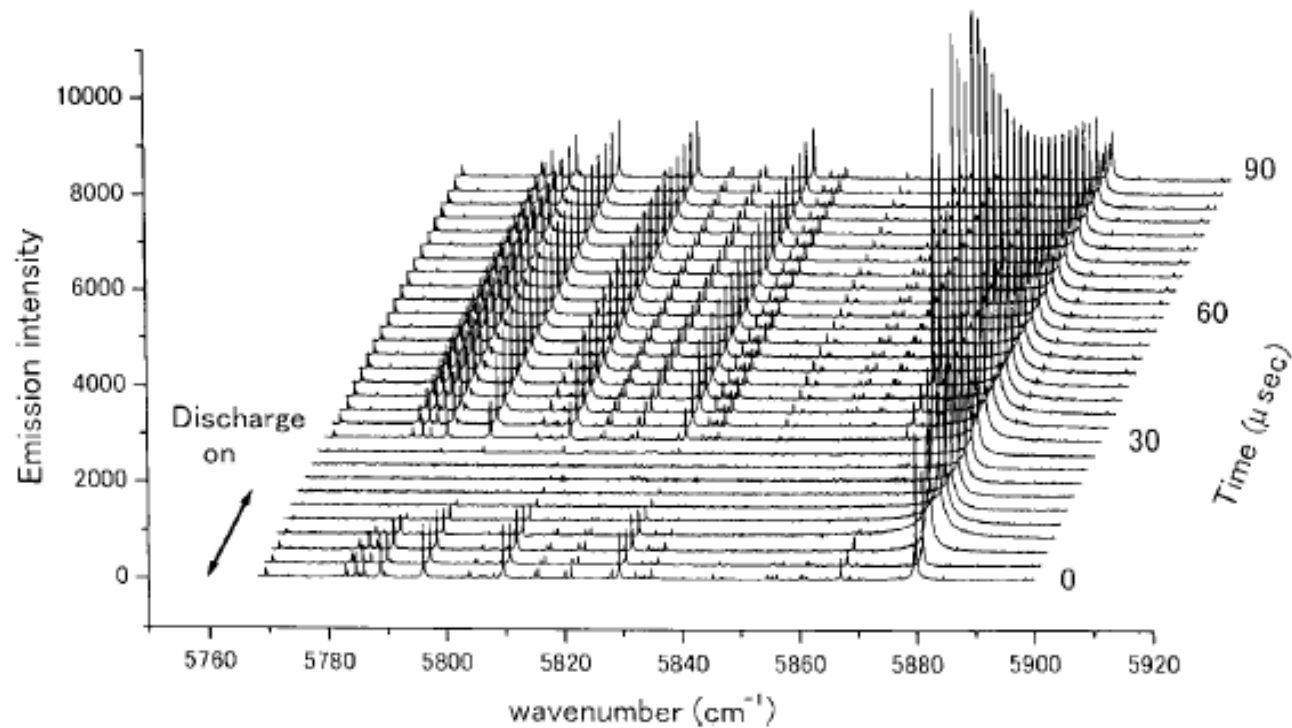


Fig. 1. A portion of the time-resolved spectrum observed by a pulsed discharge in He with a pressure of 1.33 kPa (10 Torr). The discharge was applied in the interval of 0–20 μs with a peak current of 0.5 A. The strongest peak belongs to the atomic He line (4d–3p). Other lines pertain to 4f–3d transitions of He₂.

He₂

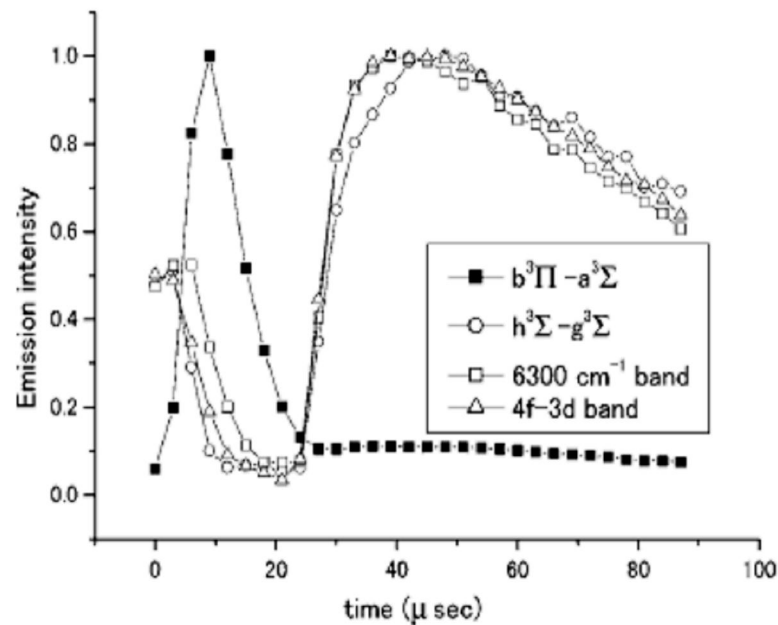
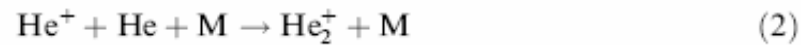
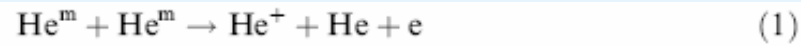


Fig. 3. Observed time profiles of emission intensities of He₂. Abscissa values show normalized intensities of the emission lines. The discharge conditions are as given in the caption of Fig. 1.

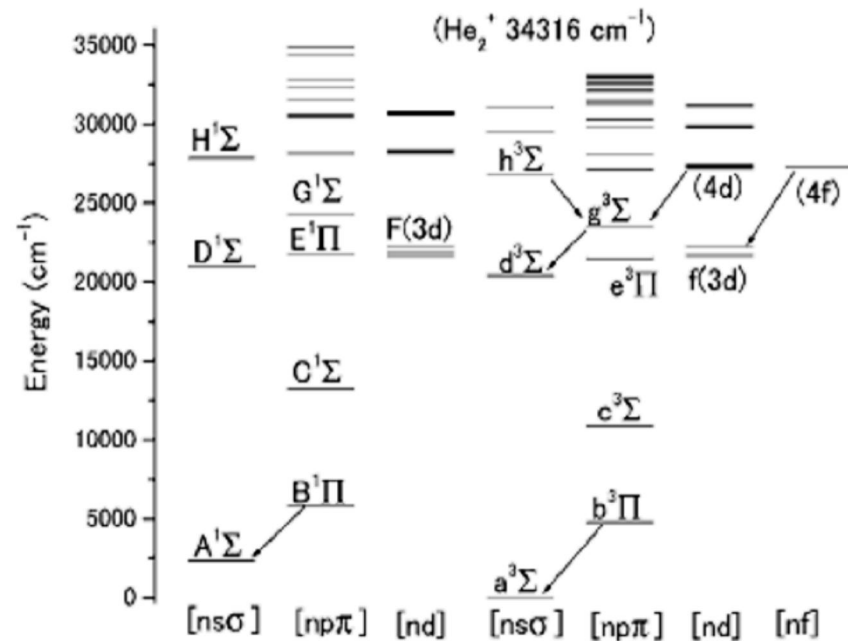


Fig. 4. Energy level diagram of He₂. The transitions observed in the present study are shown by arrows. The energy values are represented relative to the $a^3\Sigma_u^+ v = 0$ state. $n(>1)$ is the principal quantum number in a united atom molecular orbital designation. The ionization limit to He₂⁺ is 34 316 cm⁻¹.

He₂

Observed transitions of He₂ (cm⁻¹)^a

N	P(N)	o.-c.	R(N)	o.-c.
<i>h</i> ³ Σ_u^+ – <i>g</i> ³ Σ_g^+				
0	3177.5569	0.0006		
2	3134.9969	–0.0005	3206.4133	0.0018
4	3107.2556	–0.0001	3235.5447	0.0000
6	3080.1473	–0.0010	3264.8626	–0.0064
8	3053.7916	0.0000	3294.3028	0.0089
10	3028.3000	0.0012	3323.7108	–0.0054
12	3003.7698	–0.0002	3353.0067	0.0013
14	2980.2783	–0.0002		
<i>g</i> ³ Σ_g^+ – <i>d</i> ³ Σ_u^+				
1	3190.4064	–0.0019	3232.9664	–0.0013
3	3160.7770	–0.0003	3259.9345	0.0005
5	3130.2548	0.0017	3285.6522	0.0016
7	3098.9306	0.0021	3310.0069	0.0006
9	3066.8917	–0.0002	3332.8839	–0.0017
11	3034.2243	–0.0031	3354.1697	0.0005
13	3001.0169	0.0015		

^a N denotes the rotational quantum number neglecting spin in lower electronic states.

Molecular constants of He₂ in the *h*³ Σ_u^+ , *g*³ Σ_g^+ , and *d*³ Σ_u^+ states^a

	Present	Previous
(<i>h</i> ³ Σ_u^+)		^b
<i>B</i>	7.148 49(24)	7.149
<i>D</i> × 10 ³	0.505 3(24)	0.524
<i>E</i>	6368.1202(30)	
(<i>g</i> ³ Σ_g^+)		^c
<i>B</i>	7.096 423(94)	7.096 8(1)
<i>D</i> × 10 ³	0.530 70(44)	0.538(7)
<i>E</i>	3204.8589(11)	
(<i>d</i> ³ Σ_u^+)		^d
<i>B</i>	7.226 329(88)	7.228 6(15)
<i>D</i> × 10 ³	0.519 91(37)	0.532(3)
<i>E</i>	0.0	0.0

^a cm⁻¹ unit. Numbers in parentheses denote one standard deviation and apply to the last significant digits.

^b Ref. [1].

^c Ref. [12].

^d Ref. [13].

H₃⁺ FUNDAMENTAL BAND IN JUPITER'S AURORAL ZONES AT HIGH RESOLUTION FROM 2400 TO 2900 INVERSE CENTIMETERS

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Herzberg Institute of Astrophysics, NRC, Ottawa

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AND

J. CALDWELL¹

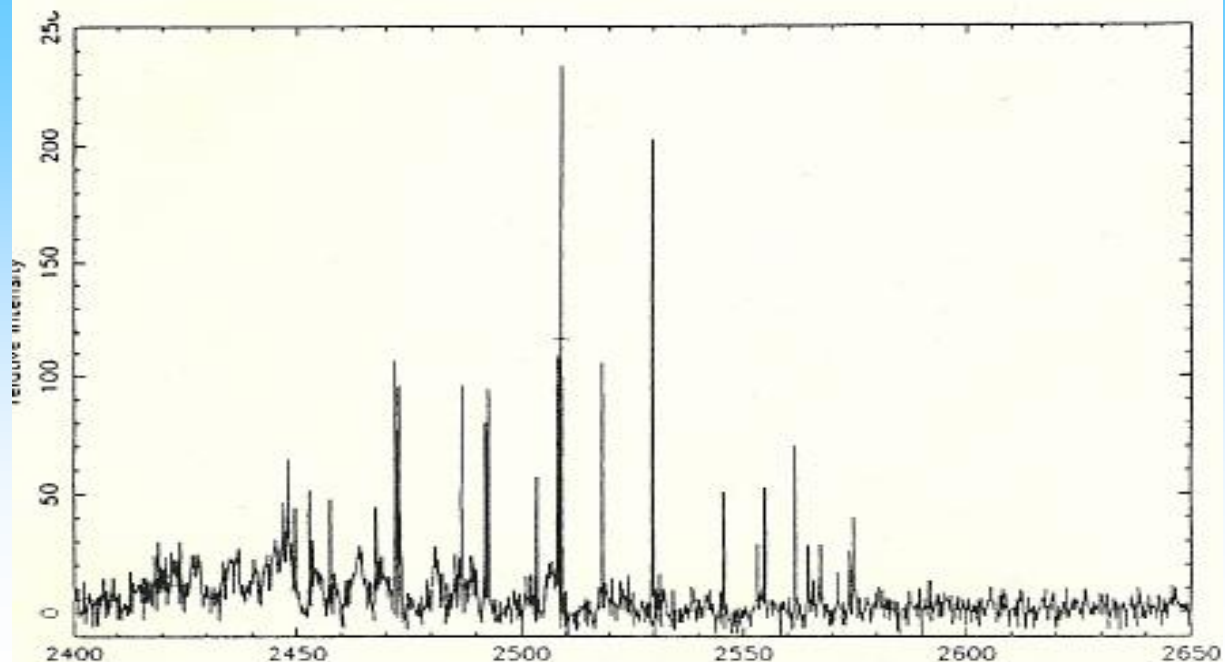
Physics Department, York University

Received 1990 June 22; accepted 1990 August 16

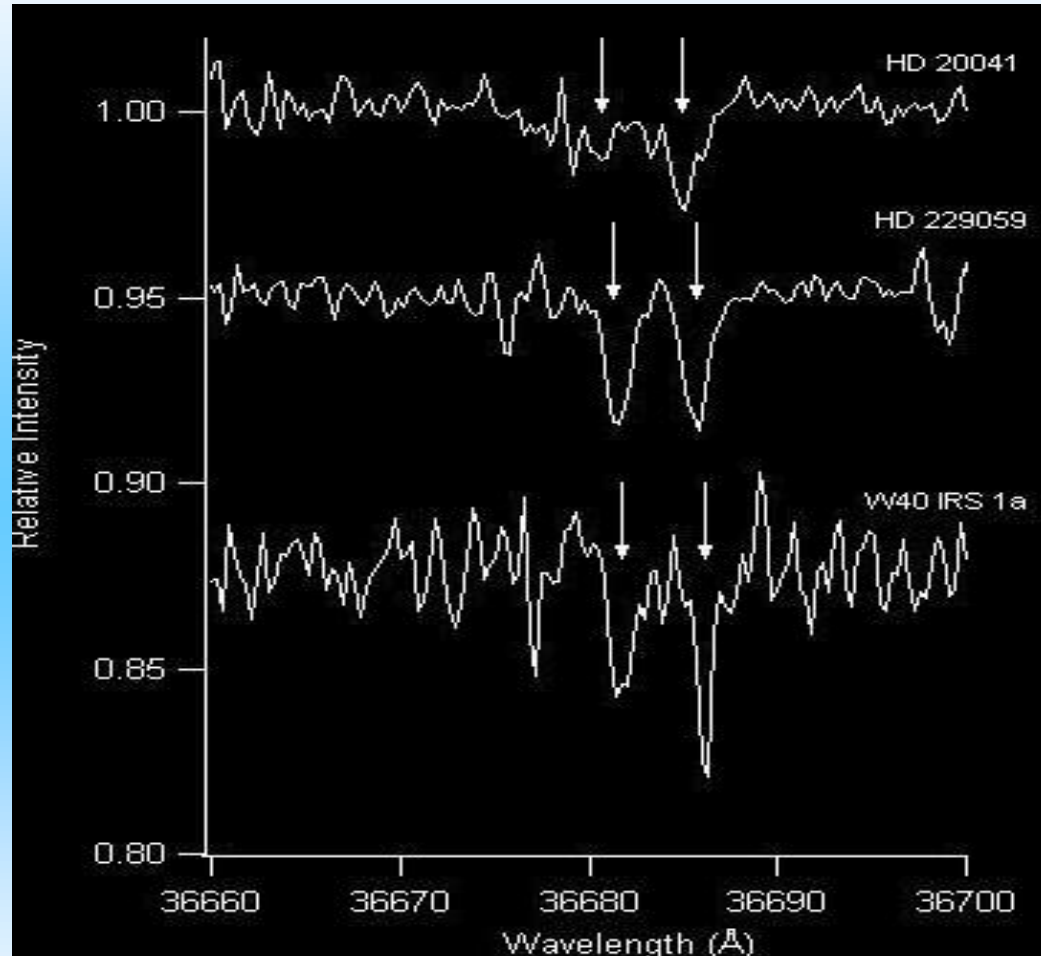
ABSTRACT

Following our previous detection of H₃⁺ in the southern auroral zone of Jupiter from its 2ν₂ band, we searched for the fundamental at 4 μm. We detected up to 42 lines of this band in emission, at high resolution, on the auroral spot of each hemisphere. A rotational temperature was derived for the southern and northern zones, respectively, of 1000 ± 40 K and 835 ± 50 K. The intensity of the lines was on the average 2 times stronger in the south than in the north. The 2ν₂ band, which was sought in the north only on this occasion, was not detectable. A purely thermal mechanism for the H₃⁺ production is implied. Spatial extension and temporal variability of the excitation is discussed.

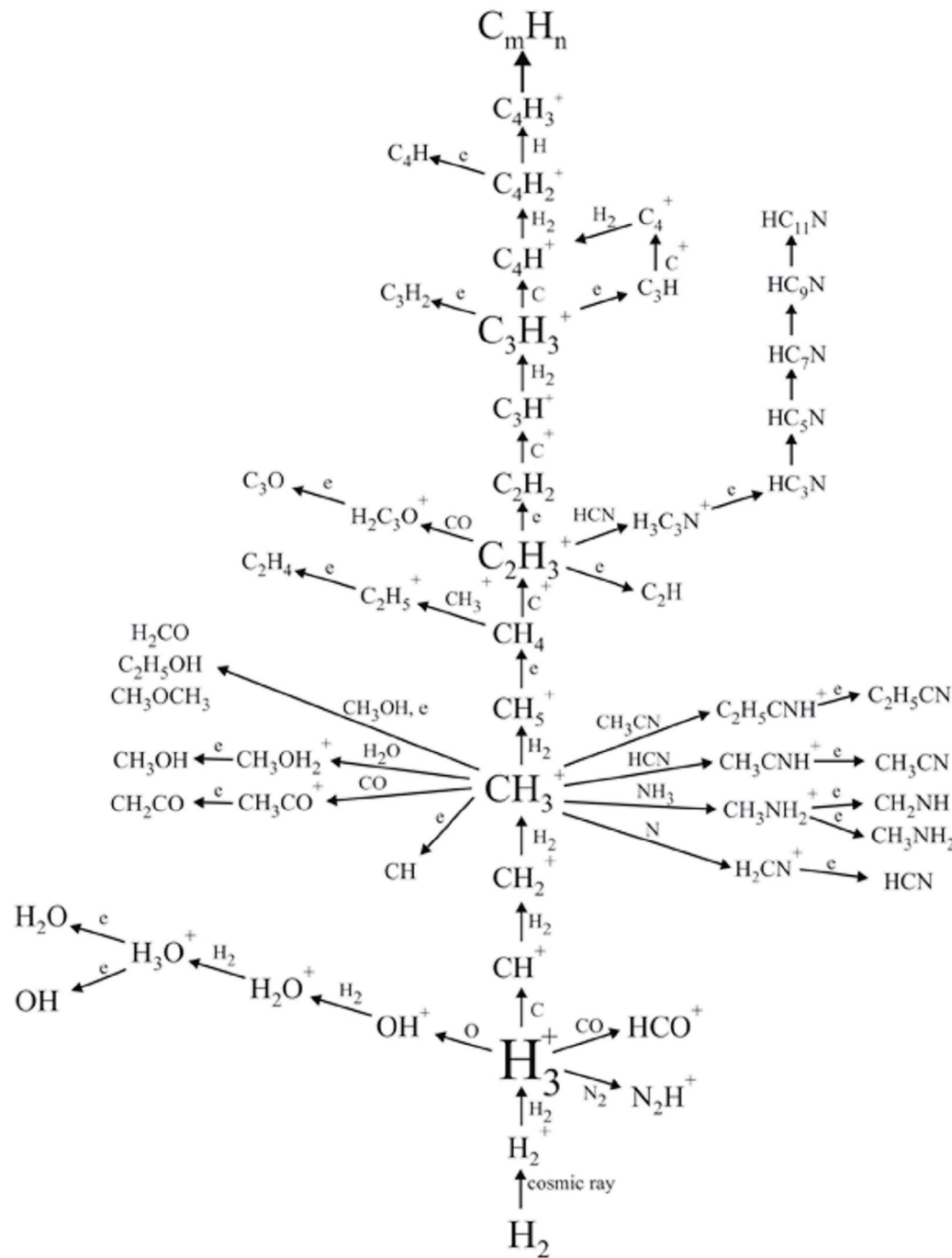
JUPITER SOUTHERN AURORAL ZONE



H_3^+ is extremely important to the chemistry of interstellar clouds. This is because H_3^+ willingly donates its extra proton to a variety of collision partners, thus laying the foundation for a large network of ion-molecule reactions.



Spectra of three interstellar clouds showing absorption lines due to the R(1,1)u and R(1,0) transitions of H_3^+ . Arrows indicate the expected positions of the absorption features from previous measurements of the interstellar gas velocity.



Why Molecular Ions?

As illustrated in the reaction chart to the left, molecular ions are responsible for driving the chemistry resulting in the formation of complex organic/prebiotic molecules. Detecting these ions in the ISM will lead to a better understanding of chemical evolution in space.

Why High-Resolution Spectroscopy?

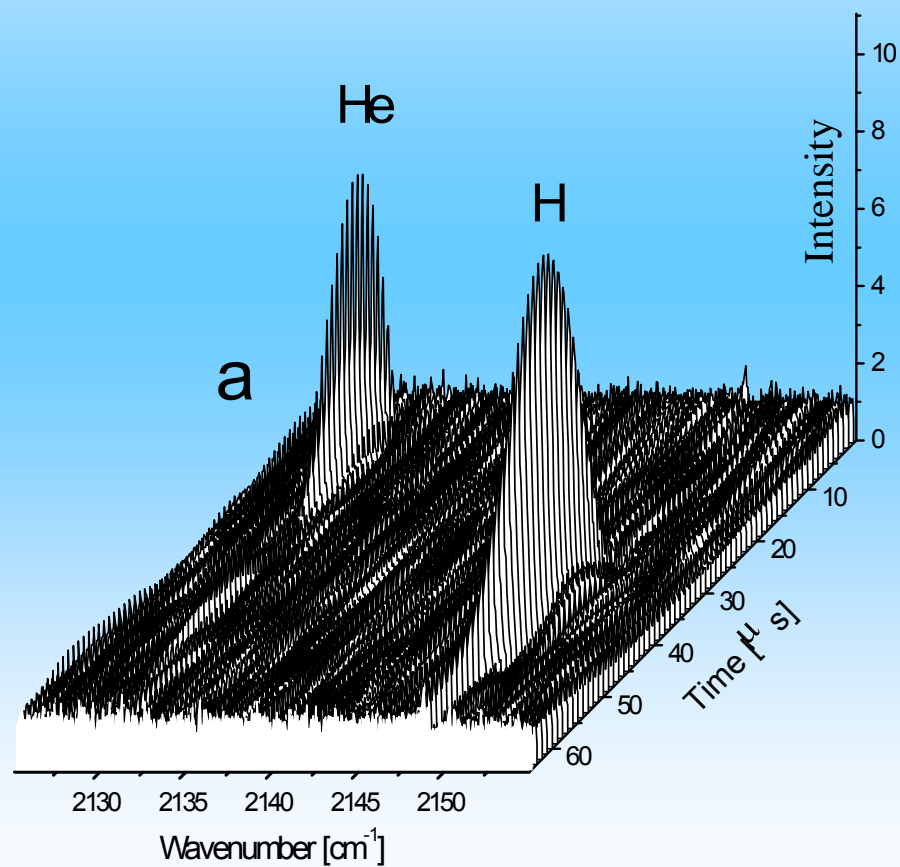
High-resolution spectra serve as molecular fingerprints, enabling astronomical searches for these species in the ISM, planetary atmospheres, and circumstellar environments.

