

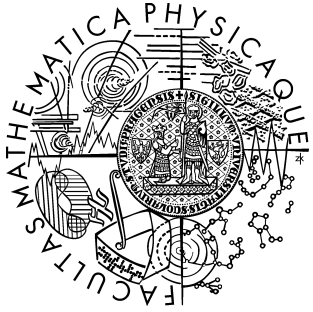
# Dijets in Diffractive Photoproduction at HERA

**Karel Černý**

Institute of Particle and Nuclear Physics

Charles University in Prague

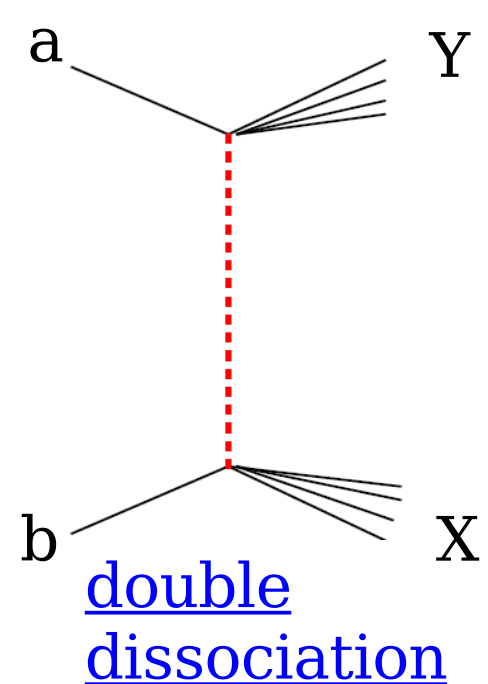
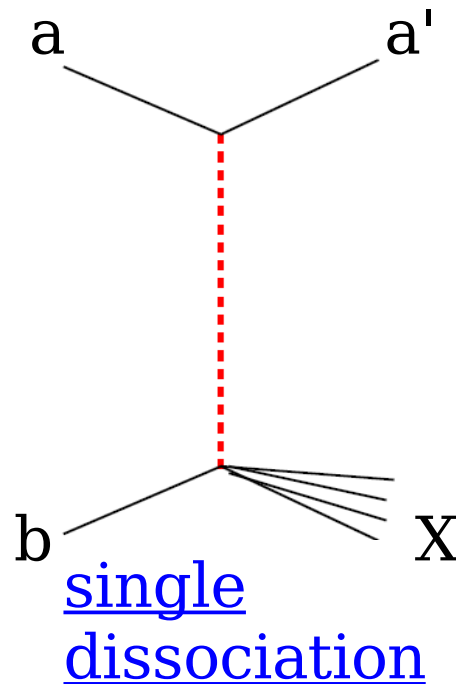
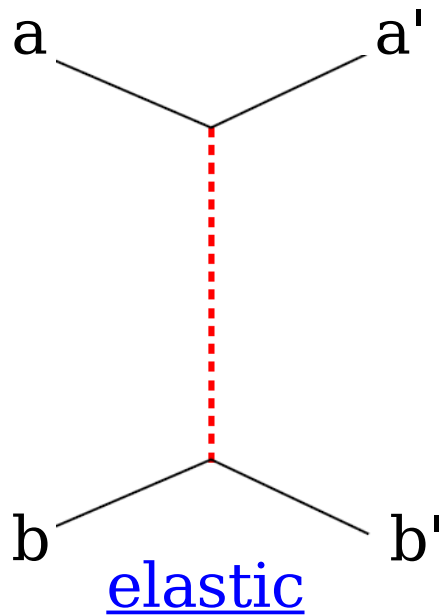
06/03/2008



- 1) Introduction
- 2) Data and Event Selection
- 3) Control Plots
- 4) Cross Section Measurement
- 5) Systematic Uncertainties
- 6) Results
- 7) Conclusions

# Introduction

**Diffraction:** interaction with vacuum quantum number exchange.

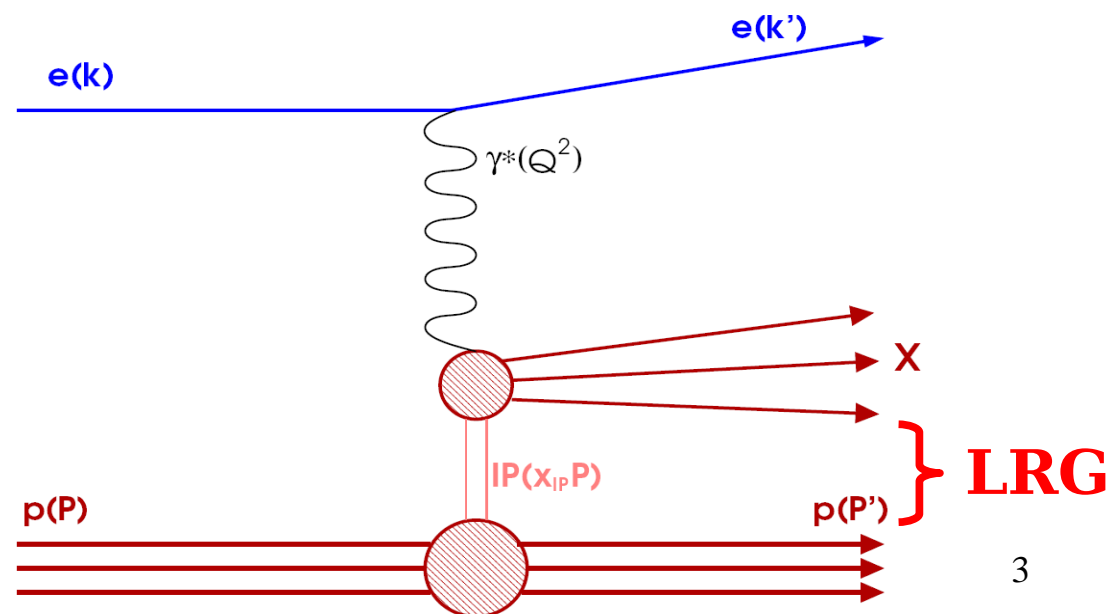
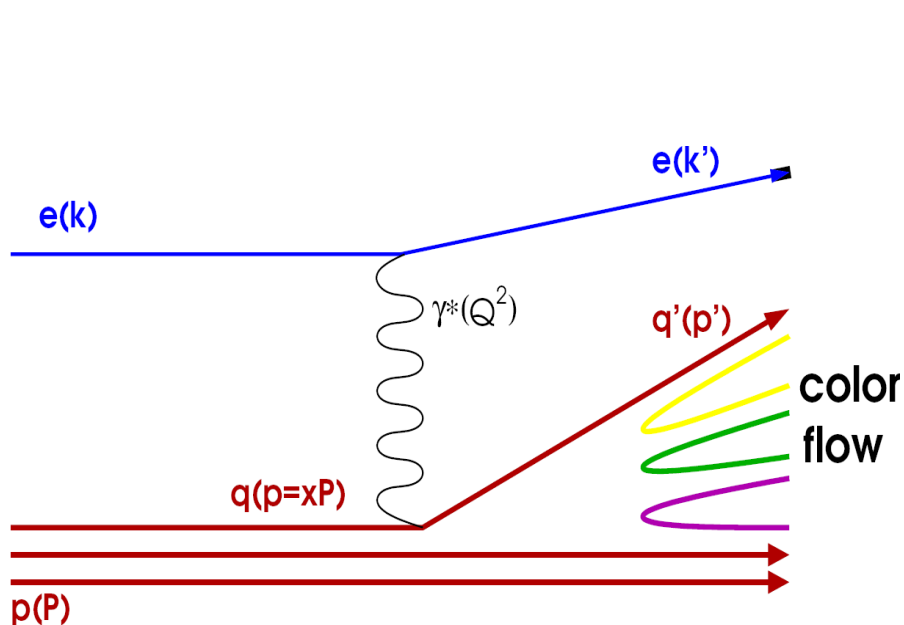