Why Infrared Spectroscopy?

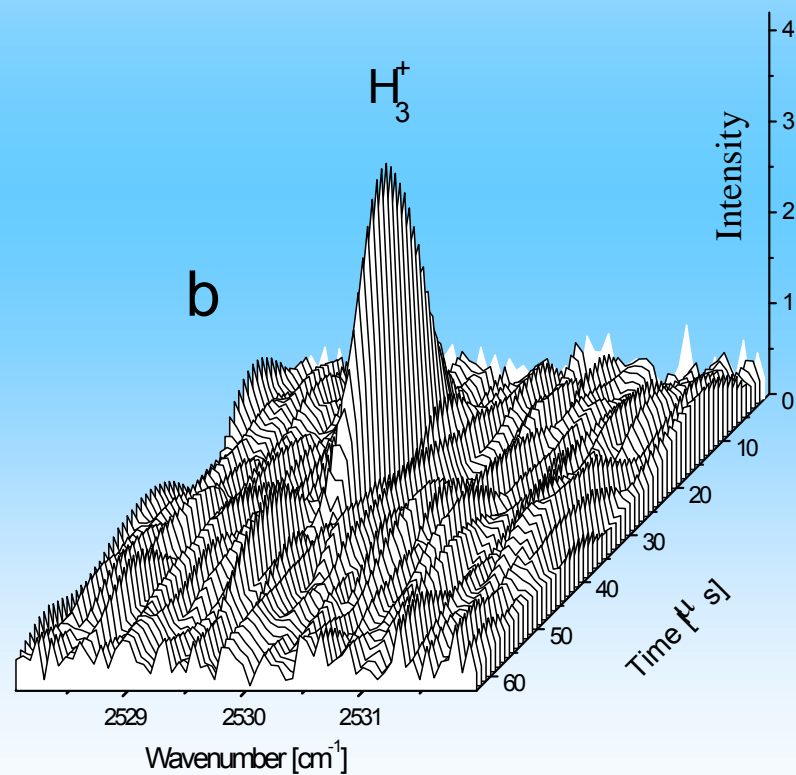
Infrared frequencies can be used to search for molecules without a permanent dipole (precluding microwave spectroscopy/radio astronomy), and can also be used to derive pure rotational transition frequencies of species which do contain a permanent dipole.

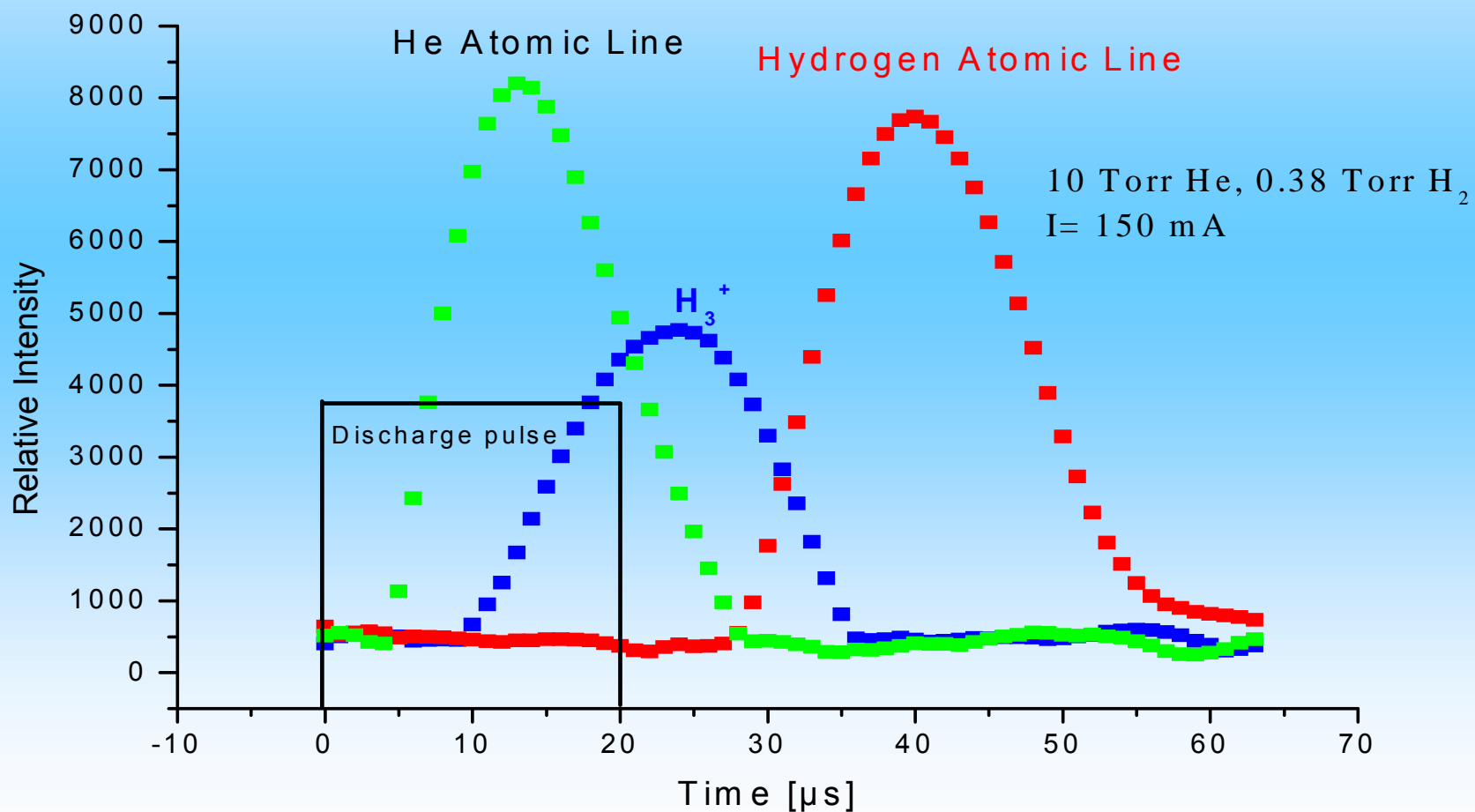
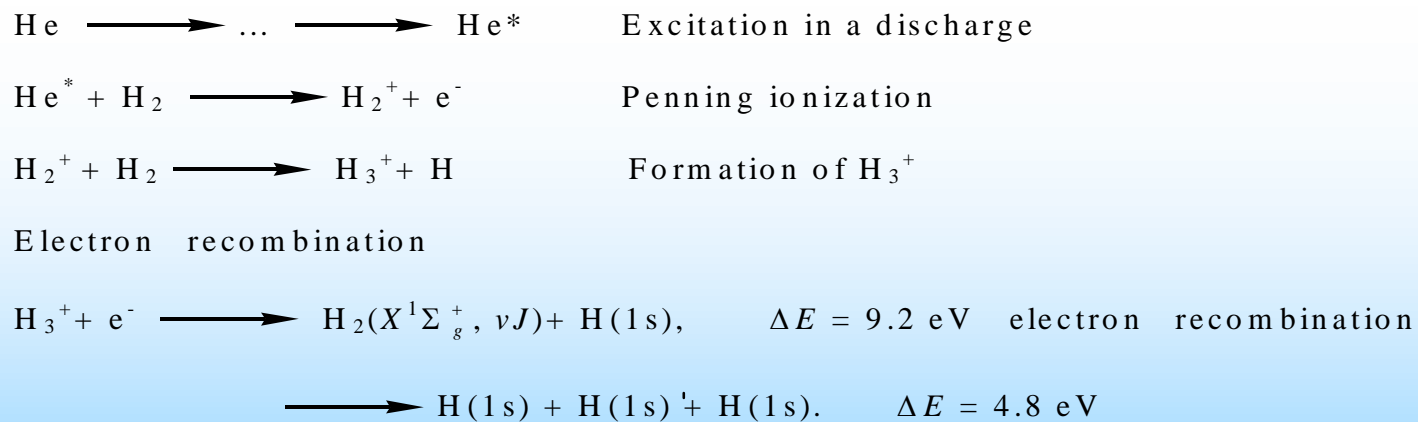
He + H₂ discharge plasma

He + H

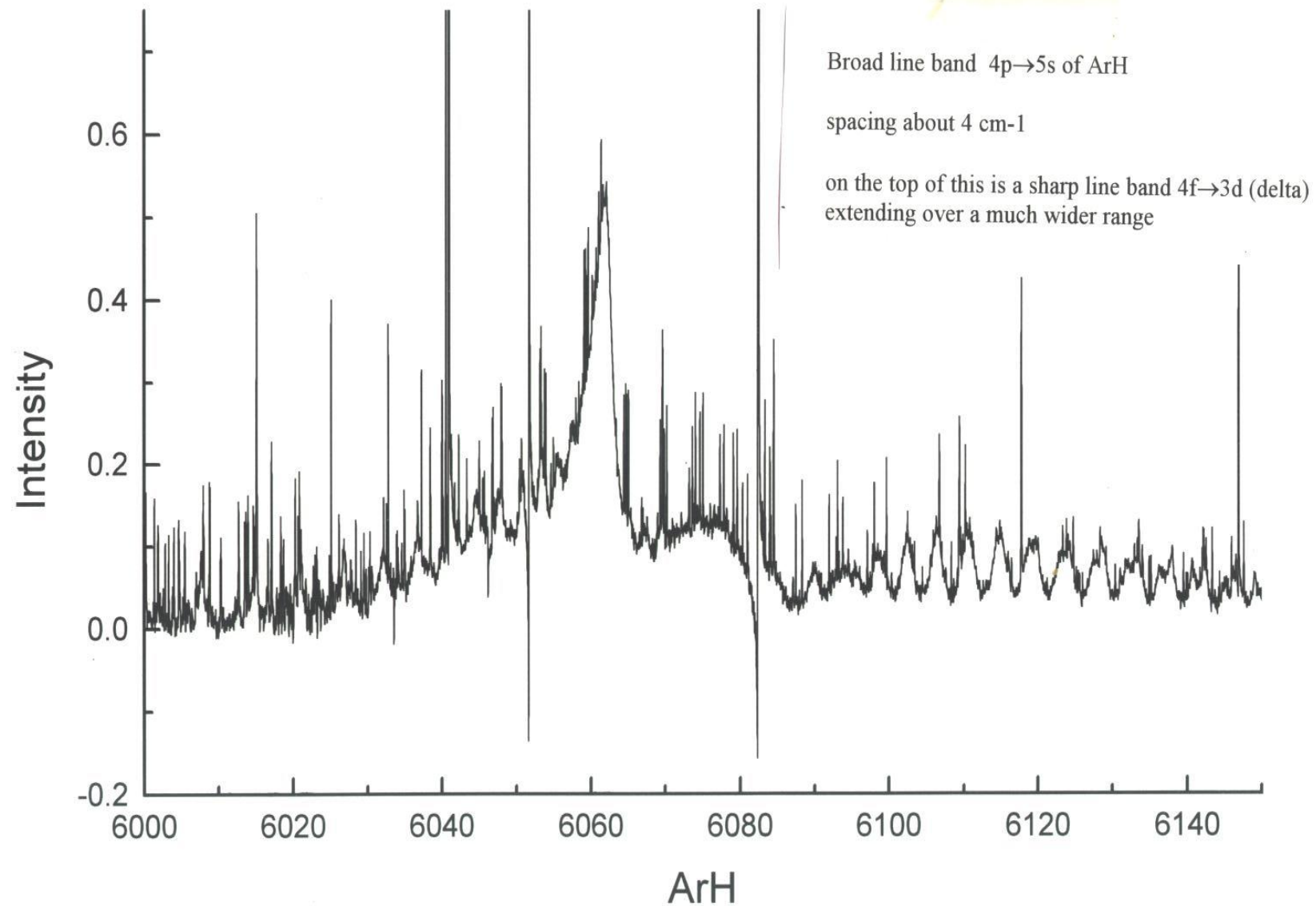


H₃⁺





Ar + H₂ discharge plasma



ArH radical

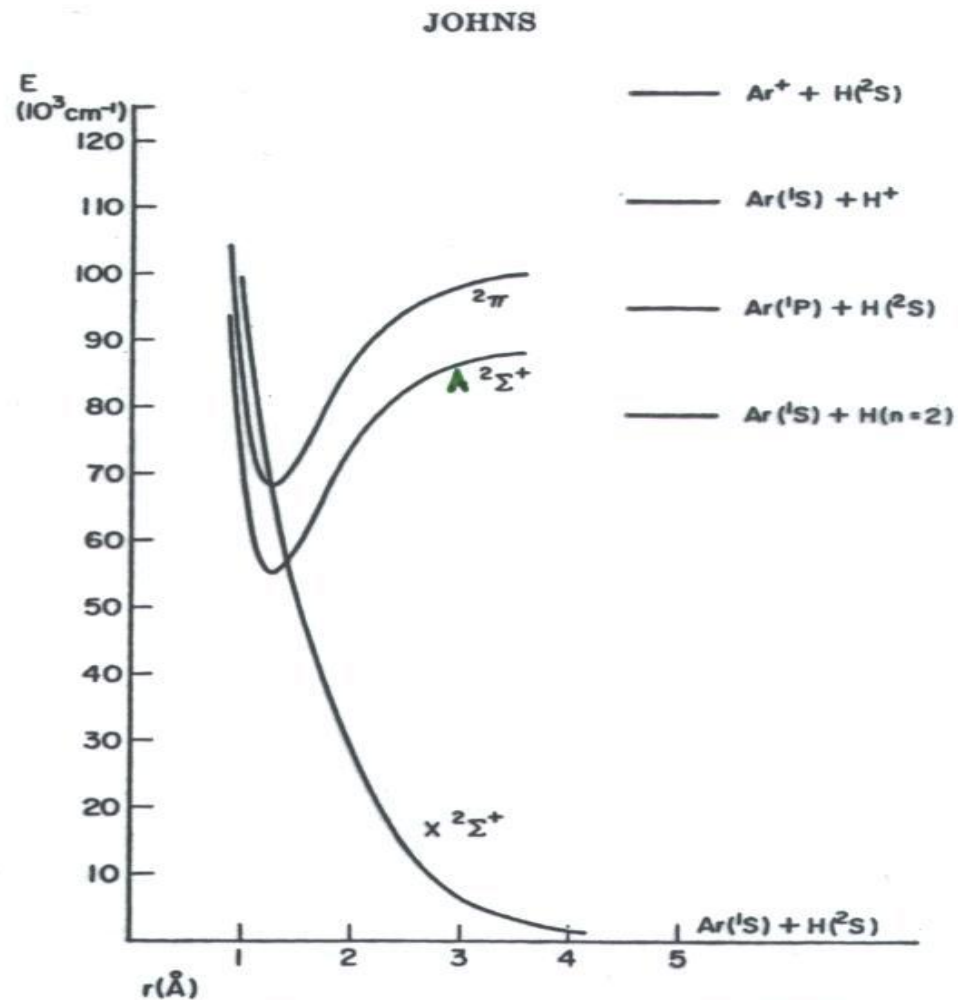
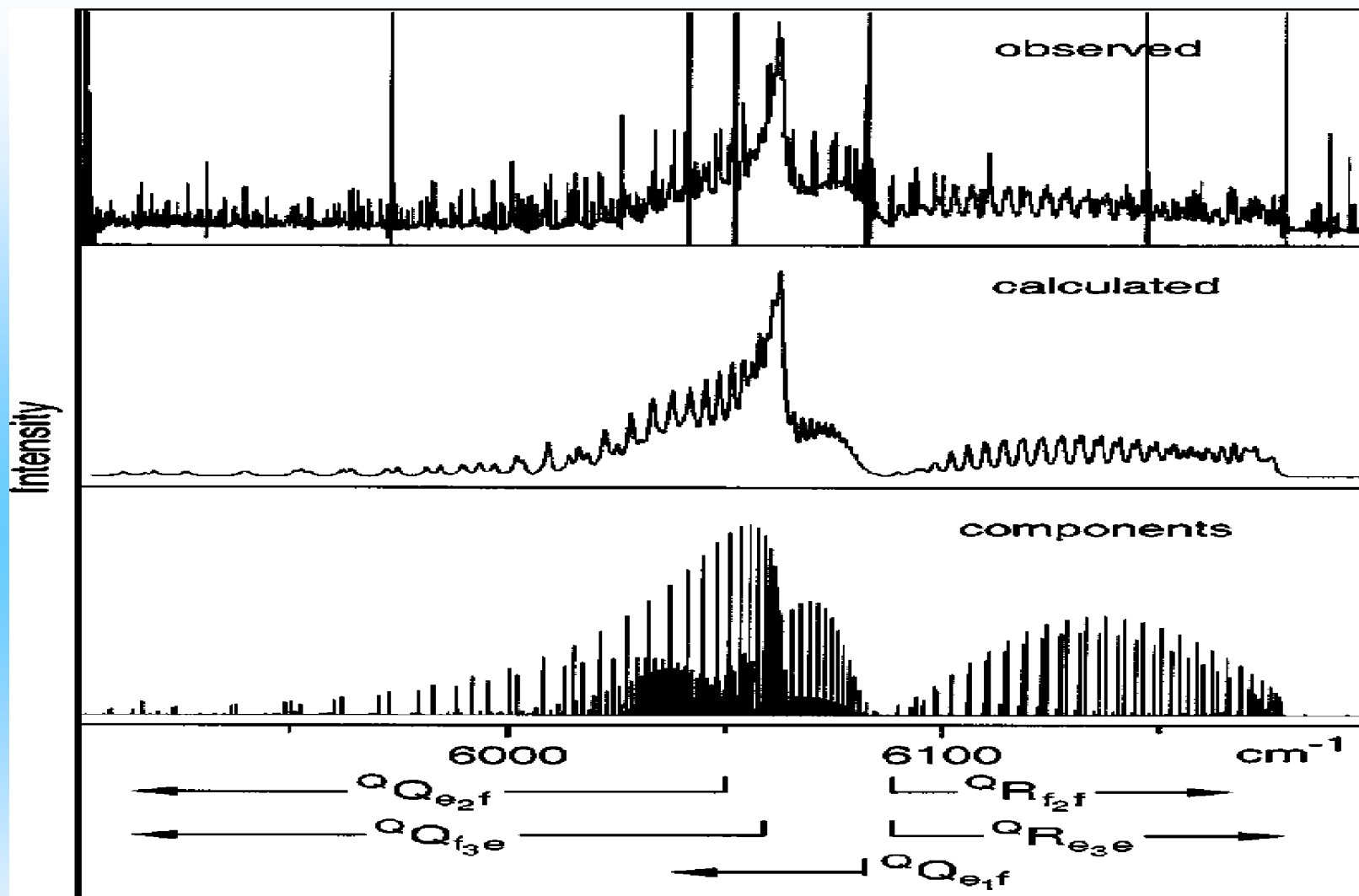


FIG. 5. Approximate potential curves for some states of ArH. The consequence of an avoided crossing between the A and X states has not been included in this figure.

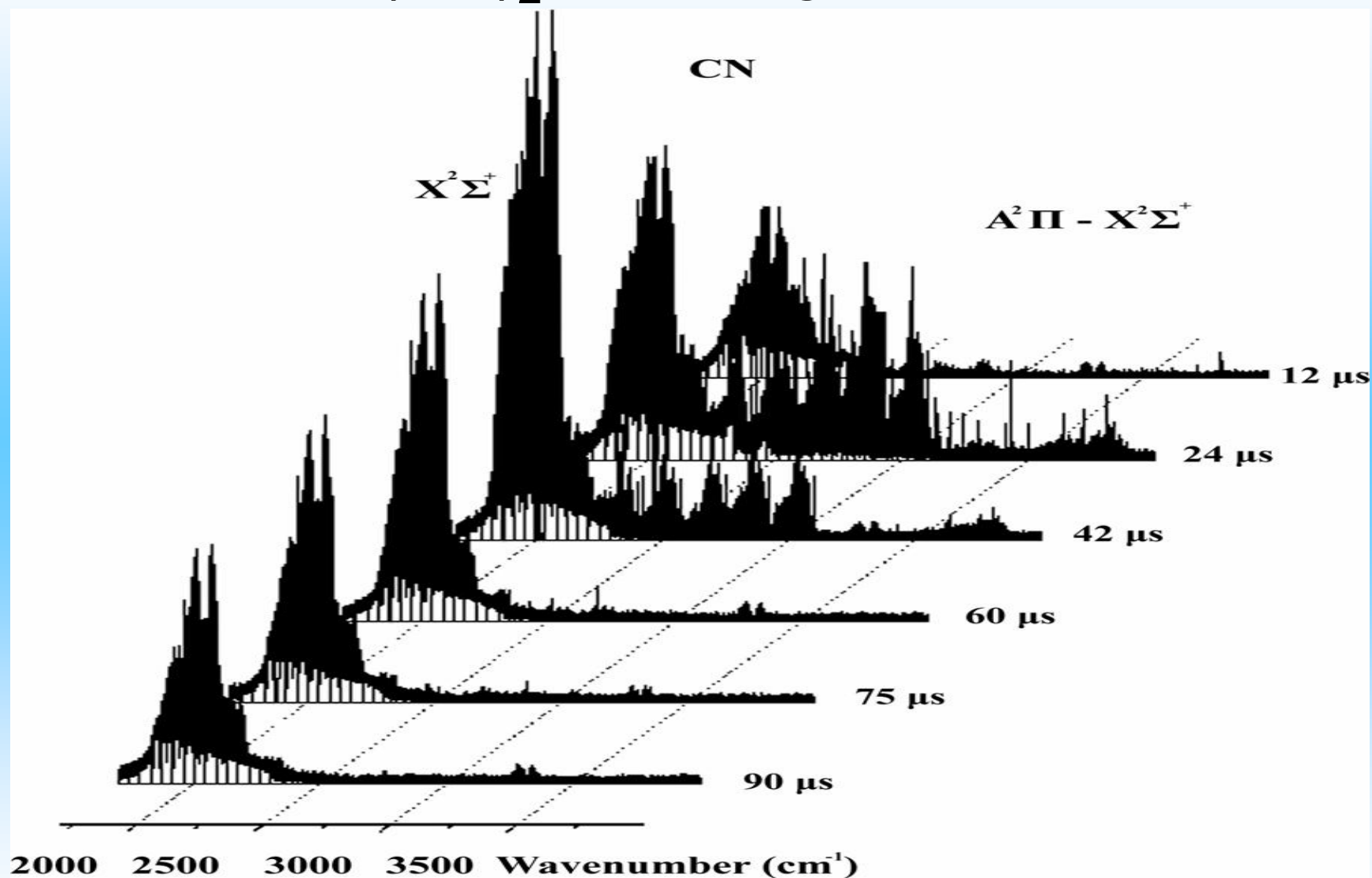


Emission spectrum of the 4p-5s band of ArH. Sharp lines belong to the 4f-3d band and probably to other unassigned bands of ArH. To calculate the spectrum a $T=2500$ K, and a halfwidth of 0.7 cm⁻¹ were assumed.