**Dissociation** into low mass states may occur:  $M_{X,Y} \ll E_{\text{CMS}}$

**Hard diffraction:** hard scale present  $\Rightarrow$  QCD calculable final states in either of the dissociated states

# Diffractive Deep Inelastic Scattering (DDIS) at HERA

- Hard diffraction in DIS was not expected to be interesting.
- $\sim 10\%$  of DIS events at low  $x$  is DDIS, however.
- Jets or vector mesons observed in the final states in DDIS.
- Due to colourless exchange the DDIS is manifested with a Large Rapidity Gap (LRG).



## Diffractive parton distribution functions (DPDF)

- In order to be able to make predictions of the cross sections of hard diffractive processes the DPDFs need to be measured.
- They are extracted from inclusive DDIS (any final state & LRG &  $Q^2$  the hard scale).
- QCD hard factorization theorem is then used for predictions.

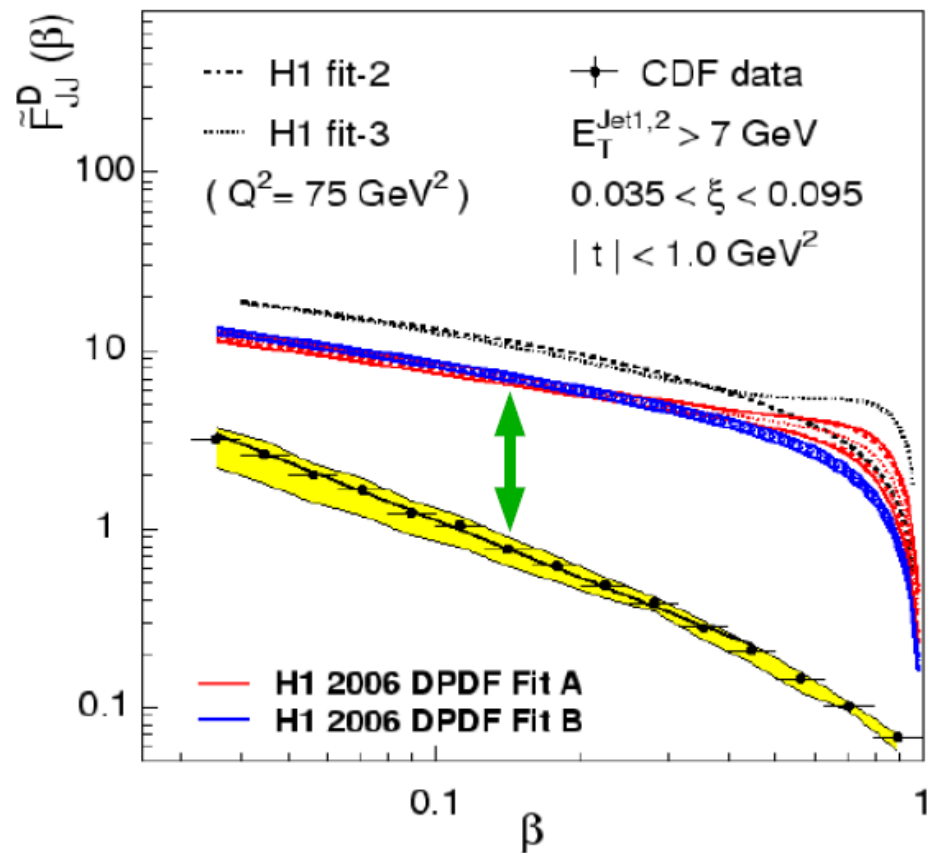
$$\sigma(\gamma^* p \rightarrow Xp) \approx f_{i/p}(z, Q^2, x_{IP}, t) \times \sigma_{\gamma^* p}(z, Q^2)$$

universal partonic cross section

Diffractive Parton Distribution Function (DPDF)

---

Diffractive parton distributions measured at **HERA** fail to predict dijet rates at **Tevatron!**

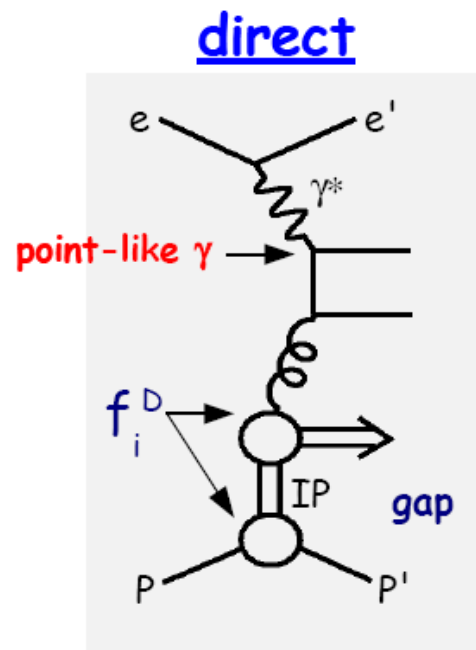


HERA DPDFs were expected to be universal but ...

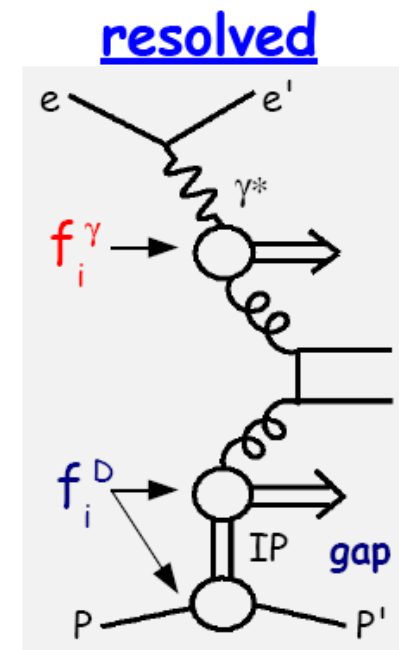
- Hard factorization **is expected** to hold in **DDIS** processes (Collins, 1998).
- Hard factorization **not expected** to hold in diffractive **h-h** scattering (Kaidalov et. al., 2003).

# Diffractive Dijets in Photoproduction

- Photoproduction =  $Q^2 \sim 0 \text{ GeV}^2$ , two leading order (LO) classes of photon interaction: **direct** / **resolved**.
- The resolved events resemble h-h scattering.



no  $\gamma$  remnant



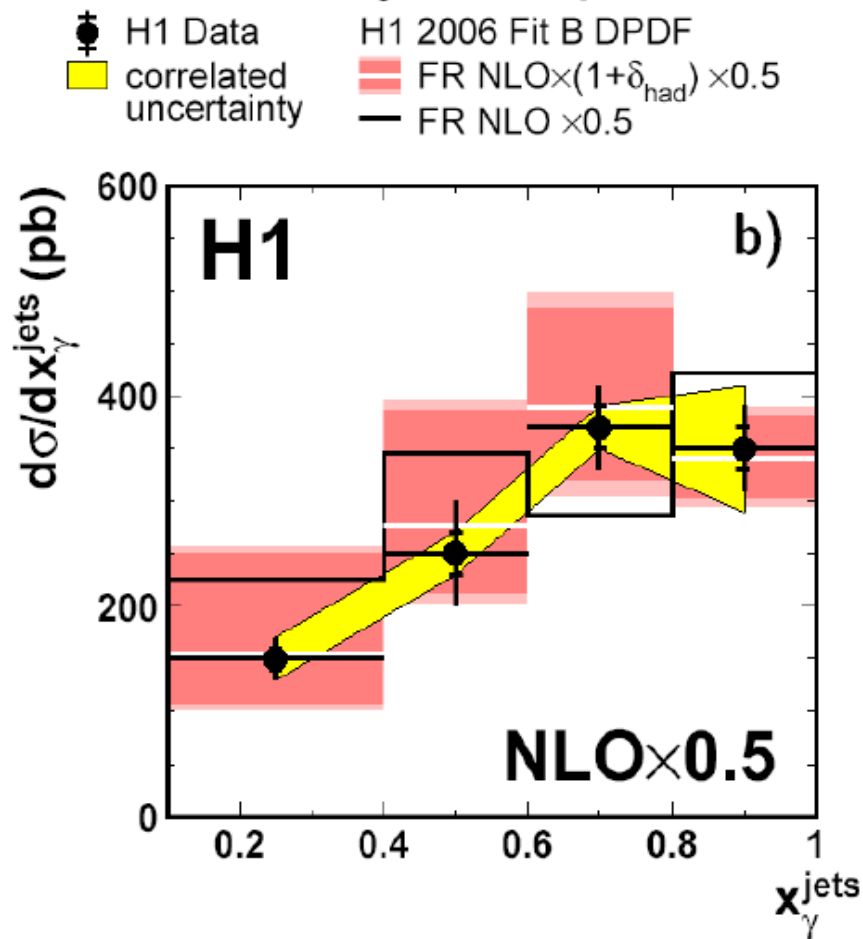
interaction with  $\gamma$  remnant can fill the Gap(LRG).