Investigated bands of ArH and ArD

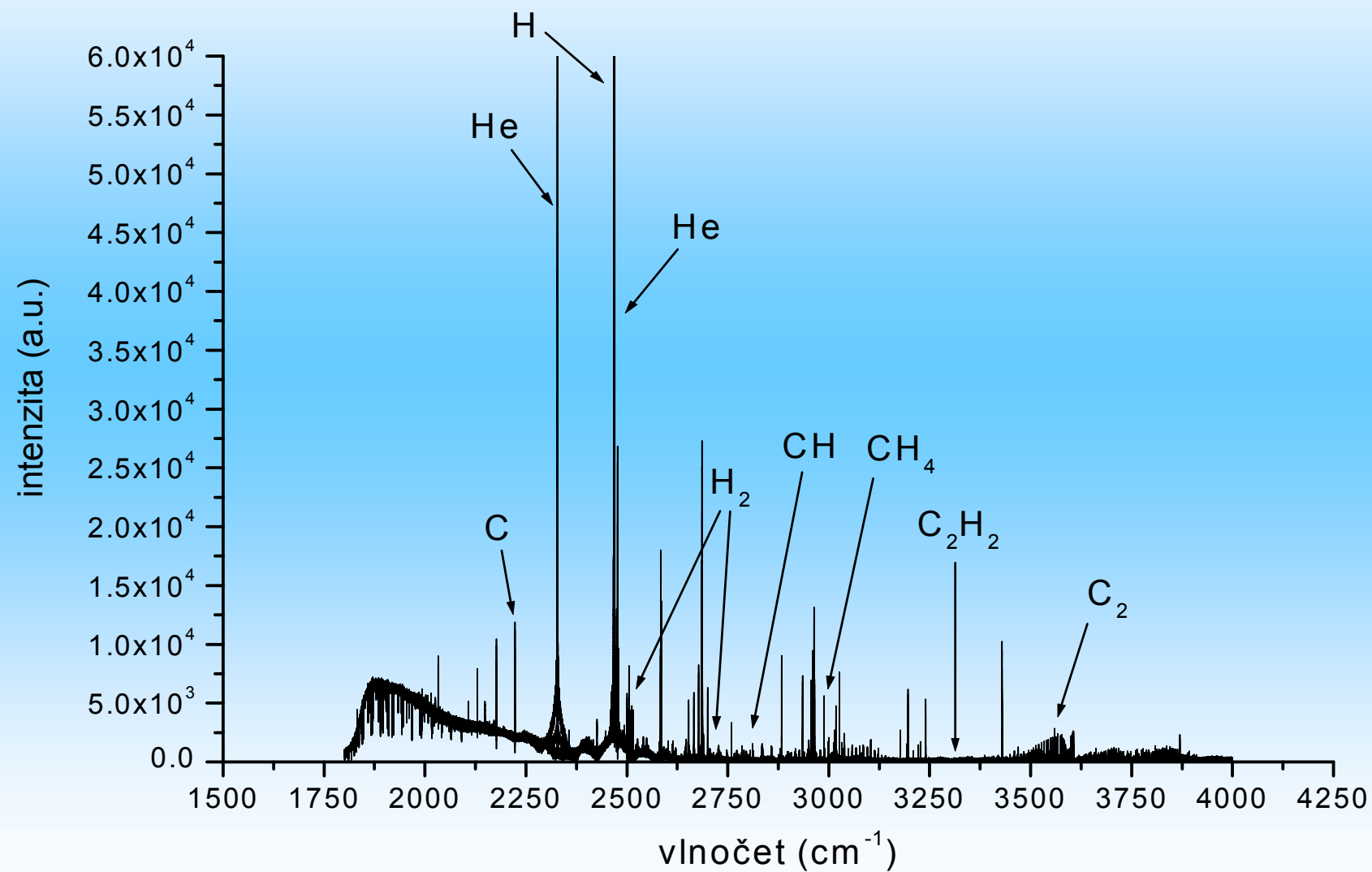
Transition	ArH	ArD	Resolution cm ⁻¹	Reference
	Band Origin cm ⁻¹			
3d π \rightarrow 5s	13024	13040	0.1	2
4d π \rightarrow 5s	19520	19570		3
π \rightarrow 5s ?	-	22575		3
5p \rightarrow 5s	-	17487	0.13	5
5p \rightarrow 6s	-	3681	0.20	5
6p \rightarrow 5s	-	21676	0.20	5
4p \rightarrow 5s	-	6120	0.05; 0.2	4
3d σ \rightarrow 4p	-	10200	0.05; 0.2	4
3d π \rightarrow 4p	6934	6900	0.05; 0.2	4,6
3d δ \rightarrow 4p	8534	8524	0.05	6
4d σ \rightarrow 4p	15097	15075	0.3	6
6s \rightarrow 4p	7703	7685	0.02	6
8s \rightarrow 4p	-	16749	0.20	6
4f \rightarrow 5s	20641	20682	0.20	7
4f \rightarrow 3d π	7627	7649	0.05	7
4f \rightarrow 3d δ	6027	6038	0.05	7
4f \rightarrow 3d σ	~ 4400	4351	0.10; 0.20	7,8

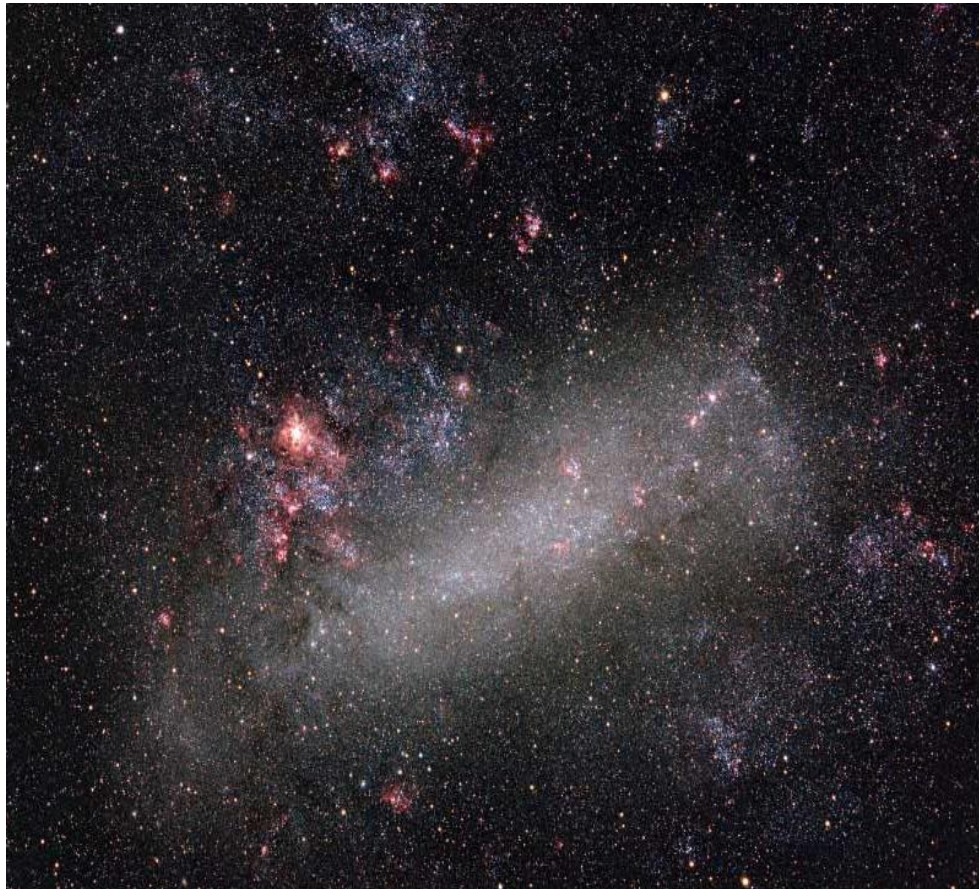
Ar + (CN)₂ discharge plasma



The time-resolved emission FT spectrum from a pulsed discharge in a (CN)₂ and He mixture. The discharge pulse duration was 20 ns. The 30 time-resolved spectra were collected from t = 0–90 μs with a step of 3 μs. The spectra of C₂H₂ and C₂ were observed at 3300 and 3600 cm⁻¹.

He + CH₄ discharge plasma





CH

*Astronomy
of the South*

Vis lines at 3878.8, 3886.4, 3890.2,
4300.3, UV lines near 1271, 1368, 1369,
1370, 1549, & 1694

$J = 1/2, \nu_{01}$ (3263.794 MHz)

$J = 1/2, \nu_{11}$ (3335.481 MHz)

$J = 1/2, \nu_{10}$ (3349.193 MHz)

Large Magellanic Cloud

[Considerations Regarding Interstellar Molecules](#)

P. Swings and L. Rosenfeld, ApJ 86:483-486 (1937)

[Evidence for the Molecular Origin of Some Hitherto Unidentified Interstellar Lines](#)

A. McKellar, Publ Astron Soc Pac 52:187-192 (1940)

[Some Results with the COUDÉ Spectrograph of the Mount Wilson Observatory](#)

W. S. Adams, ApJ 93:11-23 (1941)

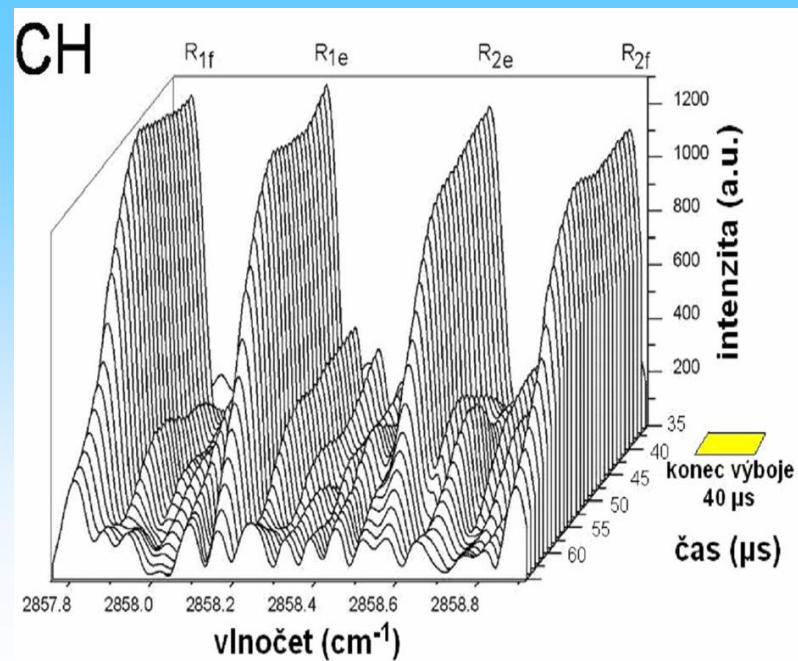
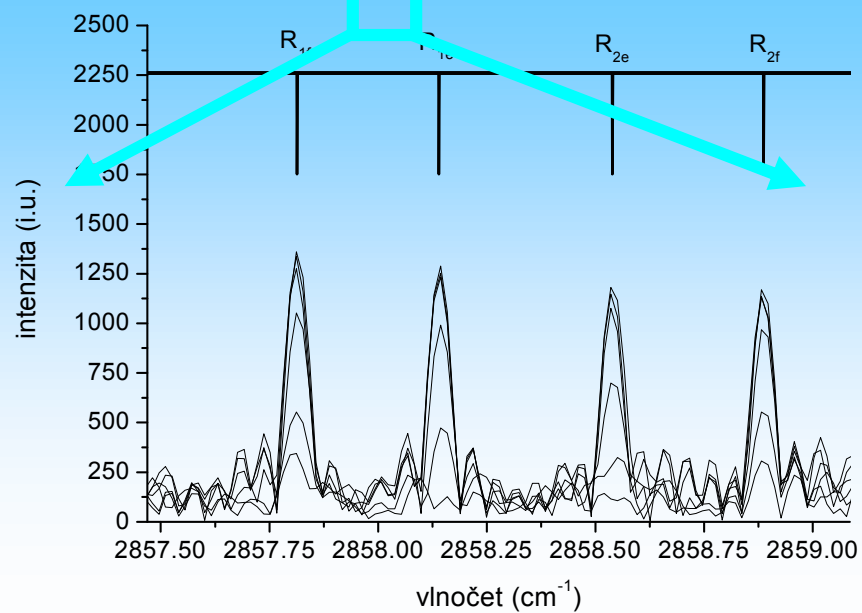
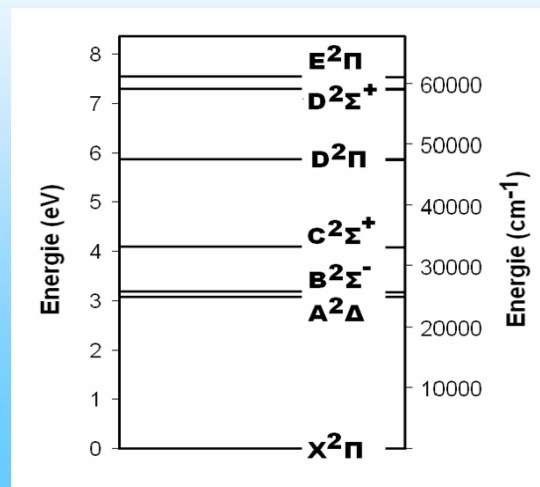
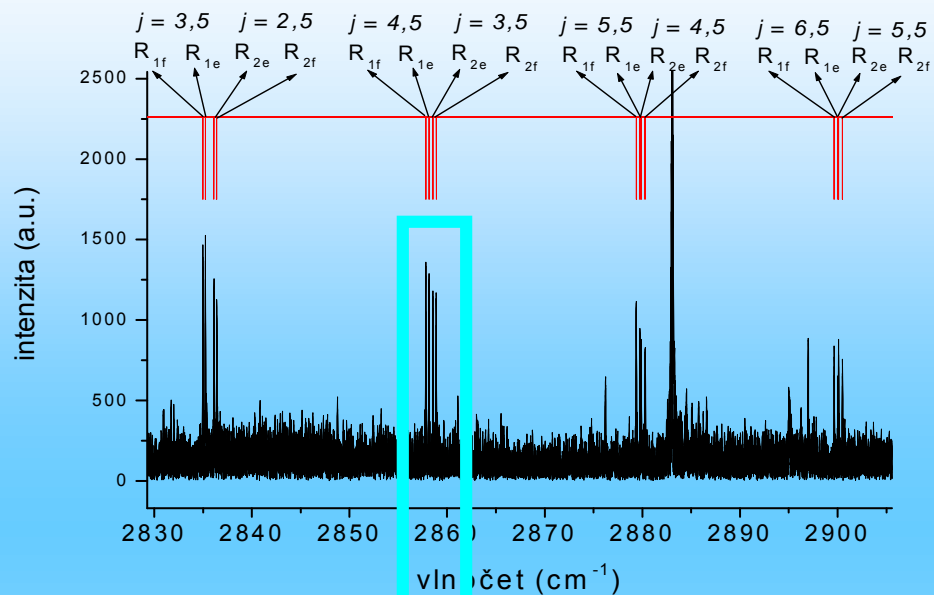
[Radio Detection of Interstellar CH](#)

O. E. H. Rydbeck, J. Eildér, and W. M. Irvine, Nature 246:246-248 (1973)

[Hubble Space Telescope Measurements of Vacuum Ultraviolet Lines of Interstellar CH](#)

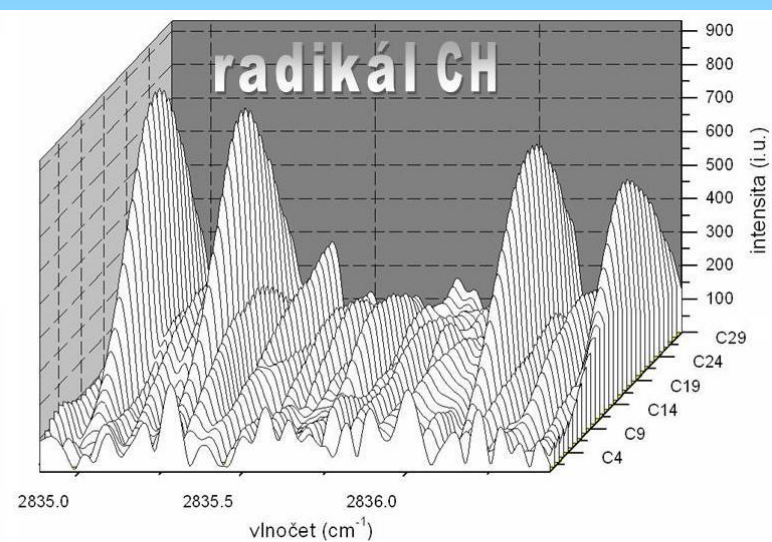
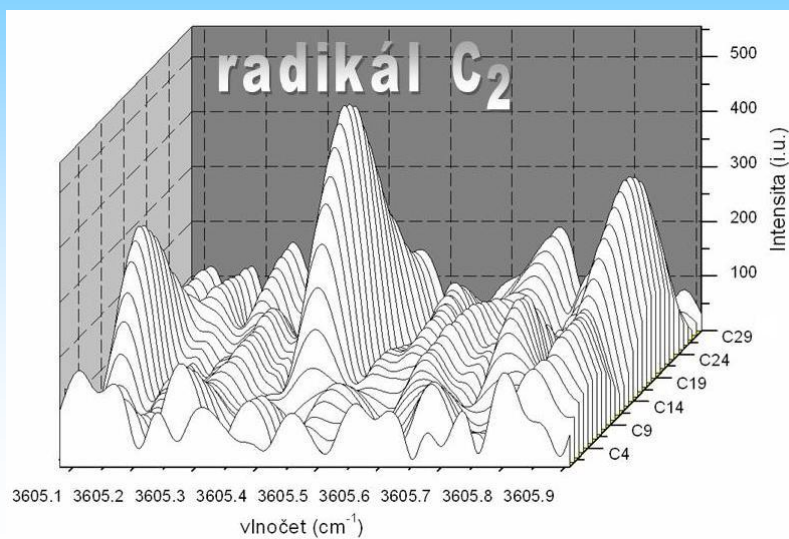
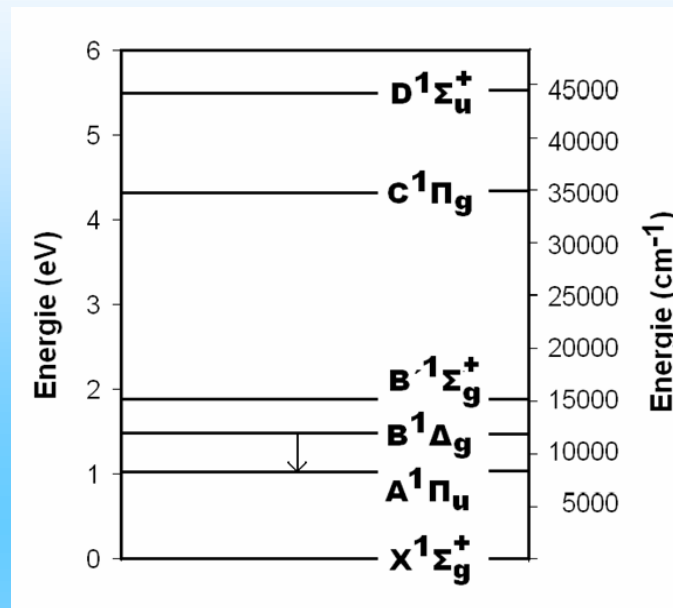
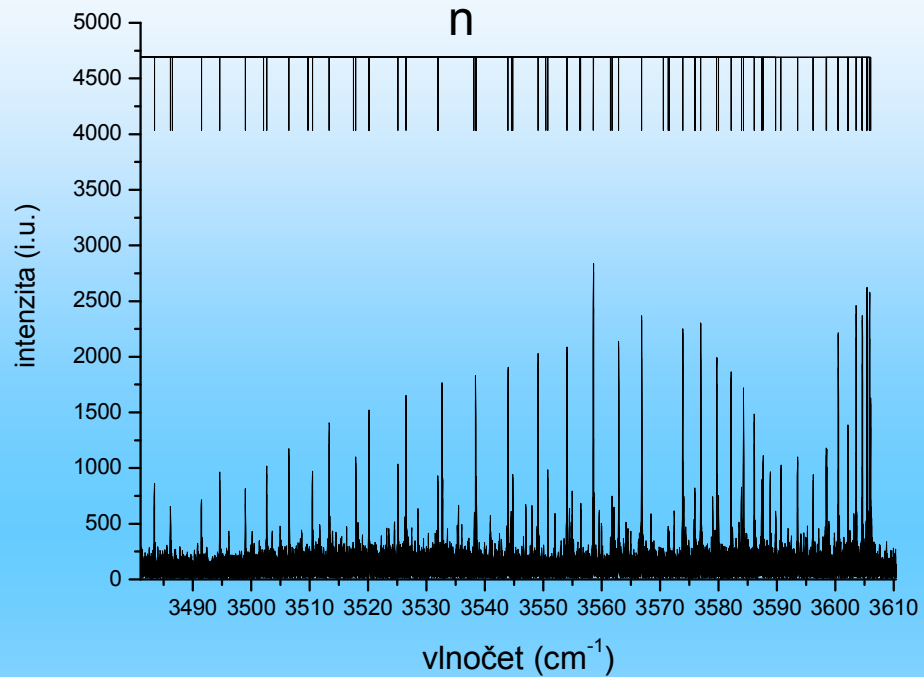
Y. Sheffer and S. R. Federman, ApJ 659:1353-1359 (2007)