## Recent Results on Diffractive Dijets in Photoproduction

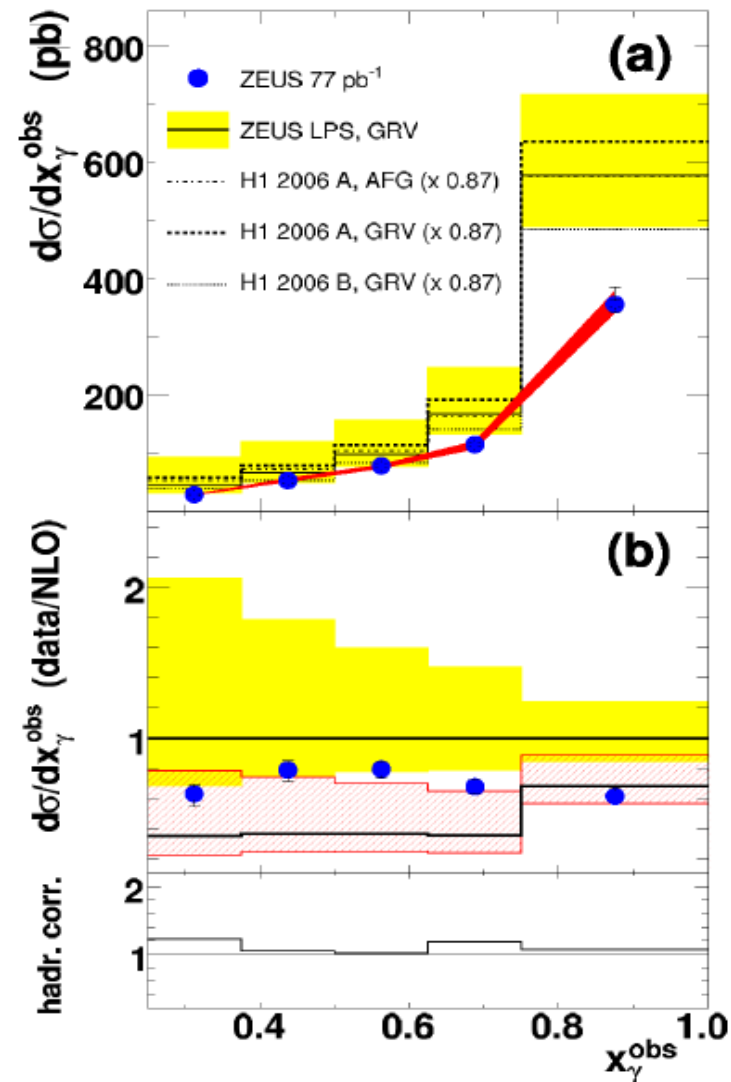
- Published **H1** results (96/97  $\sim 18 \text{ pb}^{-1}$ ) conclude that an **overall suppression factor of 0.5** is needed for NLO.
- Recent **ZEUS** results: NLOs agree with data with **weaker suppression factors** or - within the uncertainties almost without suppression.
- Both experiments observe **no special suppression of resolved w.r.t. direct.**
- There may be an  $E_T$  dependence of the suppression explaining the different H1  $\leftrightarrow$  ZEUS conclusions.
- The analysis is, therefore, carried out in two  $E_T$  cut schemes:
  - low  $E_T$  ... jets above **5** and **4** GeV in  $p_T$  (H1)
  - high  $E_T$  ... jets above **7.5** and **6.5** GeV in  $p_T$  (ZEUS)



# H1 Diffractive Dijet Photoproduction



# ZEUS



H1: suppression of NLO is needed! ... NLO  $\times 0.5$

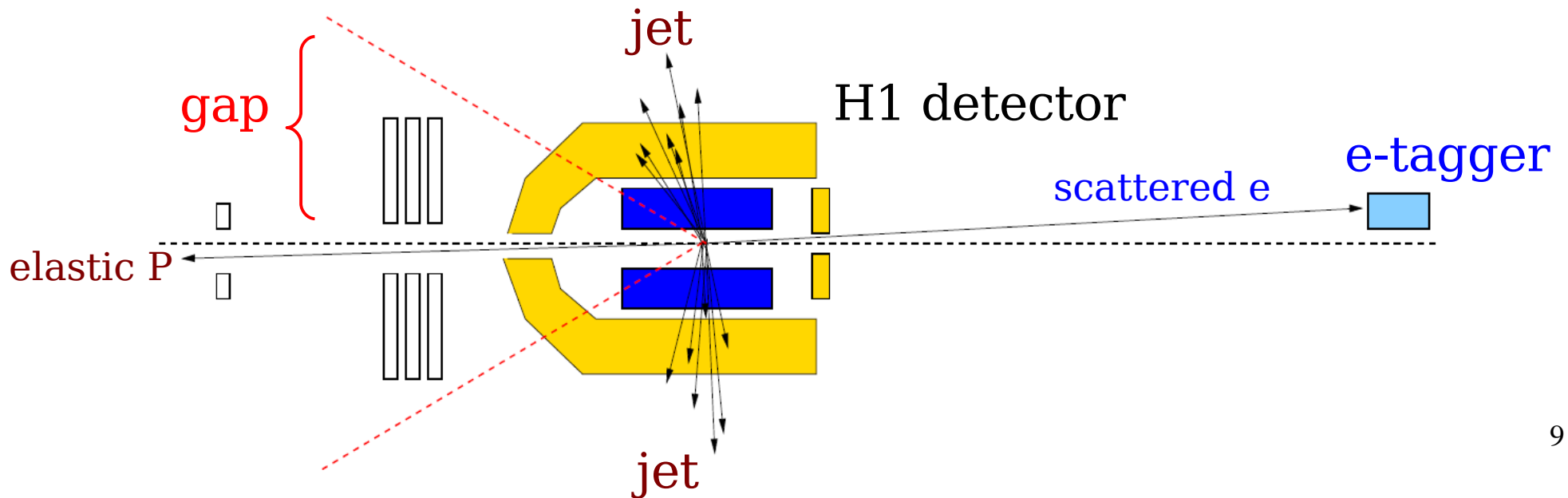
ZEUS: NLO  $\times 0.6 - 0.9$  ... varying with DPDFs

H1+ZEUS: no  $x_{\gamma}$  dependence of the suppression!