CH radical

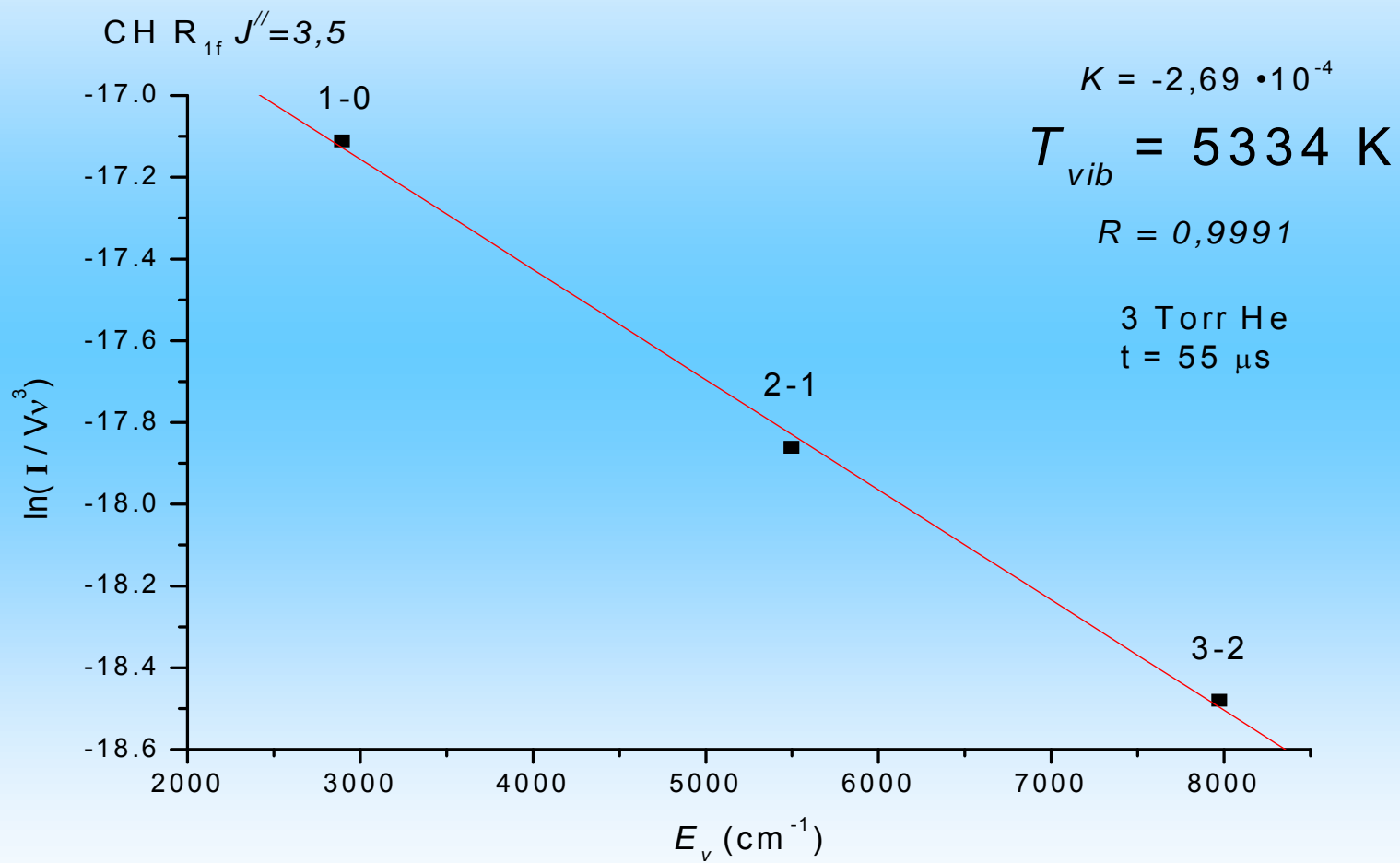


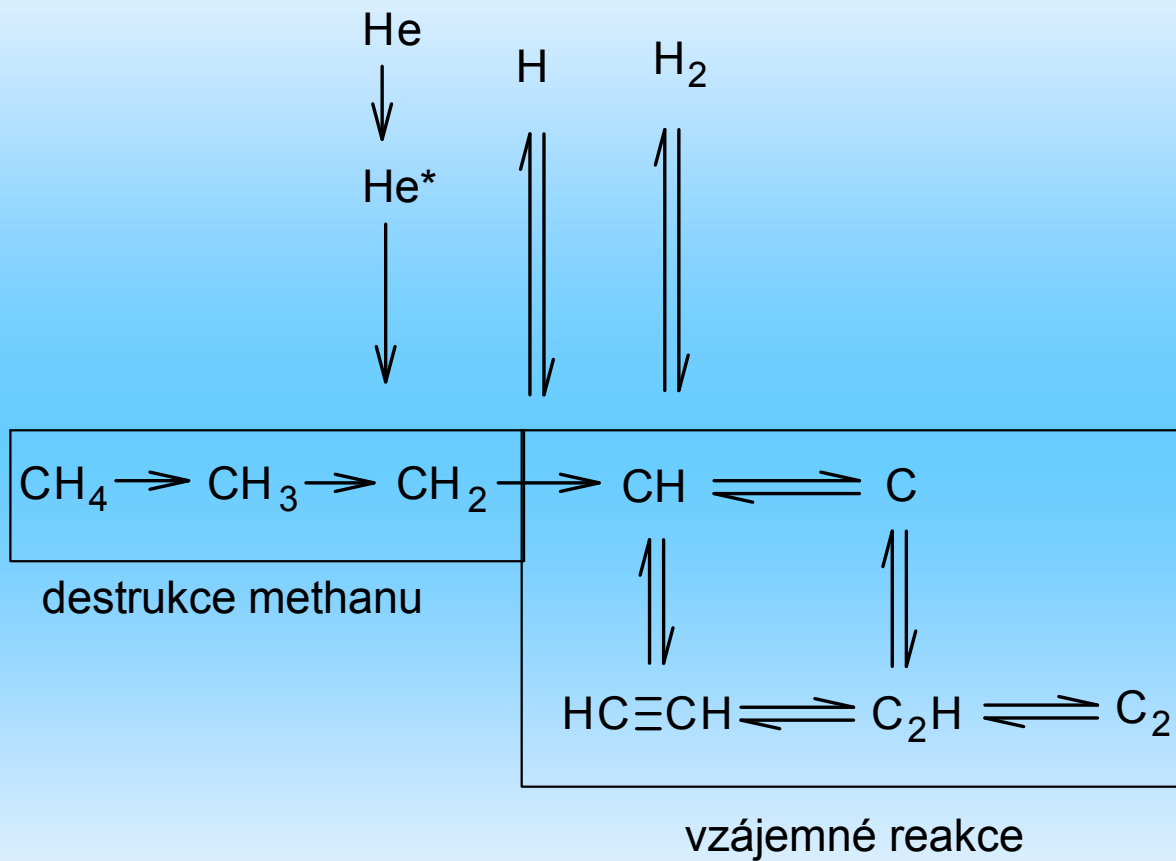
C₂ radical

n



CH vibrational temperature



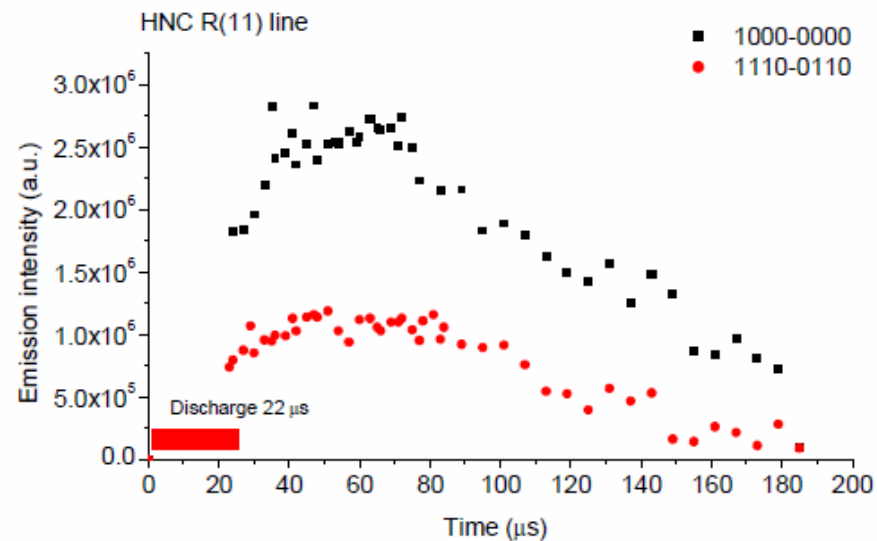
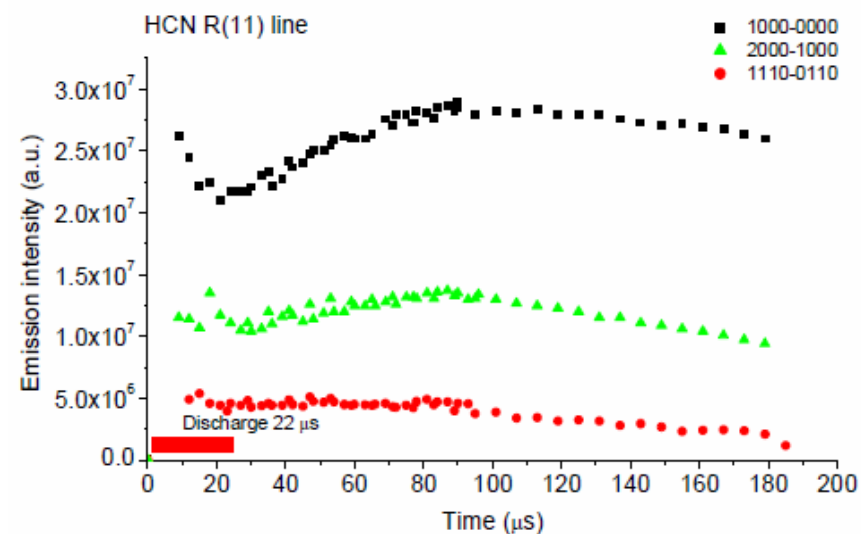
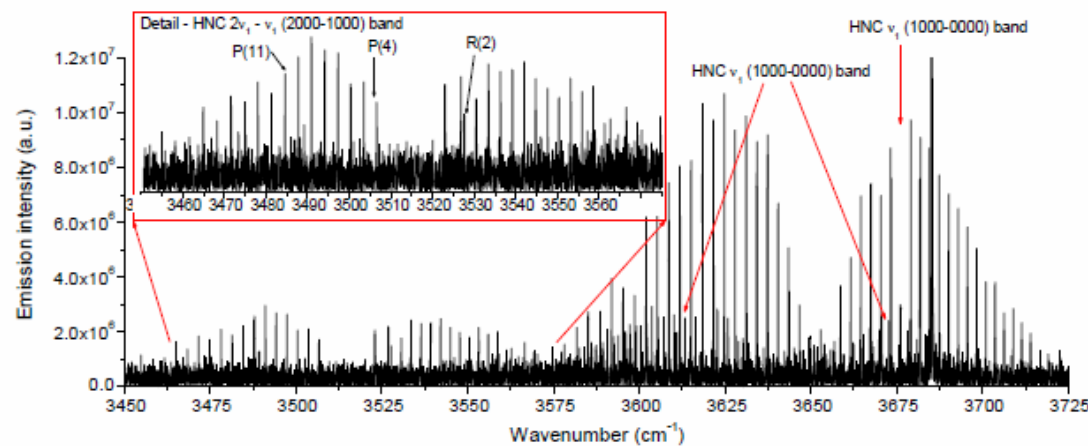
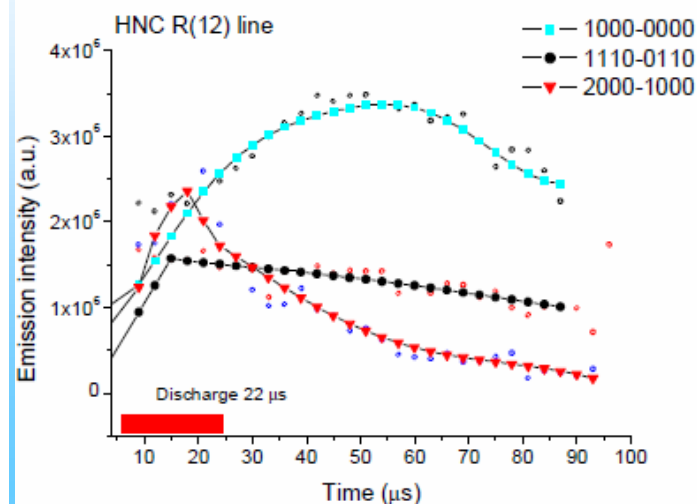


Vibrational temperature 5000 ó 6000 K

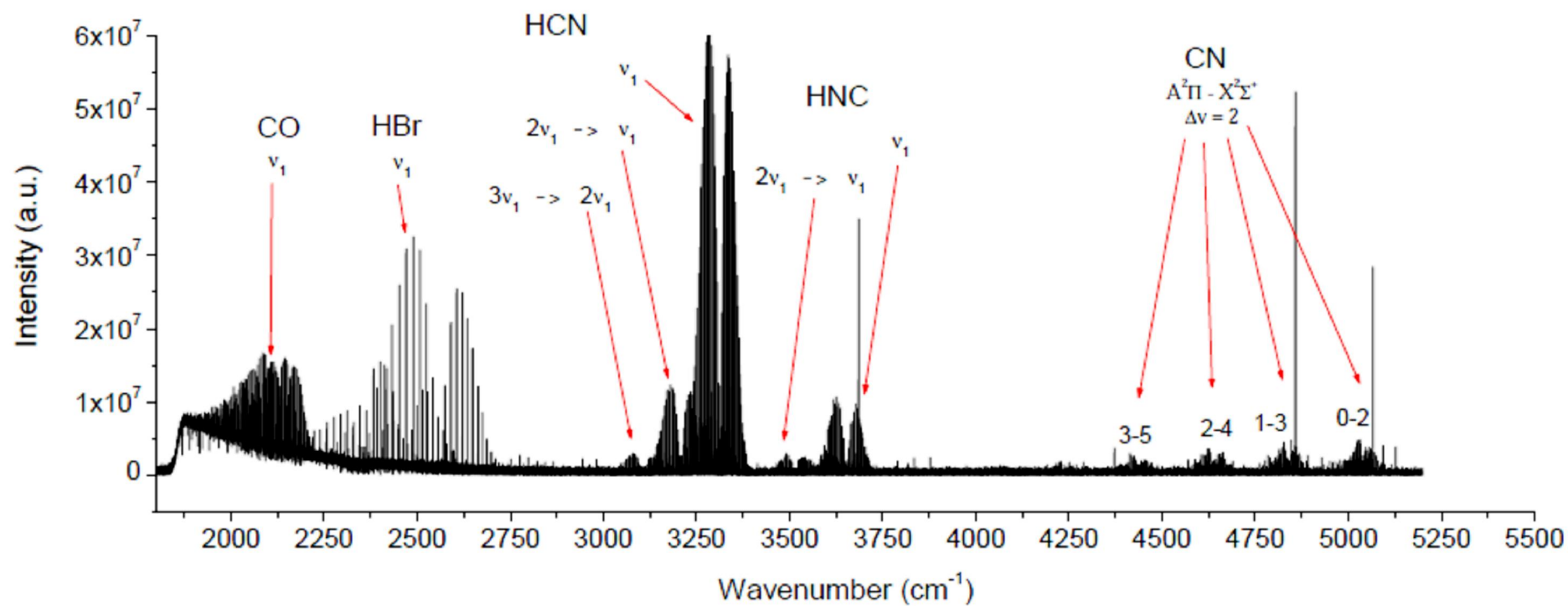
Rotational temperature 550 ó 700 K

Excitation temperature 2500 ó 3000 K

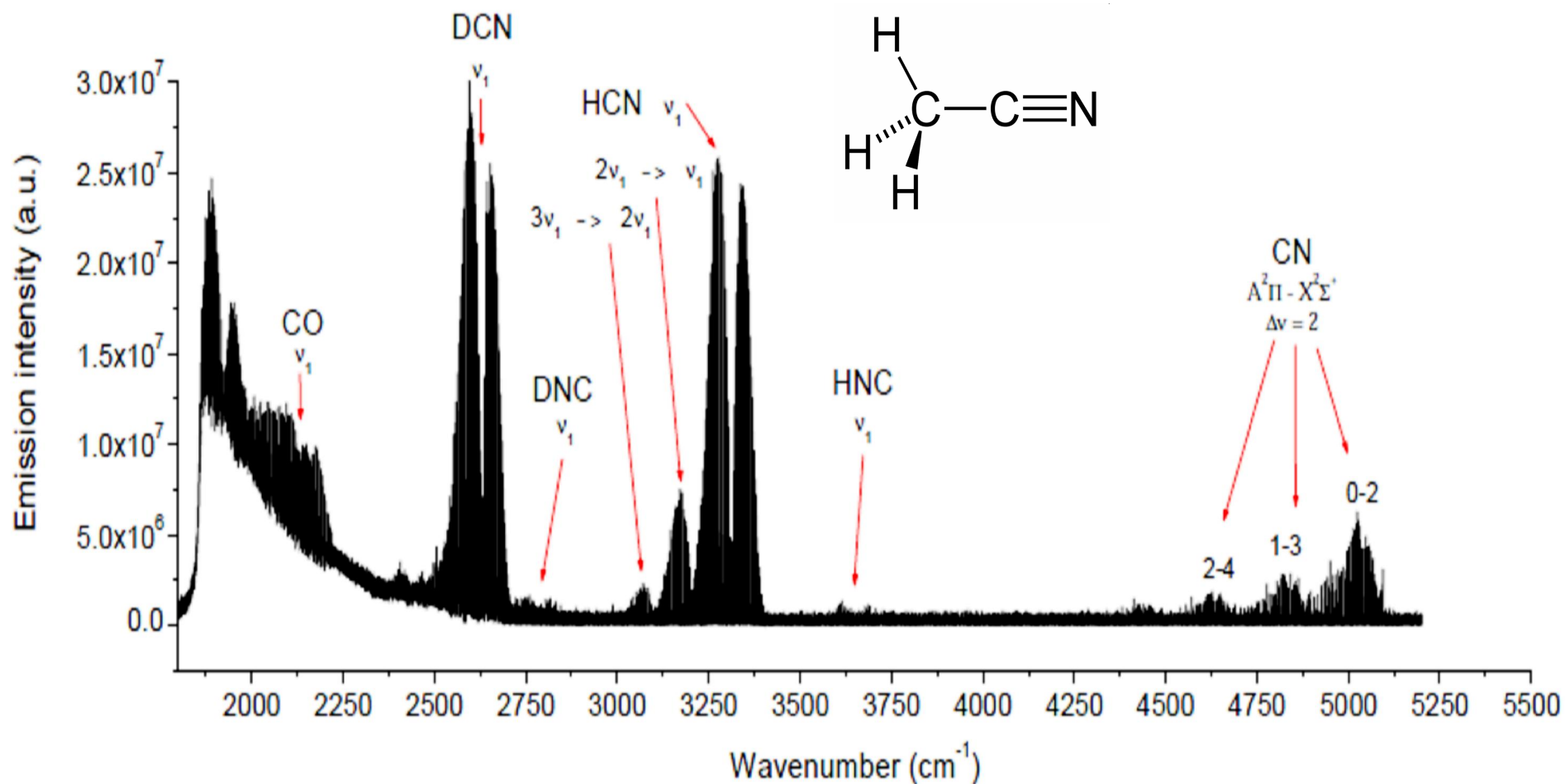
HNC/HCN emission



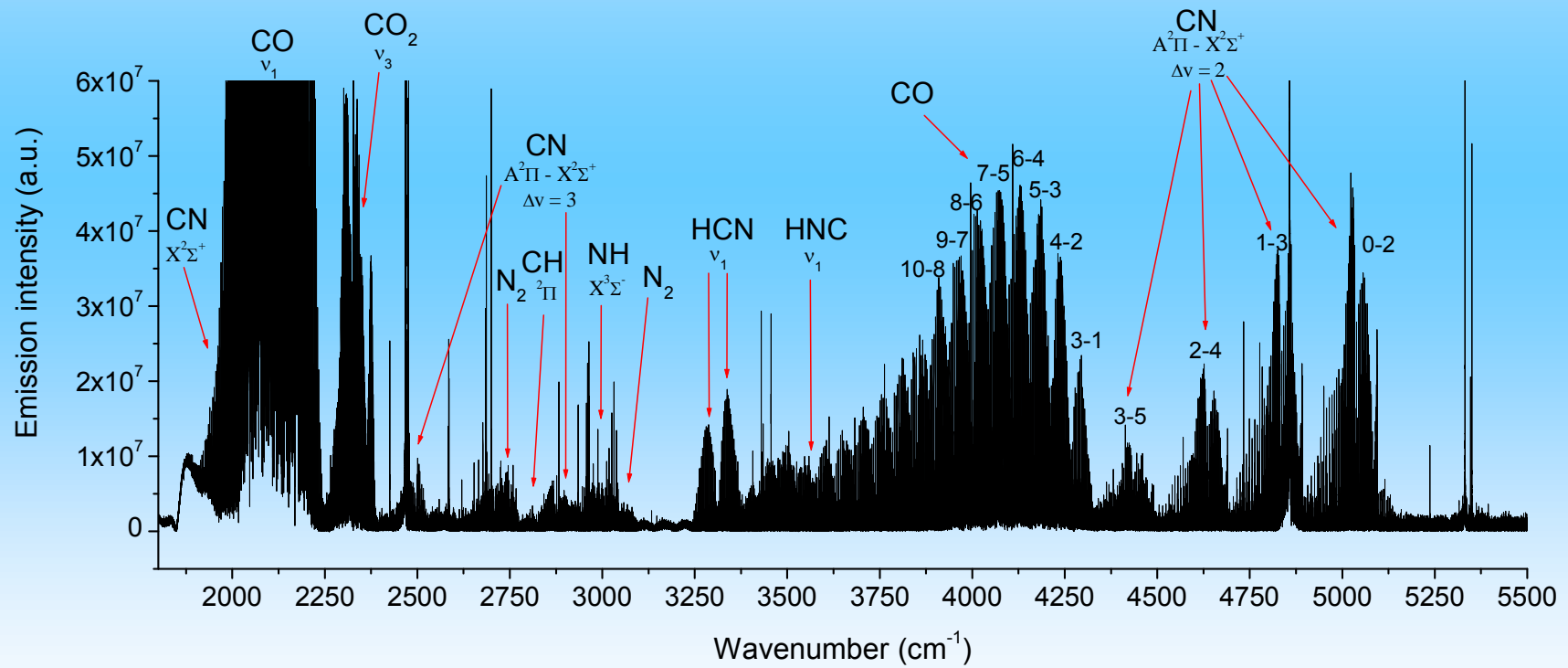
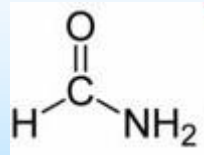
BrCN discharge