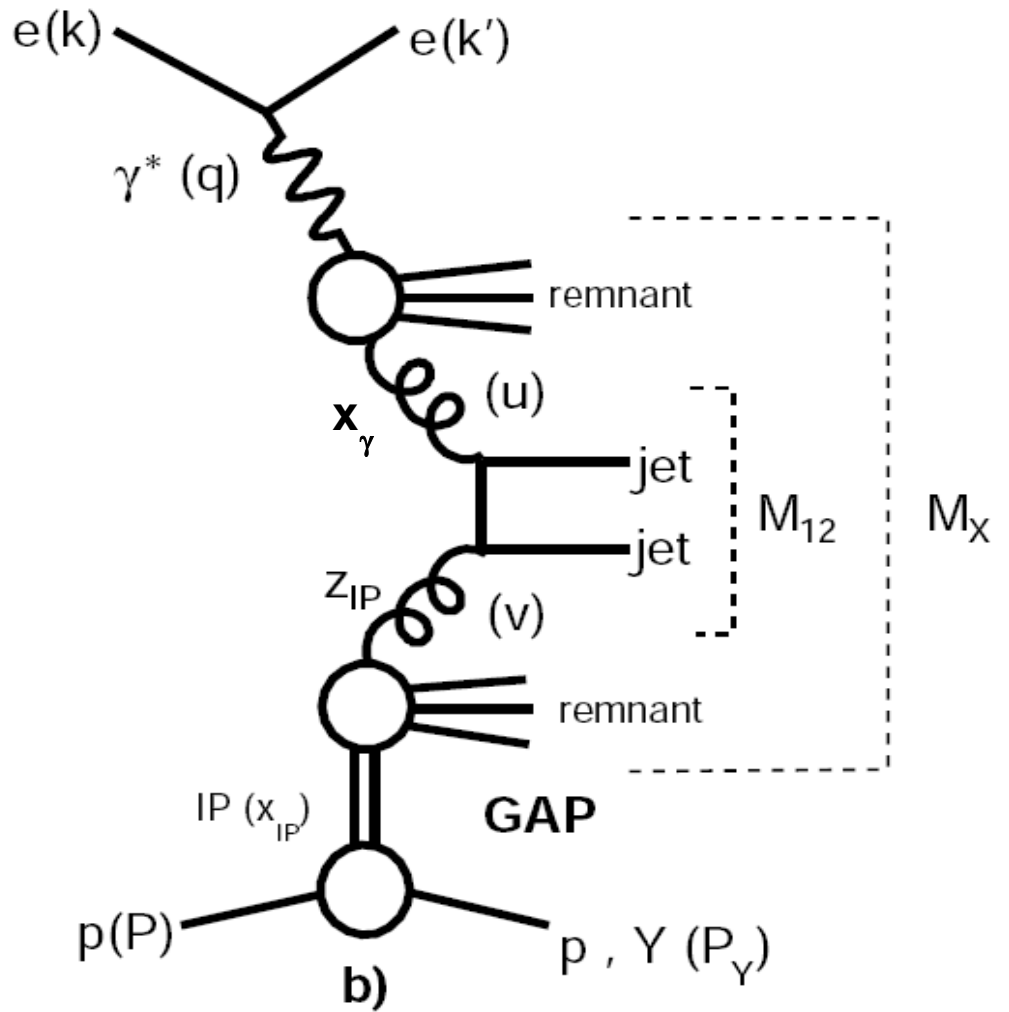
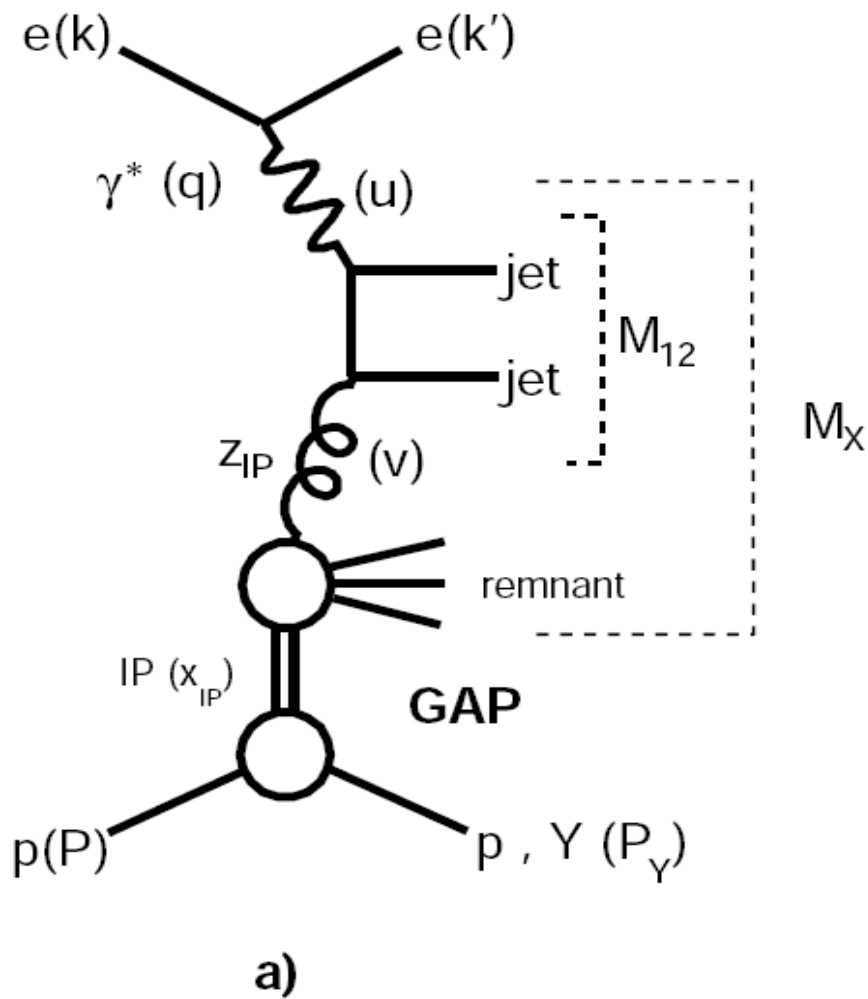


# Data and Event Selection

- Analysis is based on 1999e+/2000 HERA data.
- Integrated luminosity  $\sim 54 \text{ pb}^{-1}$ .
- Outgoing electron is **tagged** in the electron tagger (e-tag33m).
- Diffractive events selected by means of LRG method.