Acetonitrile discharge



Formamide discharge

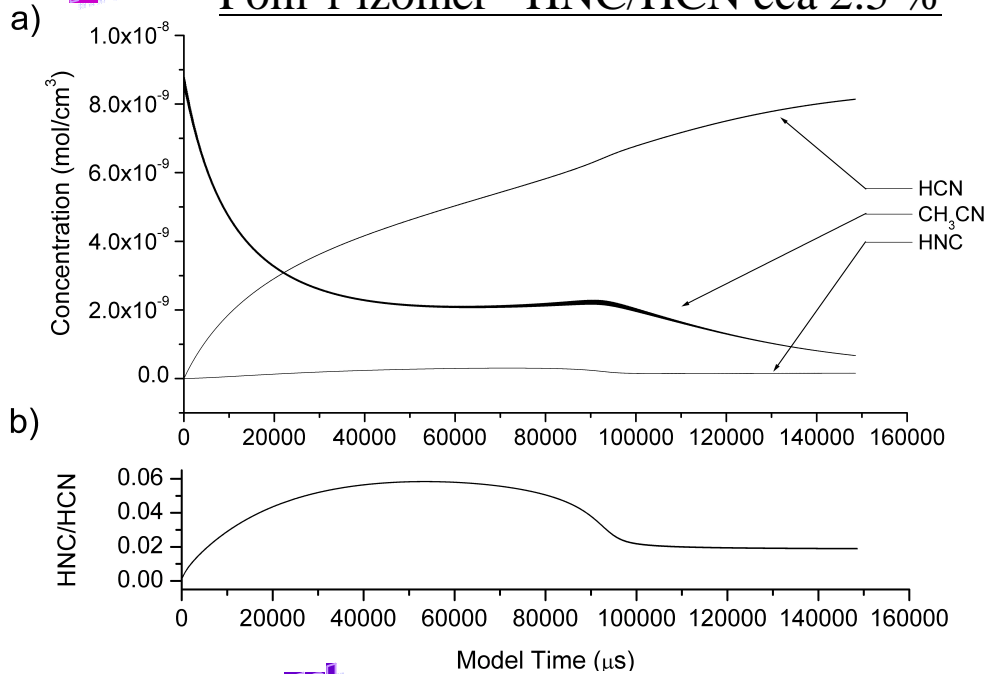


Kinetický model výboje

Model

Kumulace produktů v cele

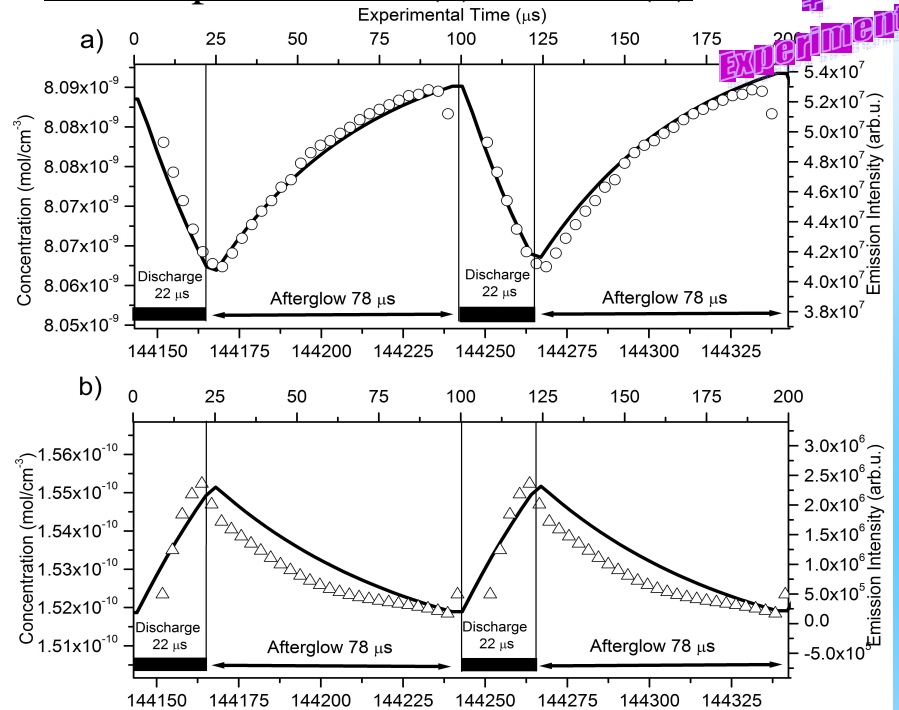
Poměr izomer HNC/HCN cca 2.5 %



Emisní profil HCN (a) a HNC (b)

Model

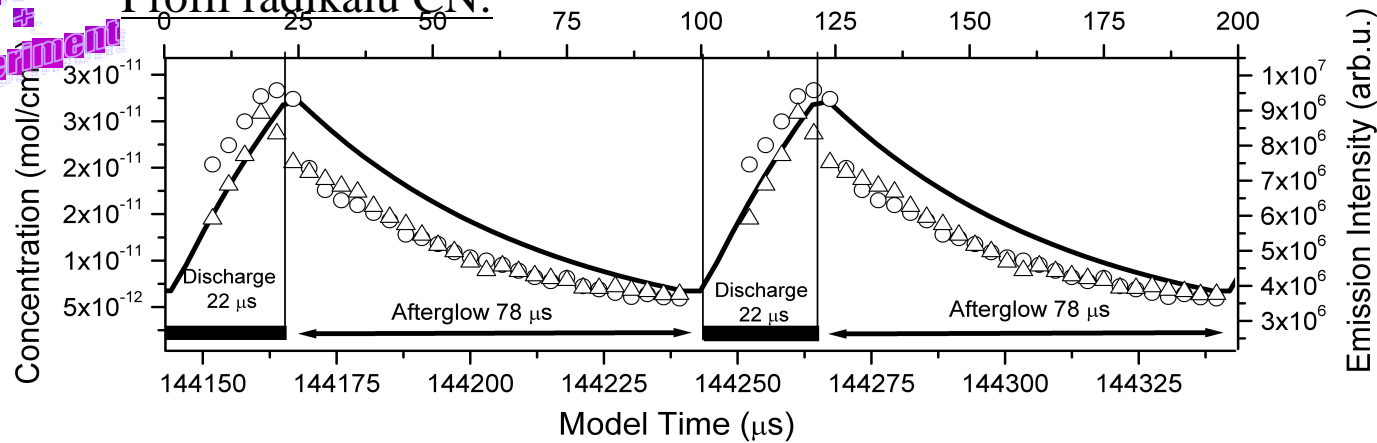
Experiment



Model
Experiment

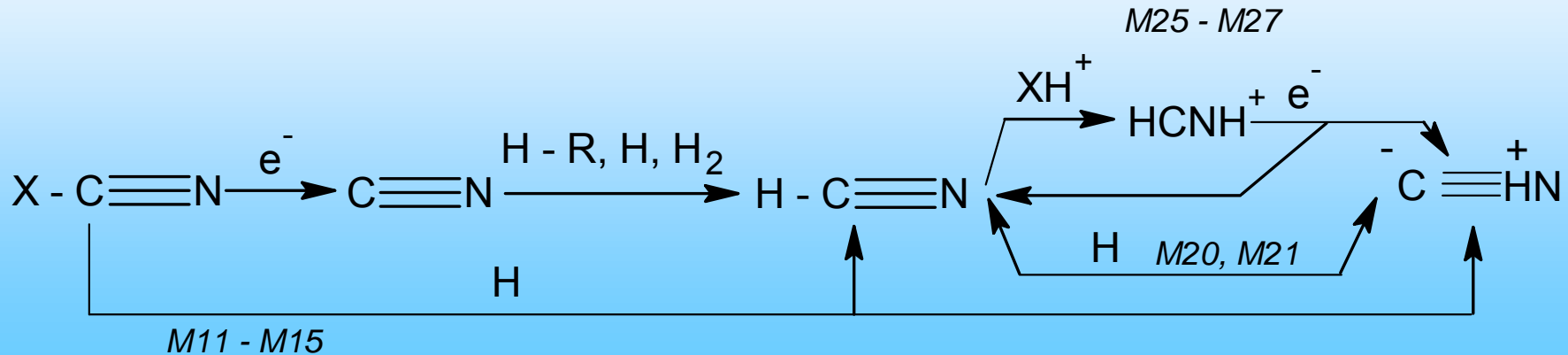
Profil radikálu CN:

Experimental Time (μs)



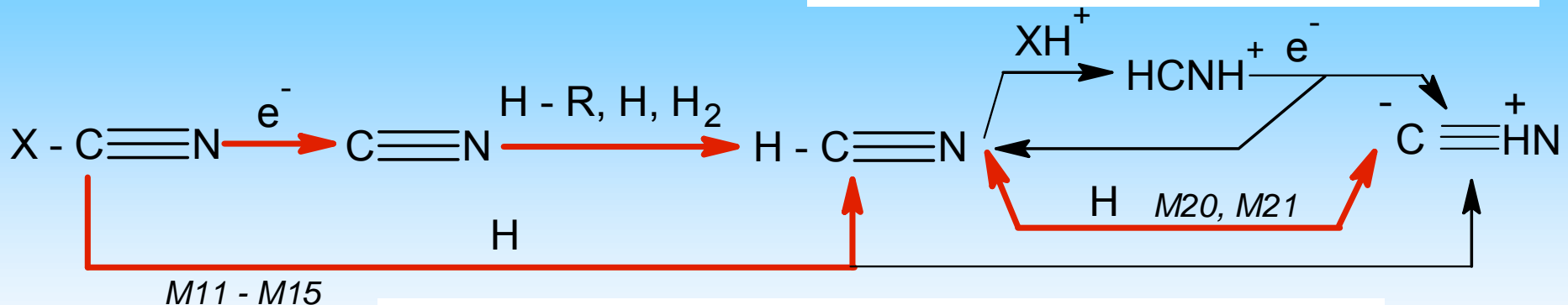
Kinetický model výboje

Mofné cesty vzniku HNC ve výboji:



Výsledek modelu vzniku HNC ve výboji:

Ionty - rychlé reakce, ale malé koncentrace

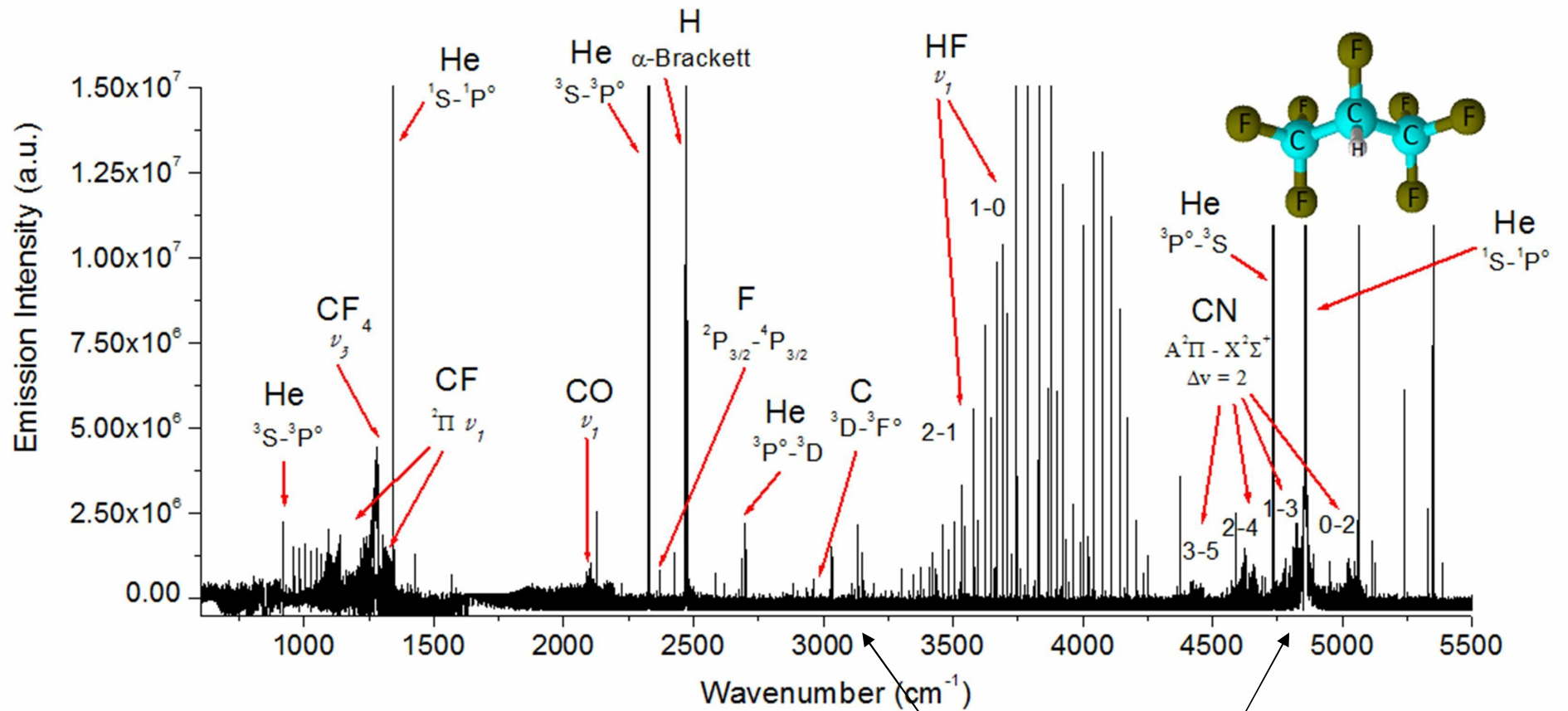


Radikály - kumulace HCN ve výbojové cele

b hem milisekund dáno rovnováhou $HNC + H \leftrightarrow HCN + H$

Potla ením radikálových reakcí vzroste pom r HNC/HCN.

Výboj - hasivo FM 200

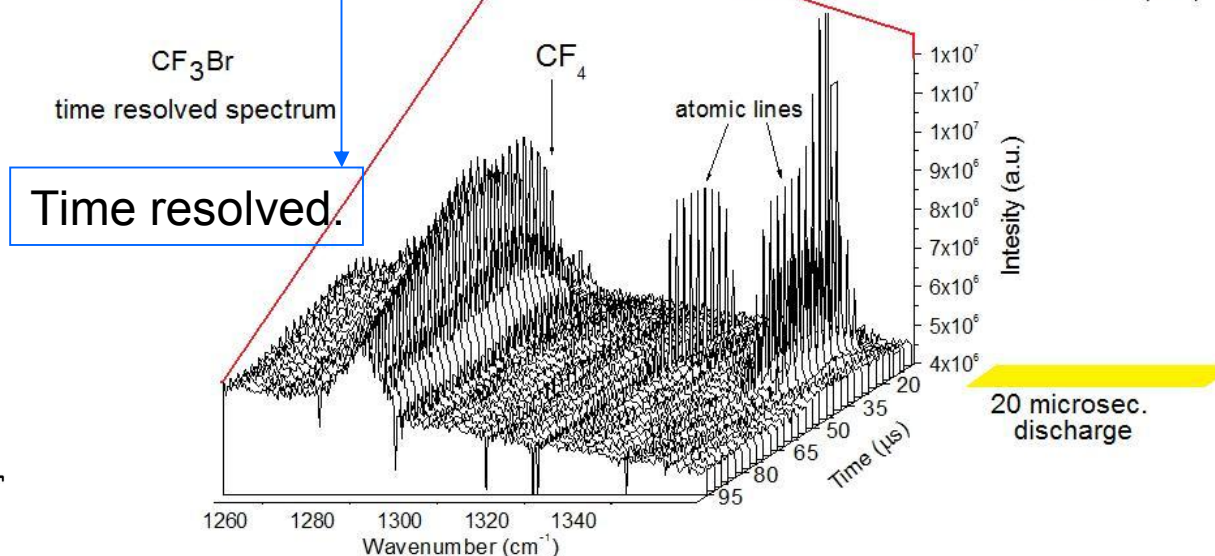
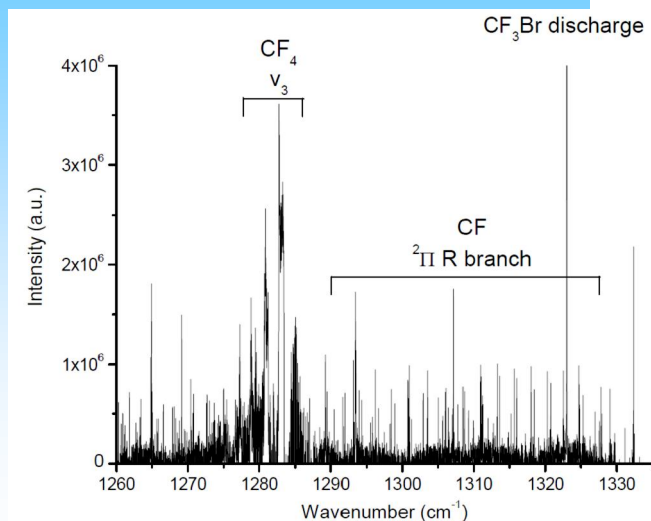
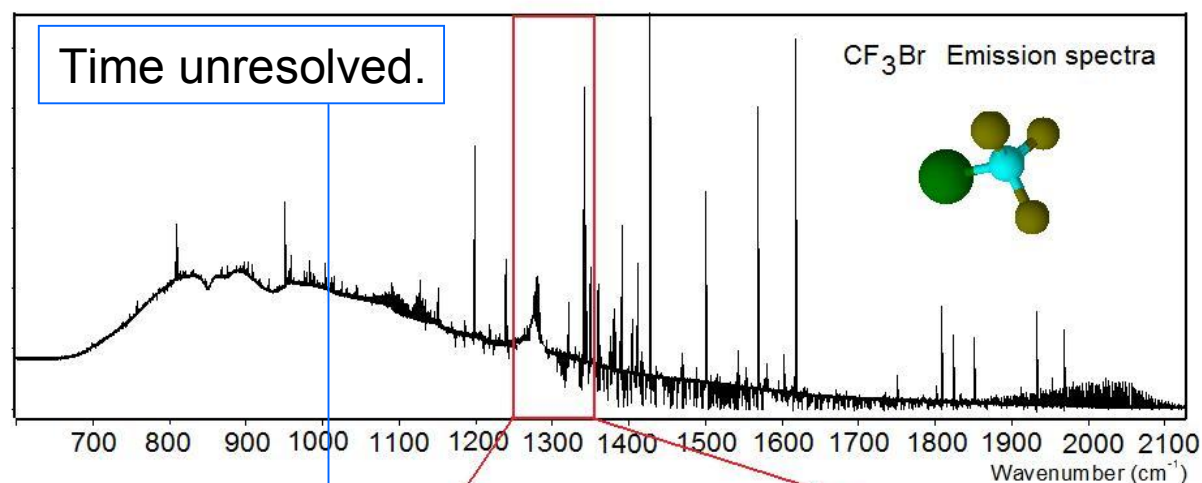


Ve výboji CN
P ídavek vody - HCN (slabý pás).