# Kinematics of the events



# Kinematics of the events

**low  $E_T$**

$$E_T^{\text{jet1}} > 5 \text{ GeV}$$

$$E_T^{\text{jet2}} > 4 \text{ GeV}$$

$$-1 < \eta^{(\text{jet 1 and 2})} < 2$$

$$x_{\text{IP}} < 0.03$$

e-tagger limitation

$$\left\{ \begin{array}{l} 0.3 < y_e < 0.65 \\ Q^2 < 0.01 \text{ GeV}^2 \end{array} \right.$$

$$|t| < 1 \text{ GeV}^2$$

$$M_Y < 1.6 \text{ GeV}$$

**high  $E_T$**

$$E_T^{\text{jet1}} > 7.5 \text{ GeV}$$

$$E_T^{\text{jet2}} > 6.5 \text{ GeV}$$

$$-1.5 < \eta^{(\text{jet 1 and 2})} < 1.5$$

$$x_{\text{IP}} < 0.025$$

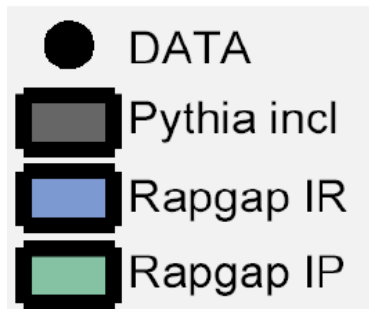
different from ZEUS

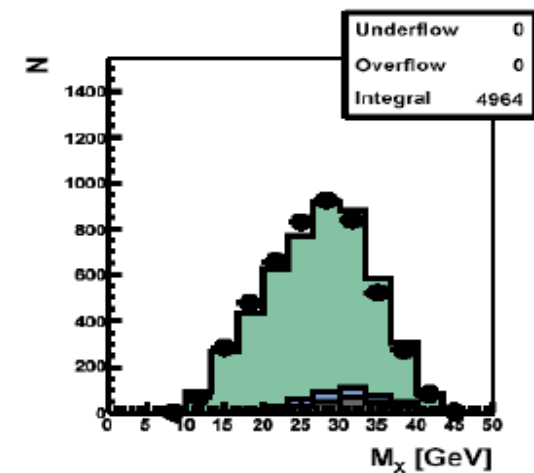
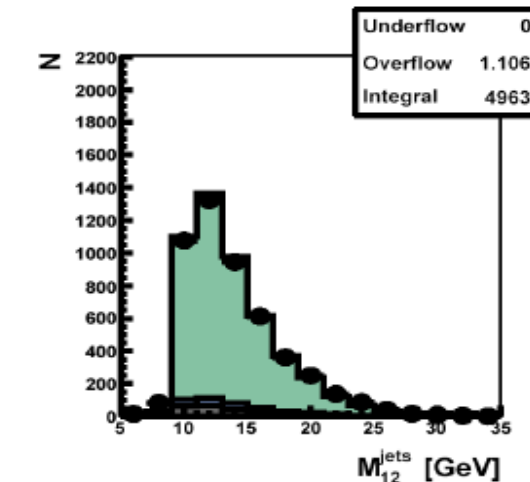
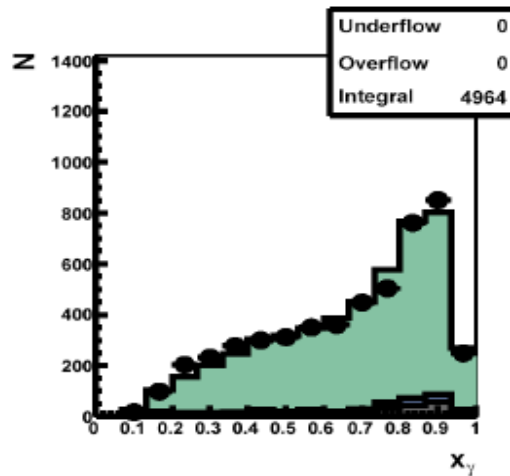
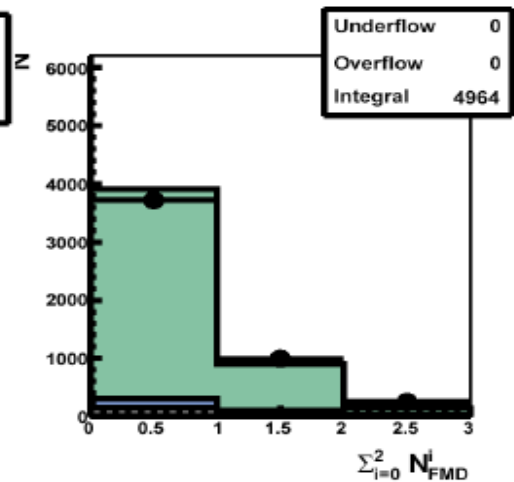
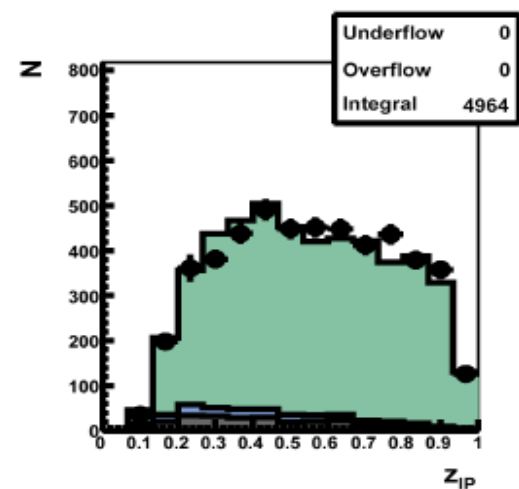
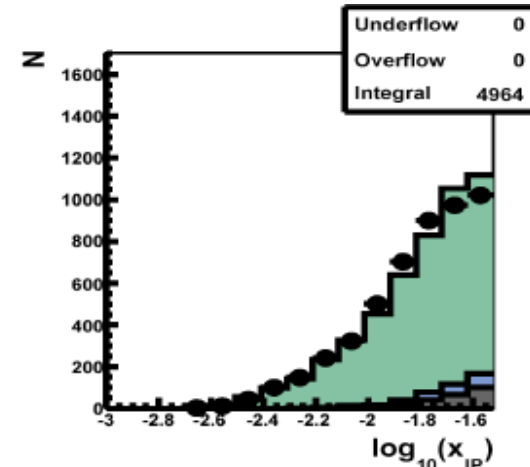
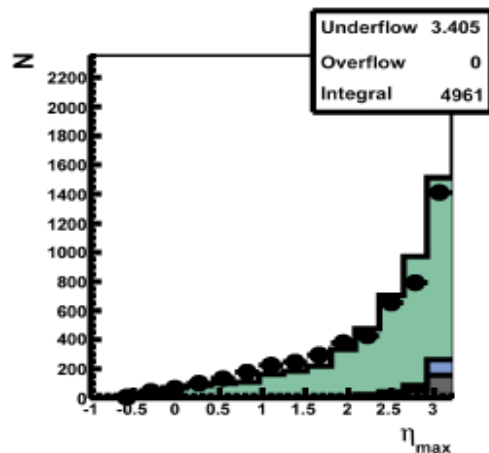
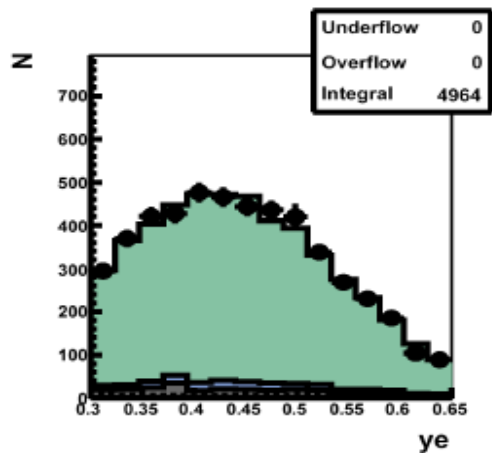
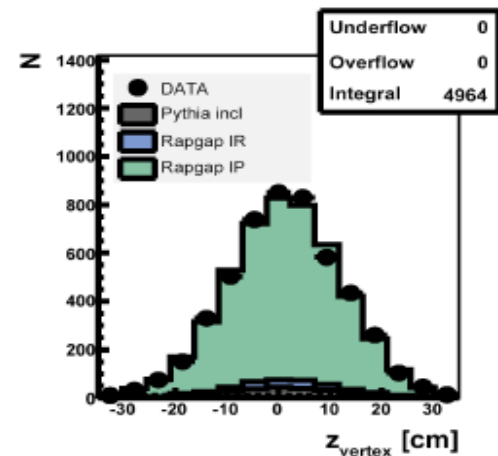
$$\left\{ \begin{array}{l} 0.3 < y_e < 0.65 \\ Q^2 < 0.01 \text{ GeV}^2 \\ |t| < 1 \text{ GeV}^2 \\ M_Y < 1.6 \text{ GeV} \end{array} \right.$$