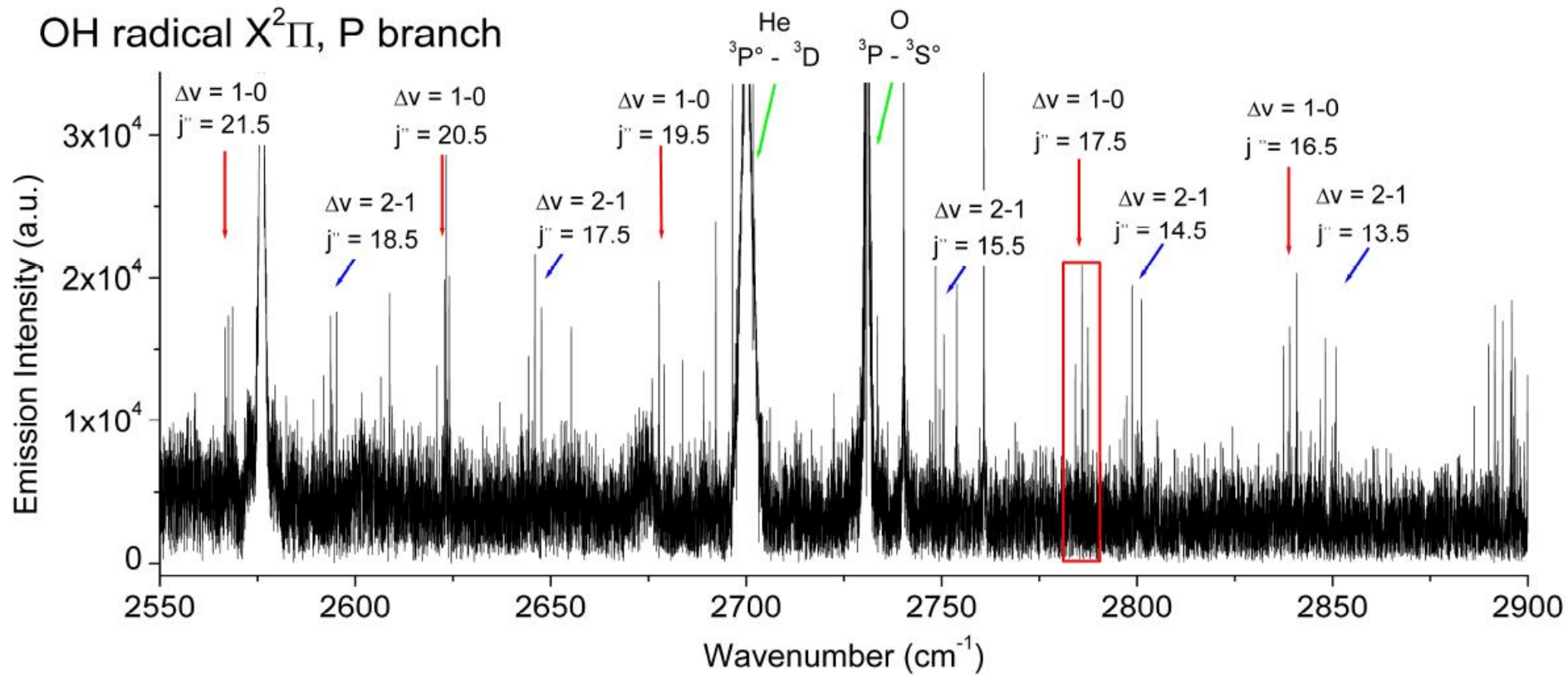
Result - information about reaction dynamics

Decomposition of extinguish agents CF_3Br and $(\text{CF}_3)_2\text{CH}$

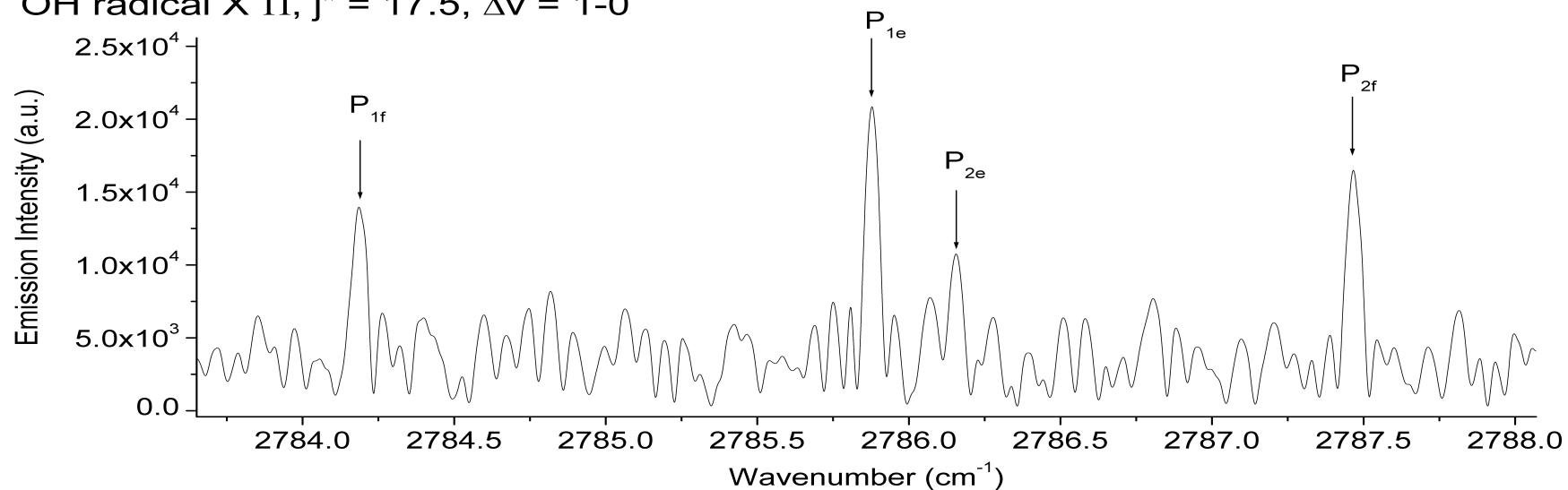
- Emission intensity time profile.
- Broad spectral range.
- High resolution (rovibrational structure)

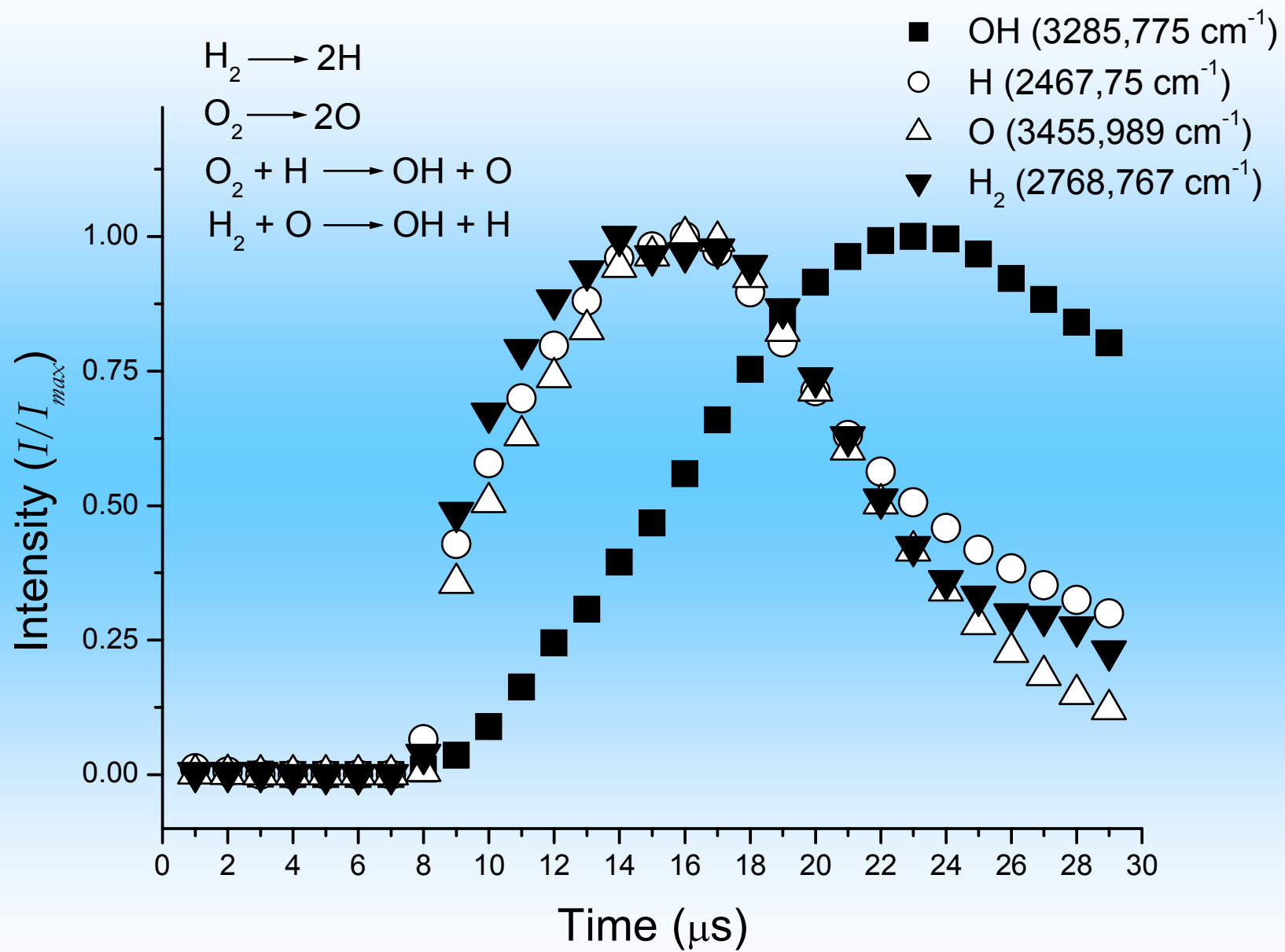


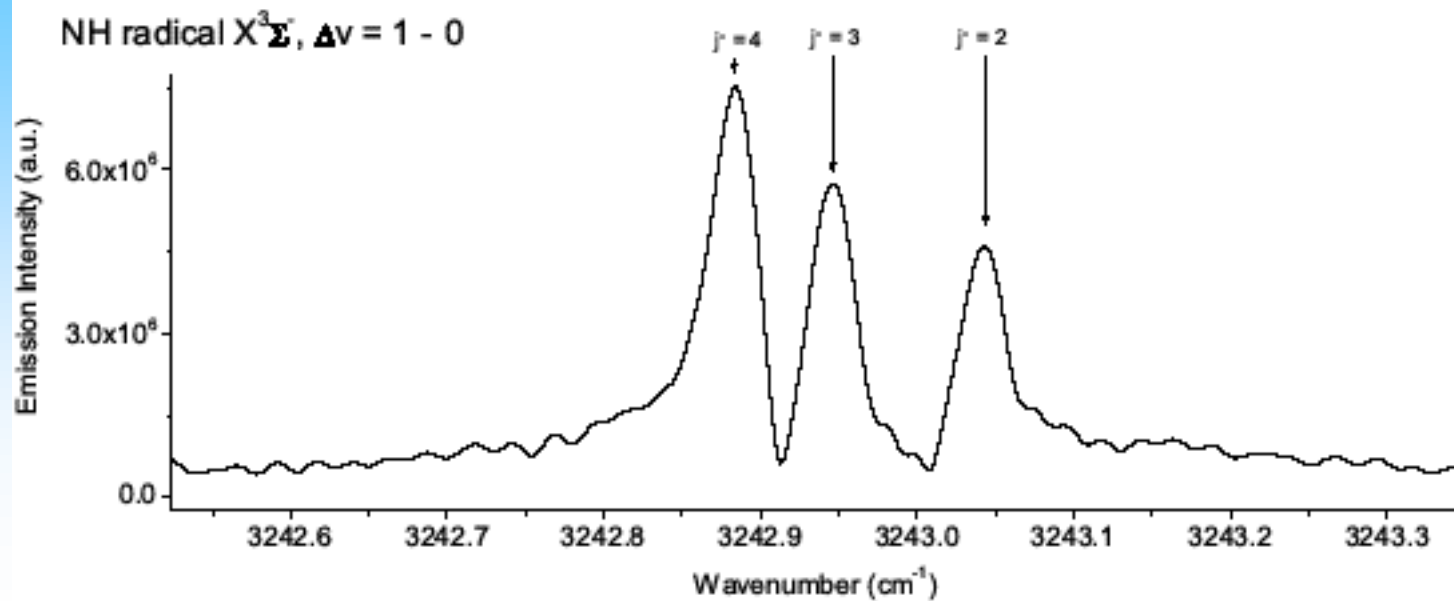
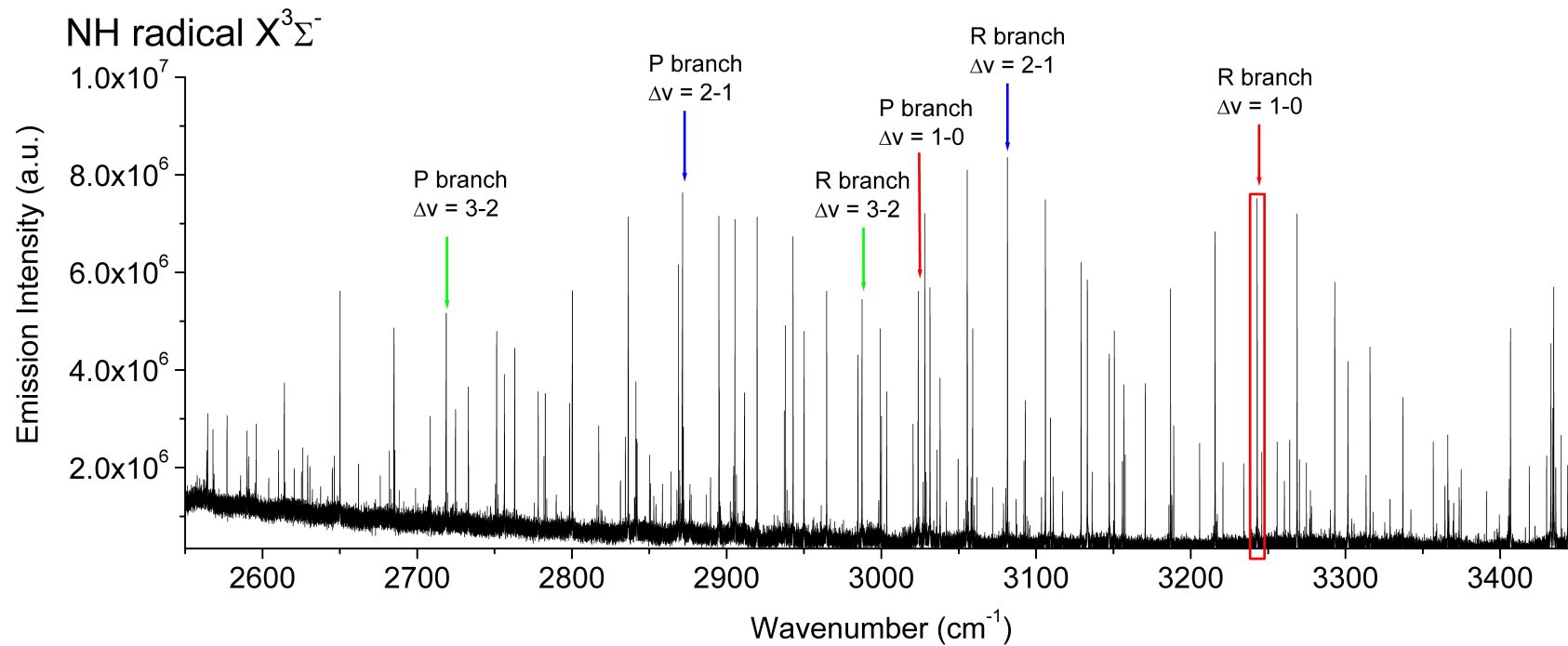
OH radical $X^2\Pi$, P branch



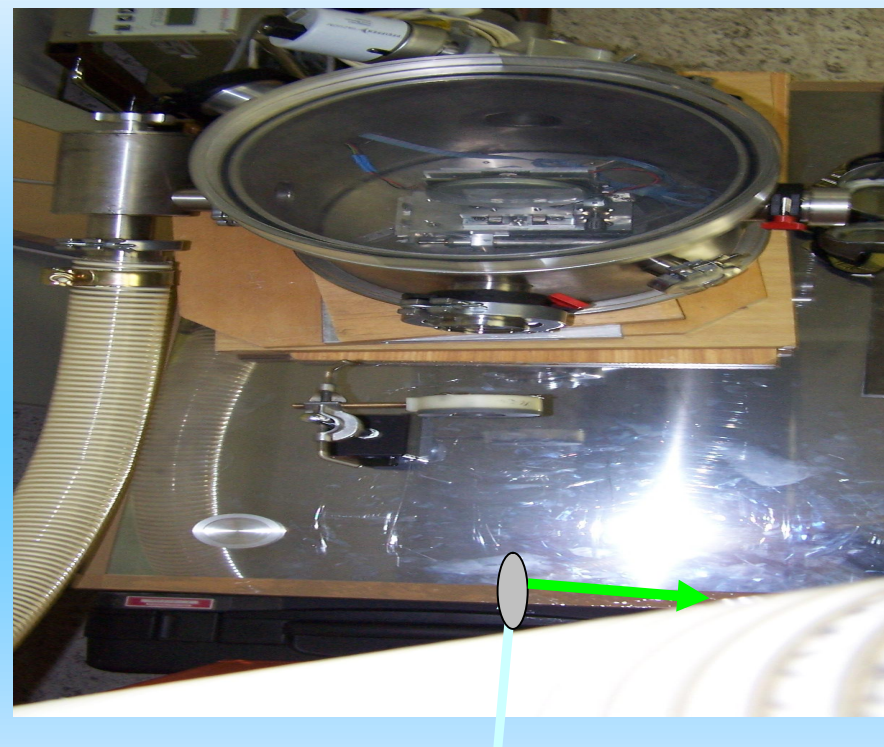
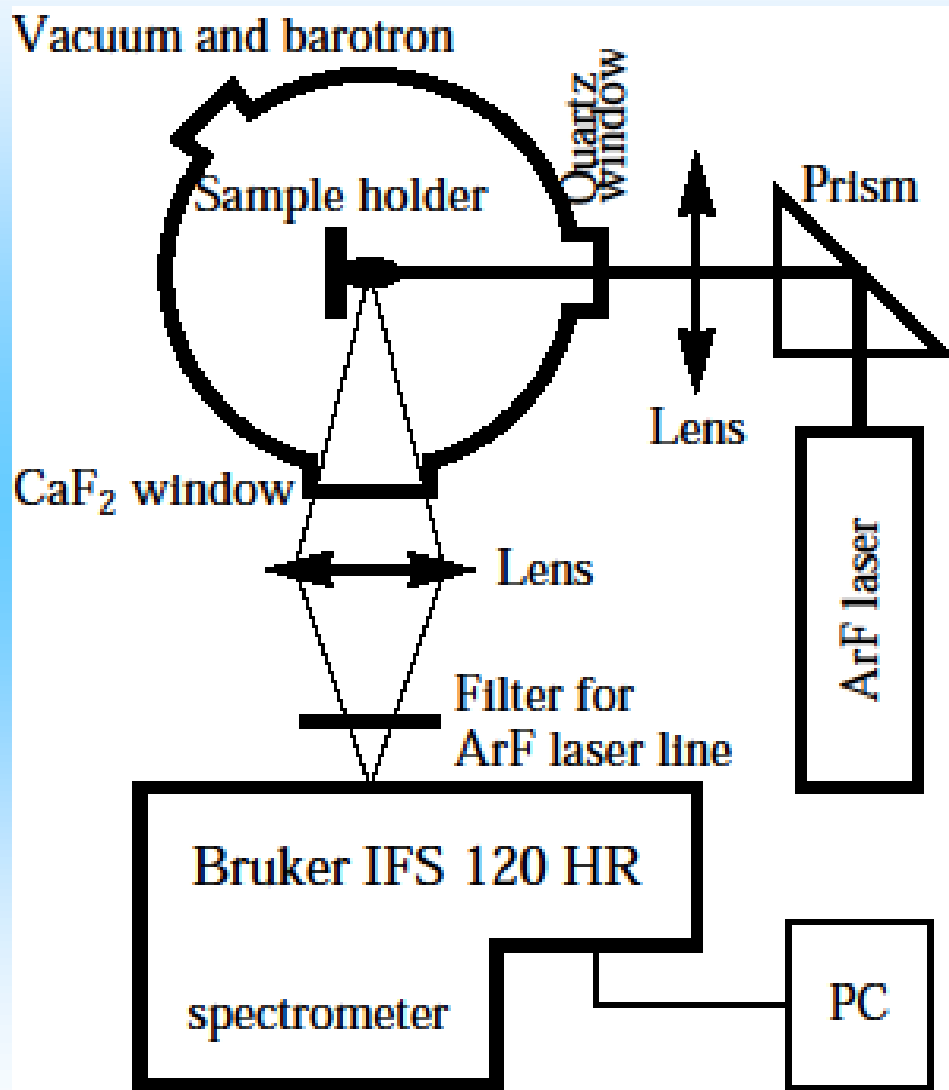
OH radical $X^2\Pi$, $j'' = 17.5$, $\Delta v = 1-0$





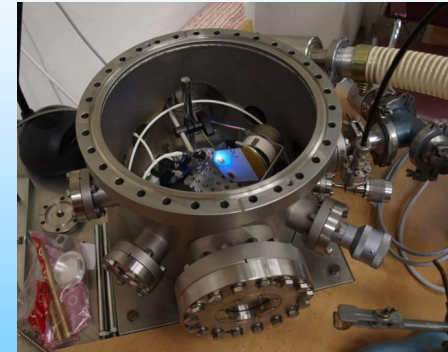
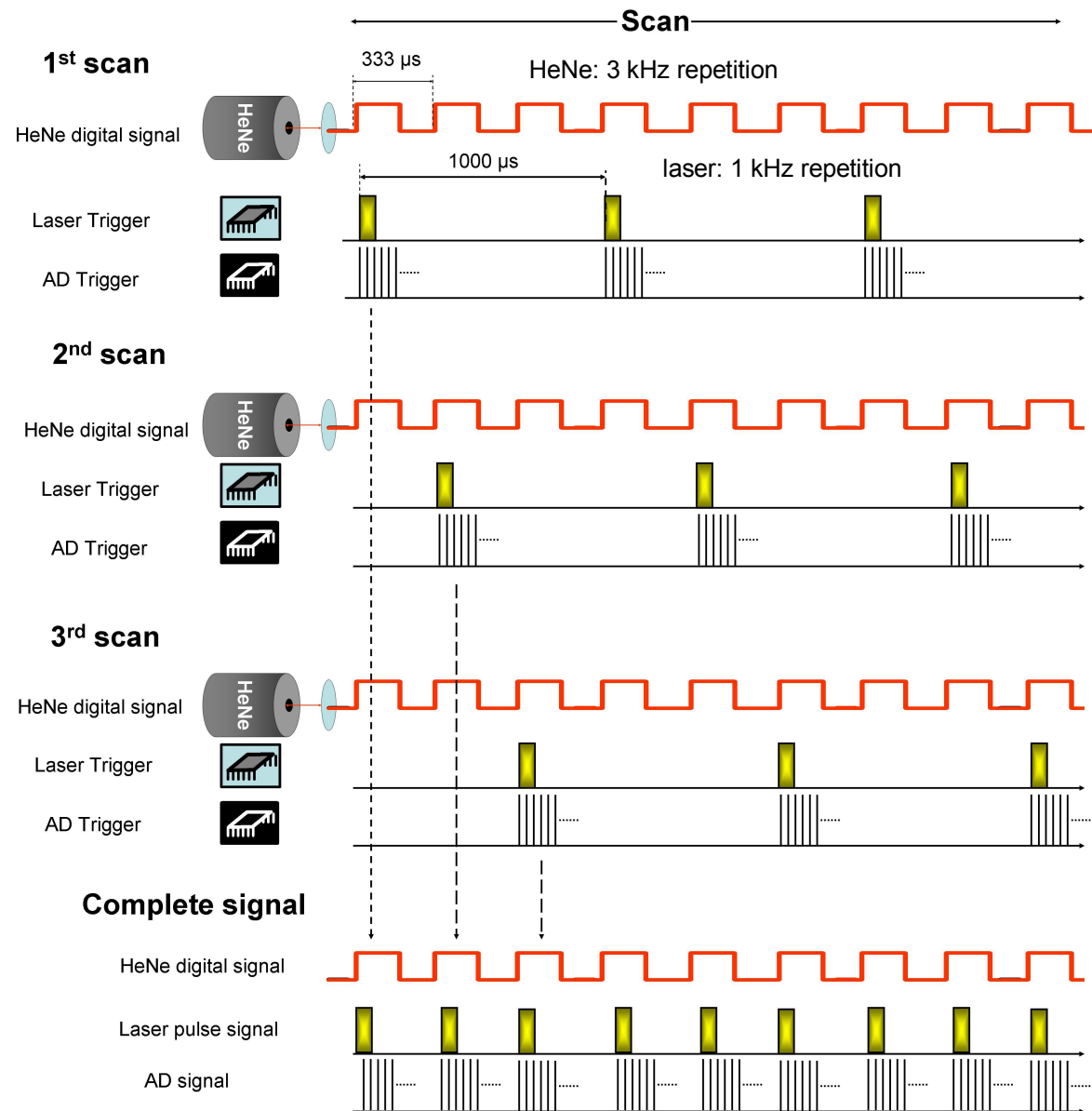