These cuts define the cross section measurement phase space 11

# Control Plots

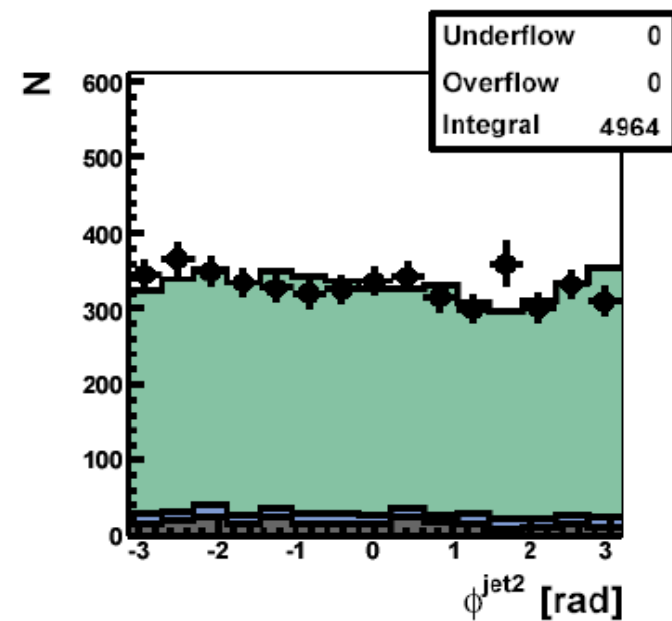
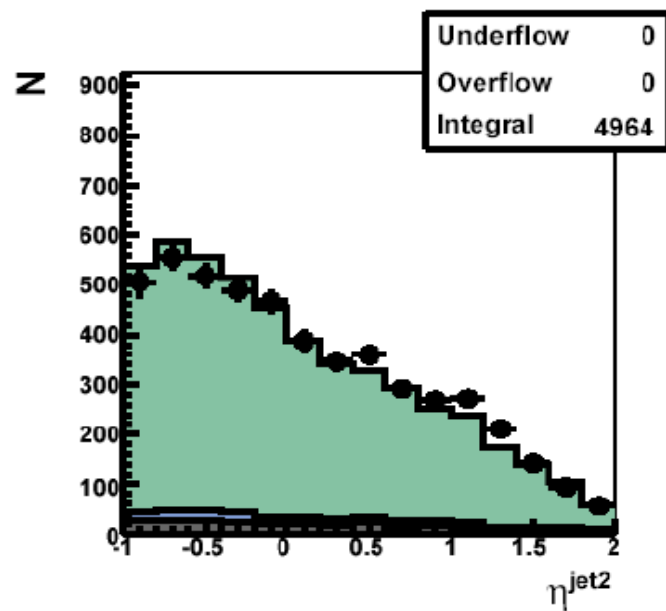
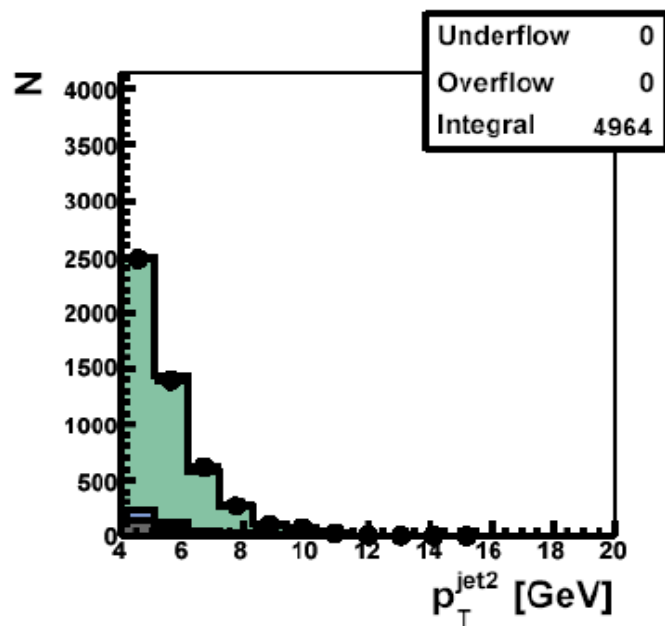
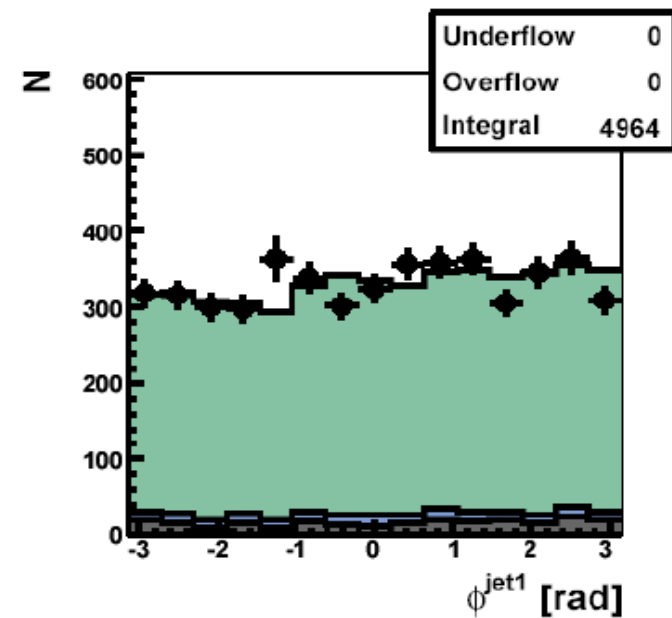
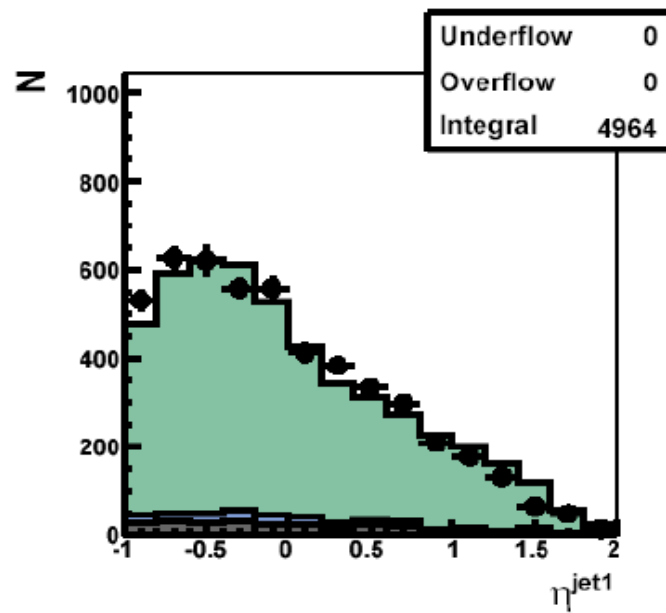
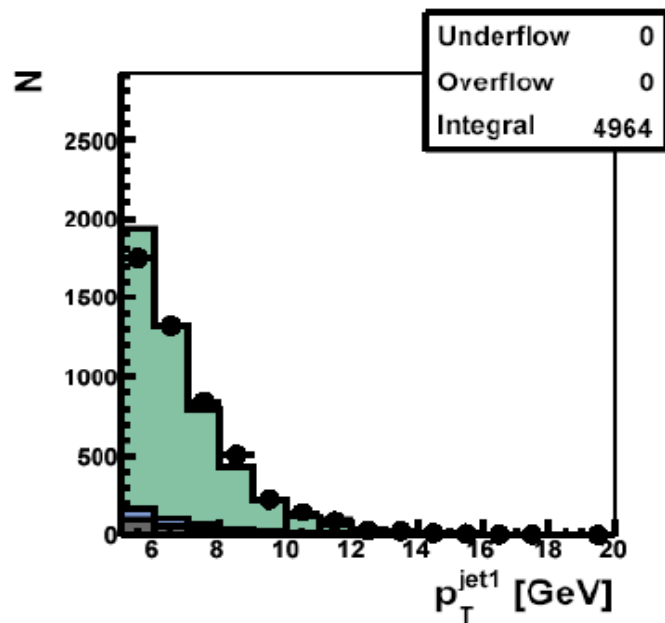