Time resolved FTIR measurement in the laser spark

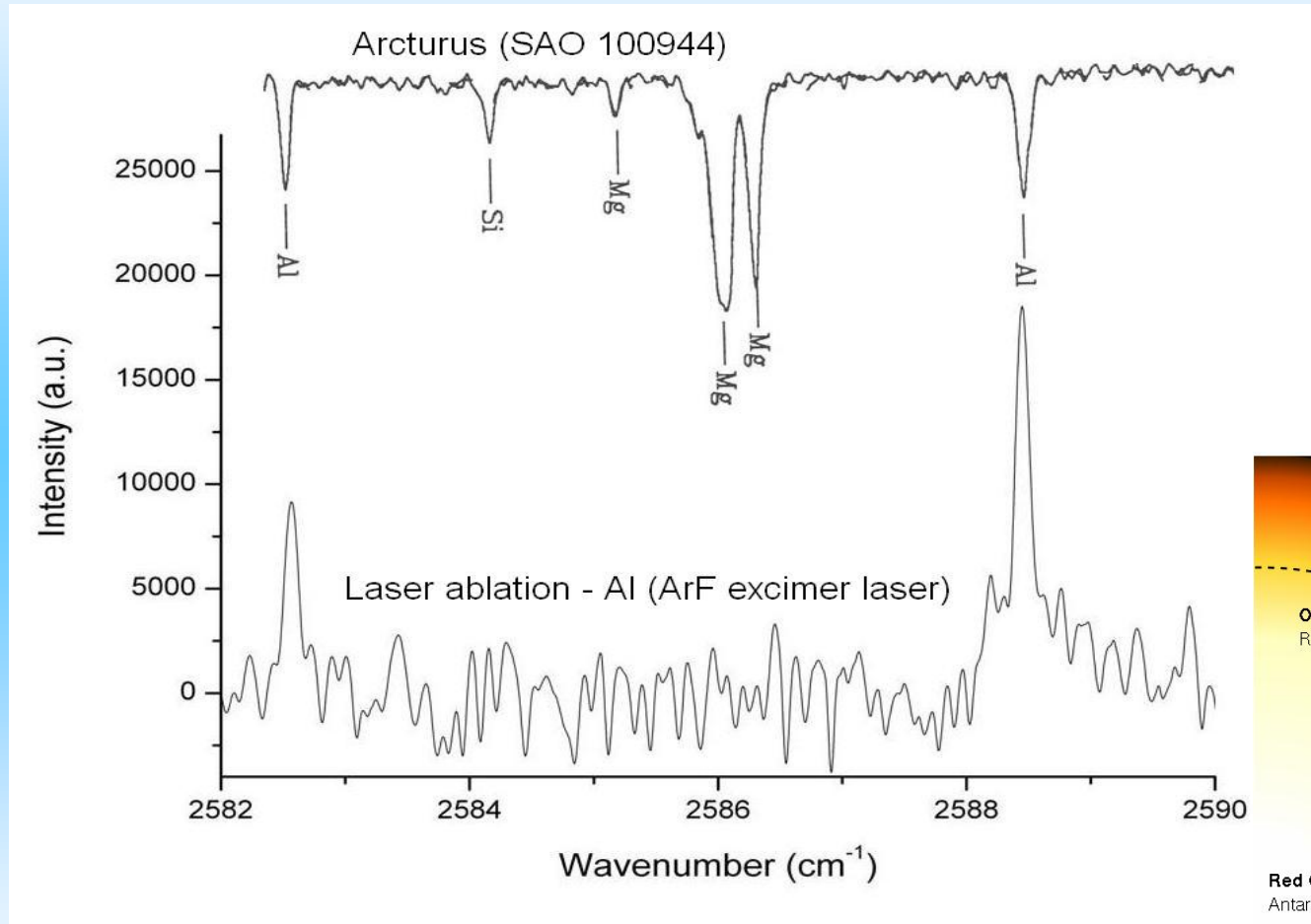


The emission arises in ablation of metal target by laser radiation of a pulsed nanosecond ArF ($\lambda = 193$ nm) laser 1 kHz repetition, at fluences around 15 mJ

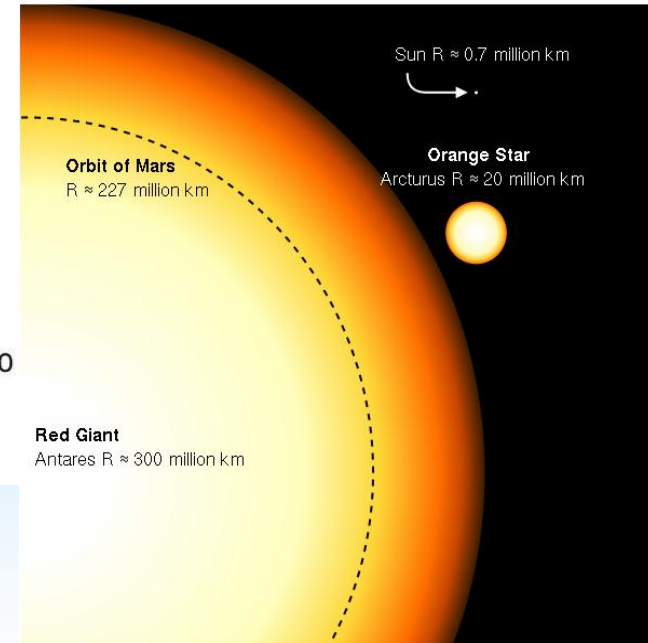
1/n system



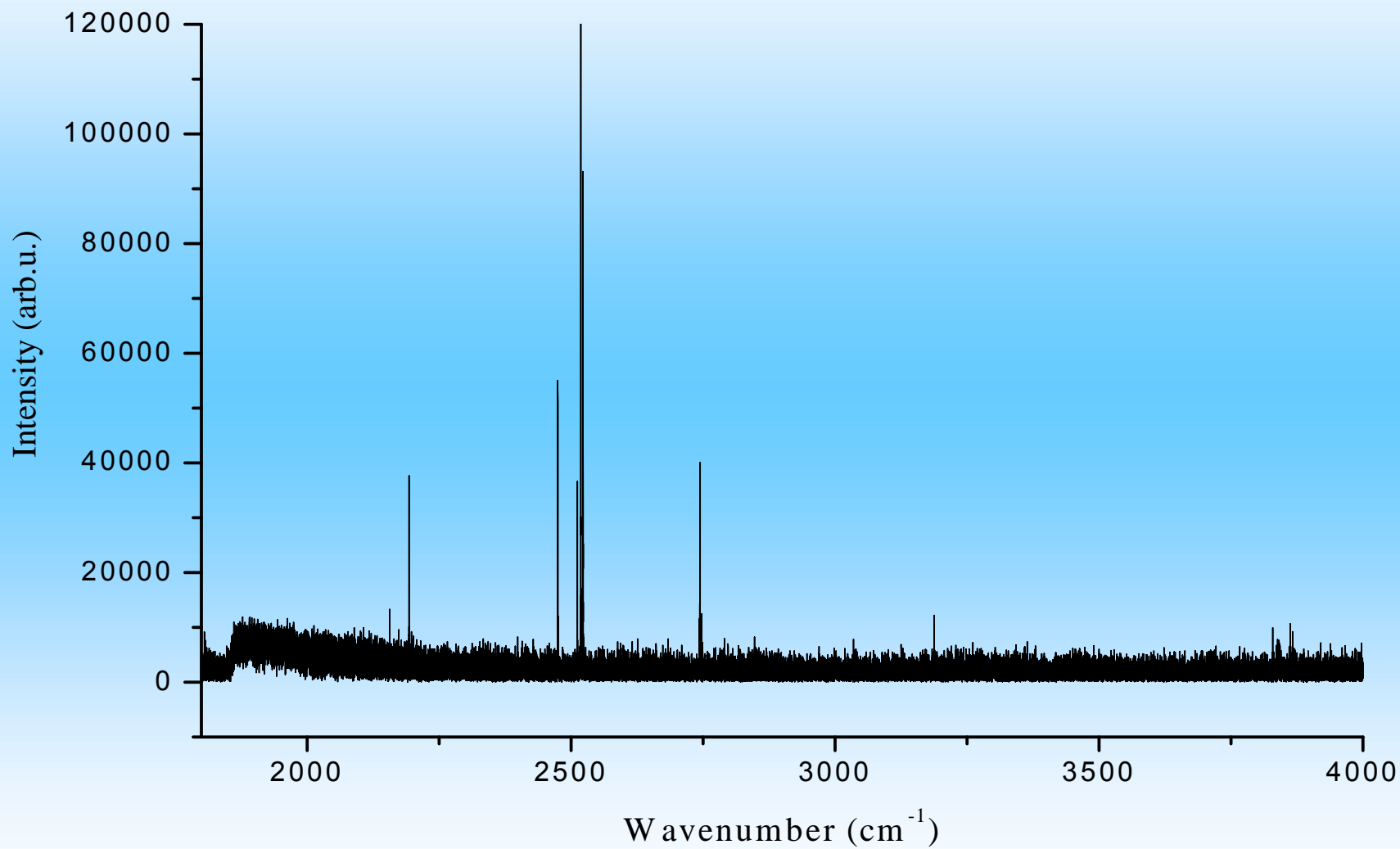
Comparison of Arcturus spectrum with FTIR emission spectrum



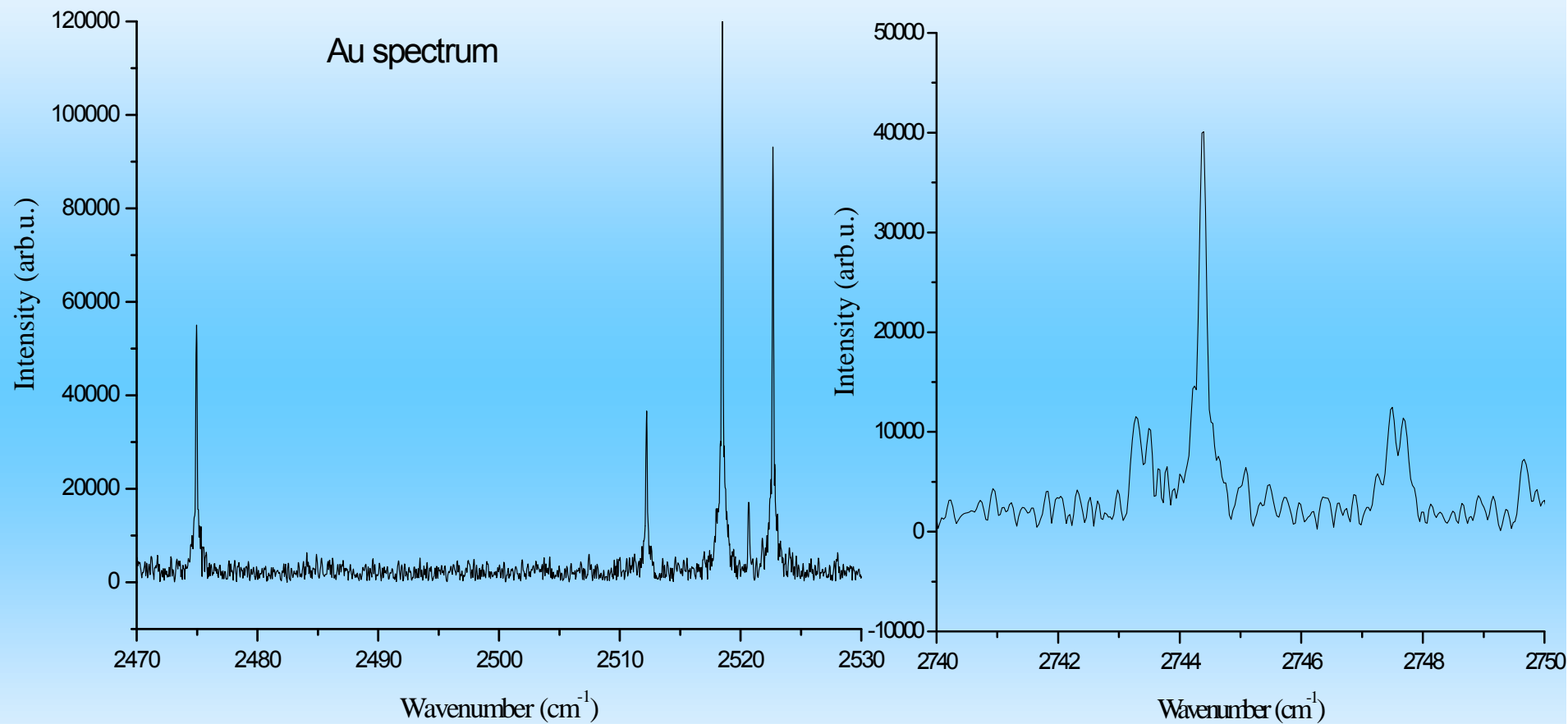
Arcturus' size relative to the Sun



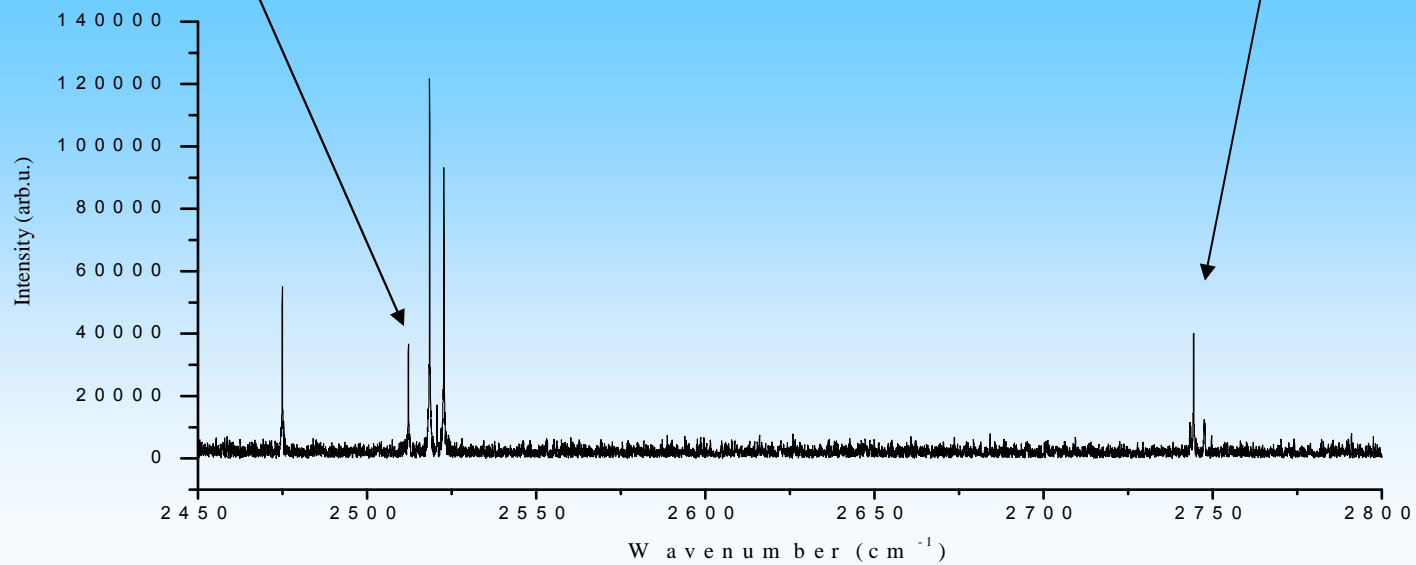
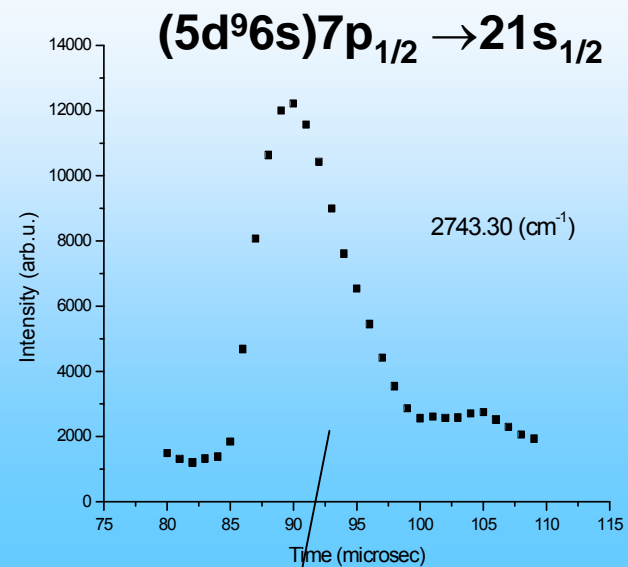
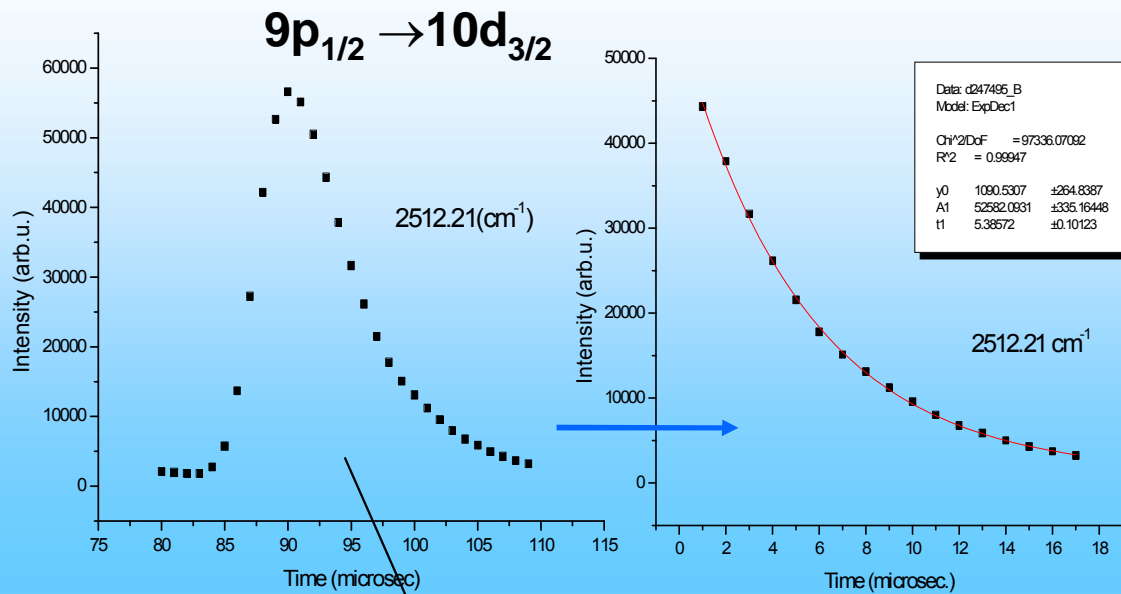
The observed IR emission spectrum of Au



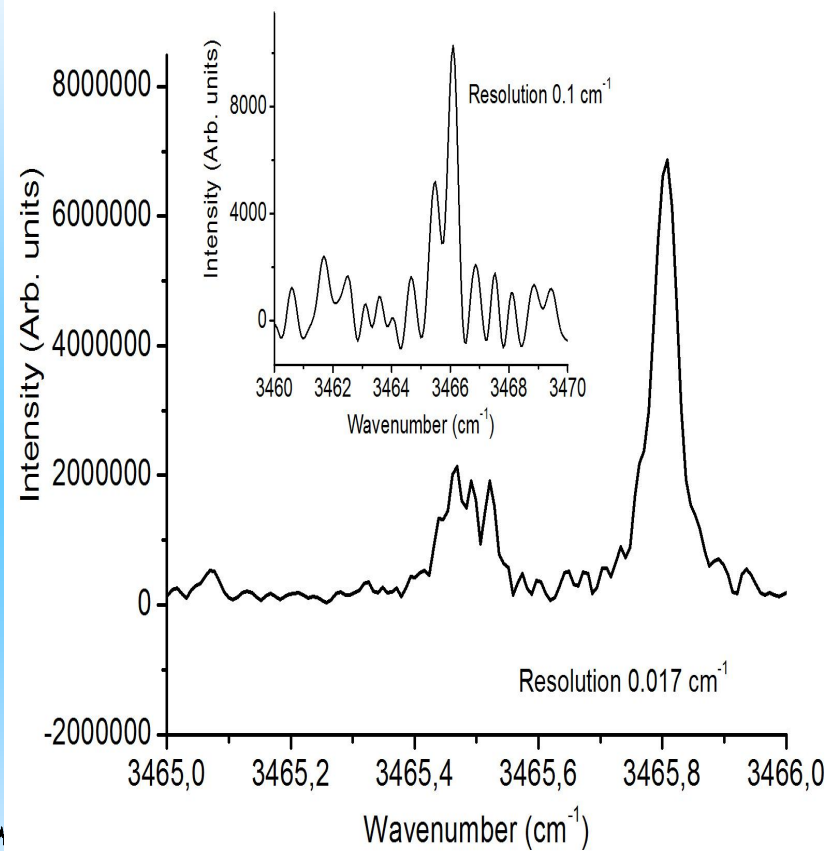
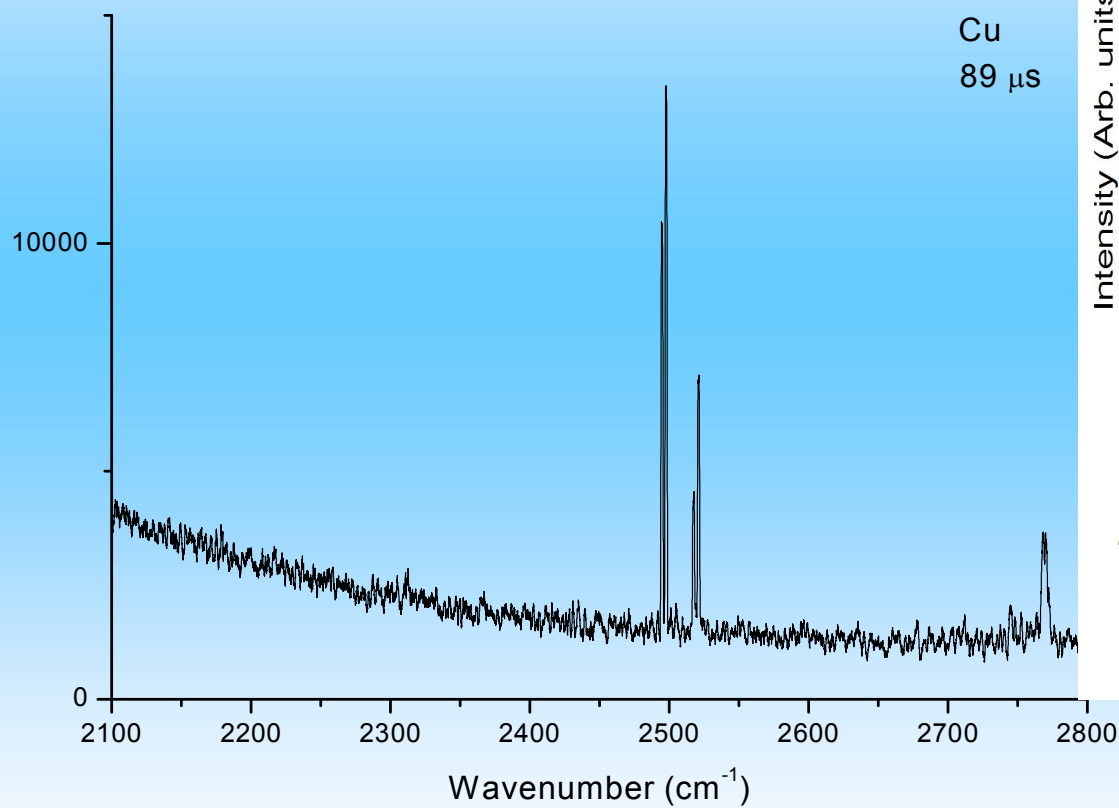
Some parts of the observed IR emission spectra of Au

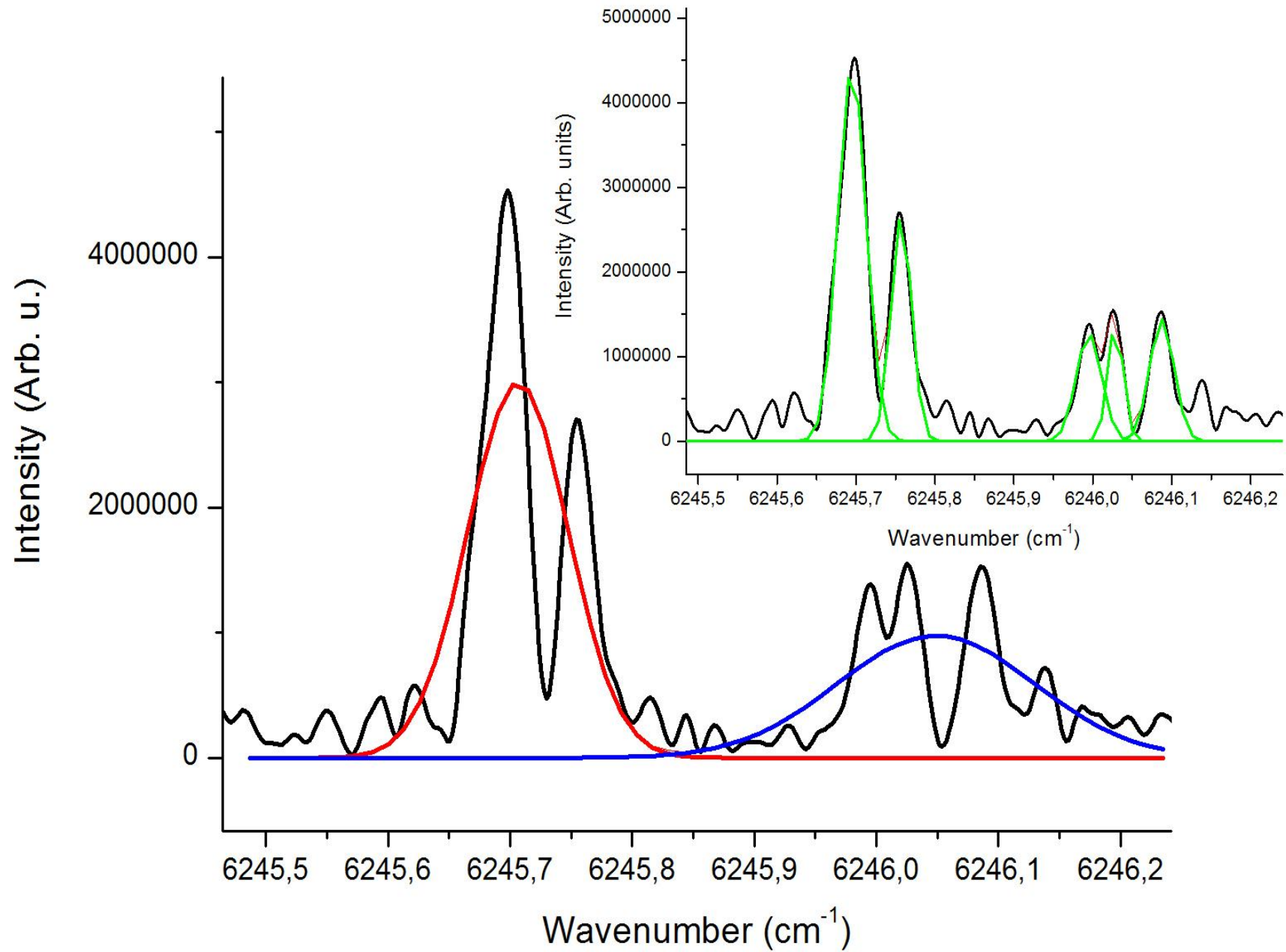


Spectral resolution 0.05 cm⁻¹



Part of the FTIR spectrum of the emission from Cu target





Conclusions

Several systems inside of discharge plasma has been studied by High resolution FT spectroscopy

New systems such extinguishing agents will be studied (C

Laboratory chemistry of HCN and HNC isomerisation

Many High resolution atomic spectra has been measured by High resolution FT spectroscopy (better understanding of the stellar spectra)

New detection technique was developed which can be applied for a study of the relaxation profiles of the individual atomic or molecular levels

THANKS:

This work is the part of the research programs of Ministry of Finance, Swiss contribution: **Spolupráce a přenos know-how v oblasti výzkumu laserové technologie, project PF 049/4V**

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Thanks

<http://www.jh-inst.cas.cz/~ftirlab/>

(Spectral library)



From the left: M. Ferus, M. Kamas, J. Cihelka, P. Kubelík, S. Civiš, K. Sovová, I. Matulková.



Bruker IFS 125 HR

Detectors

InSb
1800 – 8 000 cm^{-1}

MCT
450 – 6 000 cm^{-1}

Sample

Interferometer

Moving Mirror
Compartment

Moving Mirror

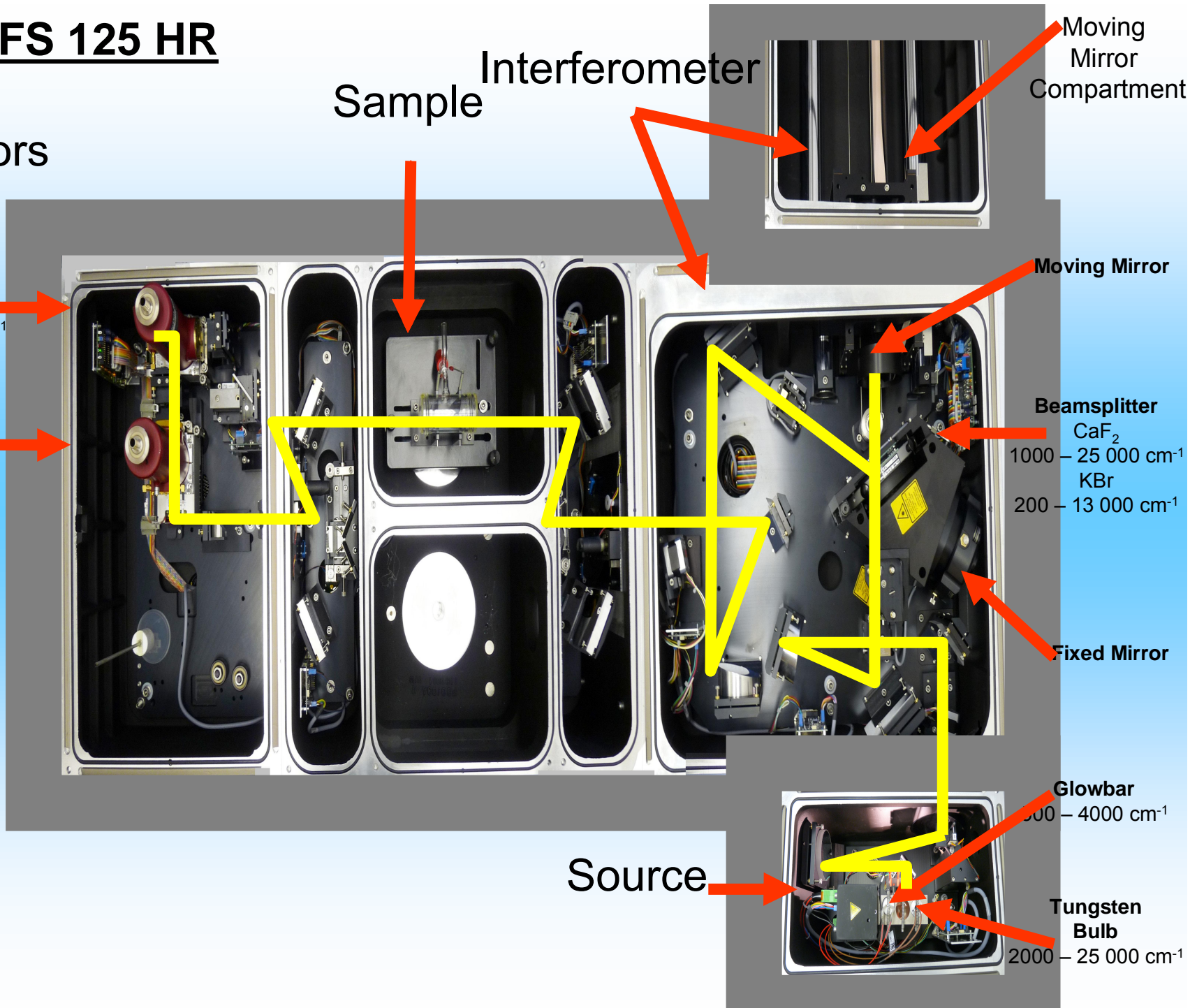
Beamsplitter
CaF₂
1000 – 25 000 cm^{-1}
KBr
200 – 13 000 cm^{-1}

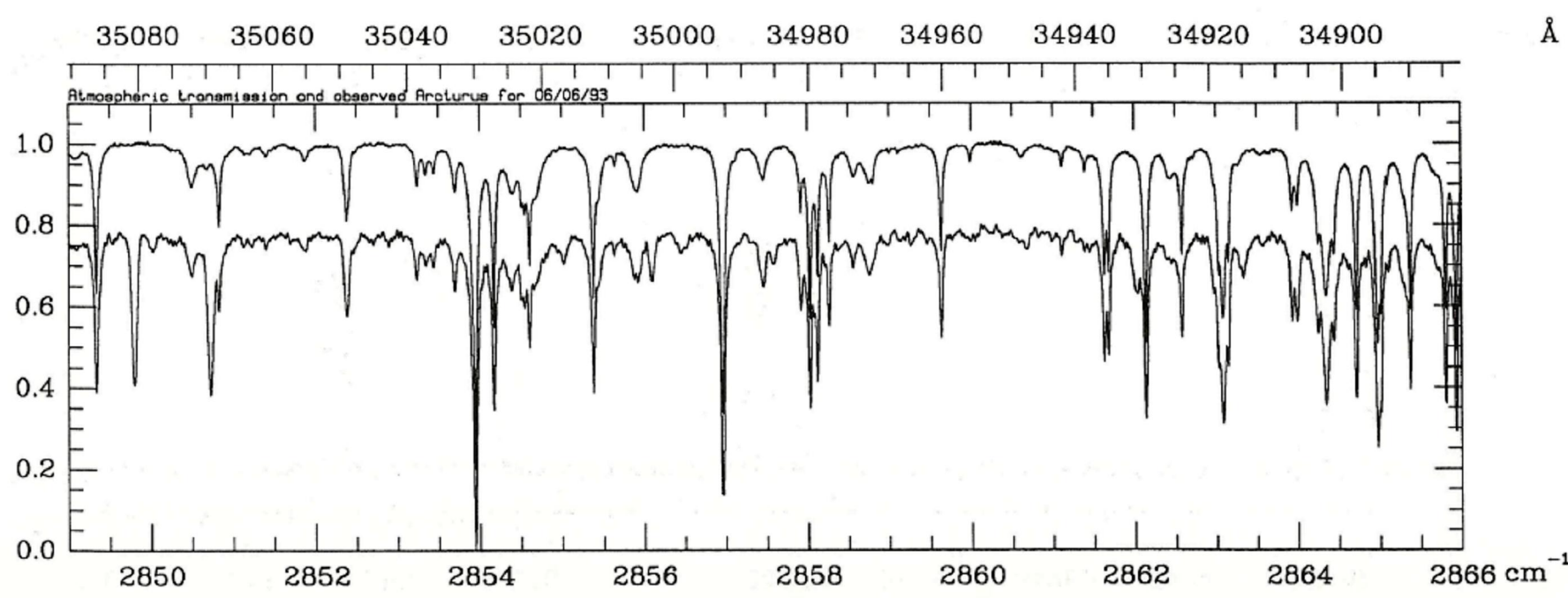
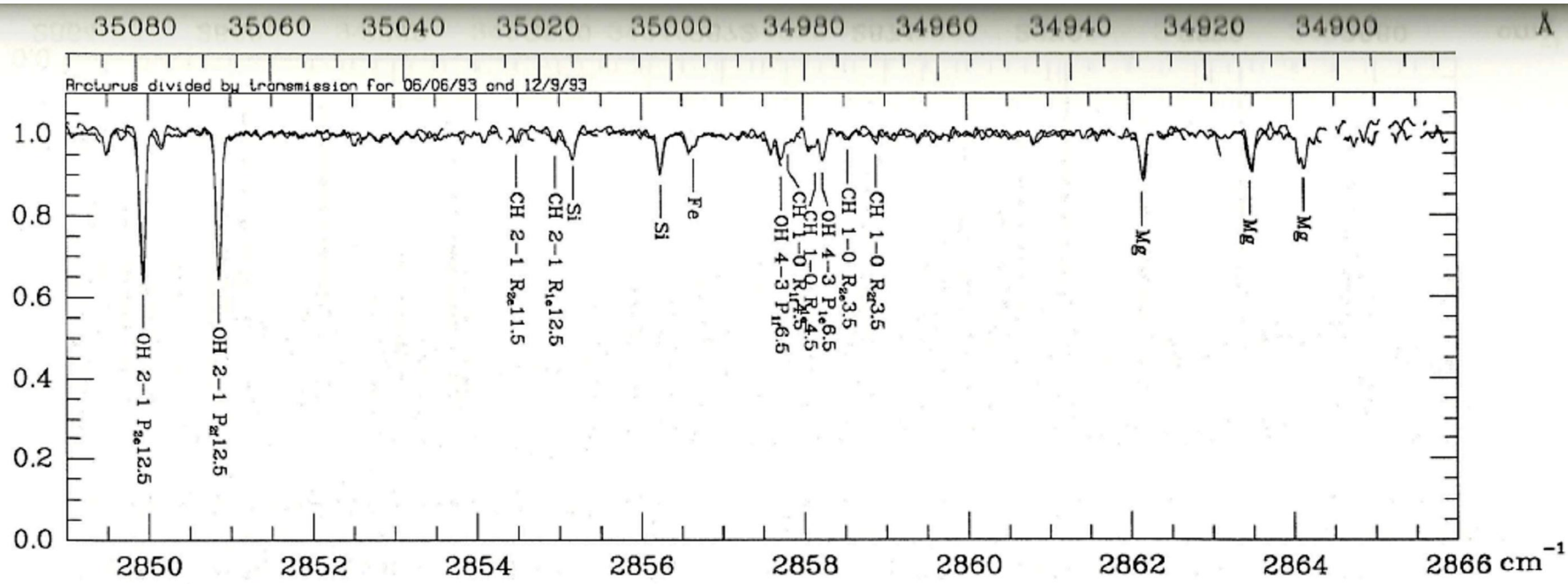
Fixed Mirror

Glowbar
500 – 4000 cm^{-1}

Source

Tungsten
Bulb
2000 – 25 000 cm^{-1}





Letter to the Editor

Radio Astronomical Determination of Ground State Transition Frequencies of CH

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Onsala Space Observatory

Received June 26, revised July 9, 1974

Summary. From a comparison of the spectra of CH and other molecules towards Cas A and dark clouds, we have determined the rest frequencies of the three allowed transitions of the ${}^2\Pi_{1/2}$, $J=1/2$, ground state A -doublet of CH to be $\nu_{01} = 3263.794 \pm 0.003$, $\nu_{11} = 3335.481 \pm 0.002$ and $\nu_{10} = 3349.193 \pm 0.003$ MHz. These frequencies cannot be checked by the sum rule, $\nu_{11} + \nu_{00} = \nu_{10} + \nu_{01}$, since the ν_{00} -transition is forbidden in this particular case. The relative errors in

these frequencies are about twice as small as those obtained earlier by radio astronomical methods for the two main lines of ground state OH. Further refinements in the determination of the CH frequencies are desirable, since these transitions have not yet been directly measured in the laboratory.

Key words: CH — Cas A — dark clouds — transition frequencies — interstellar molecules

Radio astronomical detection of the ground state ${}^2\Pi_{1/2}$, $J=1/2$, A -doublet of CH has recently been announced by Rydbeck *et al.* (1973). The upper satellite line was detected independently by Turner and Zuckerman (1974). The CH-lines were observed in southern galactic sources by Robinson *et al.* (1974). Measurements of satellite and main line intensity ratios of CH have recently been reported by Rydbeck *et al.* (1974a). The subject of the present Letter is a more accurate determination of the CH frequencies.

The equipment and observing procedure were the same as described by Rydbeck *et al.* (1974a), except that filters of bandwidth and spacing 1 kHz were also used for the present observations. Most of the integrations were done with left circular polarization, there being no evidence for polarization in any of the sources observed.

The CH transition frequencies were determined by Rydbeck *et al.* (1973) to be $\nu_{10}(F=1 \rightarrow 0) = 3349.185 \pm 0.010$ MHz, $\nu_{11}(F=1 \rightarrow 1) = 3335.475 \pm 0.010$ MHz, $\nu_{01}(F=0 \rightarrow 1) = 3263.788 \pm 0.010$ MHz. Those values have now been refined, with the following assumptions: (1) in the sources observed, CH, OH, H_2CO and NH_3 are well mixed, and have the same radial velocity; (2) the correct frequencies for the main lines of OH at 18 cm are those given by ter Meulen and Dymanus

(1972); (3) the frequency for the $F=2 \rightarrow 2$ component of H_2CO at 6 cm is that given by Tucker *et al.* (1971).

For the comparison of OH and CH spectra, we observed towards Cas A and four dark clouds which, according to Turner (1973), have either a single velocity component or at least one narrow, well-defined component.

The profile for the main $F=1 \rightarrow 1$ line of CH, in the direction of Cas A, was compared with the OH profiles taken by Davies and Matthews (1972) and with unpublished OH data taken at Onsala. The presence of two components in the Orion arm feature at $v \approx -1$ km s $^{-1}$ (Fig. 1) and the characteristic structure of the Perseus arm feature allows the CH and OH spectra to be compared with an accuracy of about ± 1 kHz. We estimate a correction, $\delta\nu_{11}$, of 6 ± 1 kHz to our previously published value of ν_{11} (Rydbeck *et al.*, 1973).

Another estimate of $\delta\nu_{11} = 5.5 \pm 2$ kHz was obtained by comparing the main line spectrum of CH with that of H_2CO (Davies, 1973; Davies, 1974; Zuckerman *et al.*, 1970) in the direction of Cas A.

A third, independent, estimate of $\delta\nu_{11}$ was obtained from a comparison of our CH spectra with those of other molecules towards four dark clouds. Table 1 gives a compilation of the radial velocities for CH and other molecules observed at approximately the same positions. The OH values are averages of the data of

* On leave from the Department of Physics and Astronomy, University of Massachusetts, Amherst.

Title: Les bandes de CH et la présence de l'hydrogène dans les Comètes.

Authors: Nicolet, M.

Journal: Zeitschrift für Astrophysik, Vol. 15, p.154

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**Les bandes de CH et la présence de l'hydrogène
dans les Comètes.**

Par **M. Nicolet**, D. Sc.

(Reçu janvier 12, 1938.)

Il est montré que l'étude de la structure de rotation est nécessaire pour l'identification des bandes cométaires et que le groupe de raies cométaires vers λ 4300 Å doit probablement être attribué à la molécule CH.

Les raies de la série de Balmer n'apparaissent certainement pas dans les spectres des comètes et les bandes de résonance de H_2 , étant en dehors de la région spectrale astronomique, ne peuvent être observées. Quant à la présence des bandes cométaires attribuées au système de Raffety, elle ne peut être envisagée. Rien ne permet donc d'établir jusqu'ici la présence de l'hydrogène dans les Comètes. Cependant, si l'identification de la molécule CH était assurée, la présence d'hydrogène dans les Comètes serait du même coup démontrée.

F. BALDET ¹⁾, R. C. JOHNSON ²⁾ et N. T. BOBROVNIKOFF ³⁾ ont examiné la possibilité de l'identification de CH, par la comparaison du maximum d'intensité de la bande λ 4300 Å de CH avec λ 4313 Å. F. BALDET avait considéré la coïncidence avec la radiation cométaire à λ 4314 Å comme fortuite (Comète Brooks) parce que la comète ne montre pas une bande dégradée vers le violet, mais bien une ligne piquée. Cependant, N. T. BOBROVNIKOFF avait indiqué (Comète Halley) que cette circonstance rendait l'identification quelque peu incertaine, mais que la faible dispersion employée et les transitions restreintes aux faibles nombres quantiques de rotation pourraient donner une bande aigue de forme analogue à celle d'une raie. Bref, l'identification était extrêmement douteuse et n'était généralement pas acceptée⁴⁾.

Le but de cette note est d'indiquer que le spectre de la molécule CH est très probablement présent dans les spectres cométaires. Le composé CH

¹⁾ Thèse, Paris 1926. — ²⁾ M. N. **87**, 625, 1927. — ³⁾ Public. Lick Obs. **17**, 443, 1931. — ⁴⁾ Voir K. WURM, Handb. d. Astrophys. **7**, 305, 1936 et R. MIN-KOWSKI, P. A. S. P. **49**, 276, 1937.