- Monte Carlo (MC) samples are generated and simulated
- Good description of the reconstructed level is needed.
- Rapgap 3.1 MC generator is used for signal events.
- Pythia 6.1 is used for background estimation.

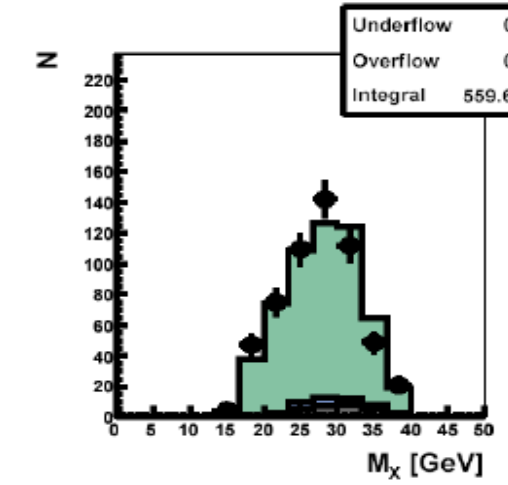
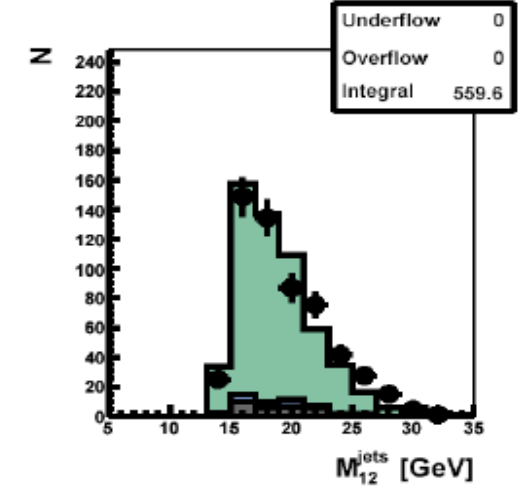
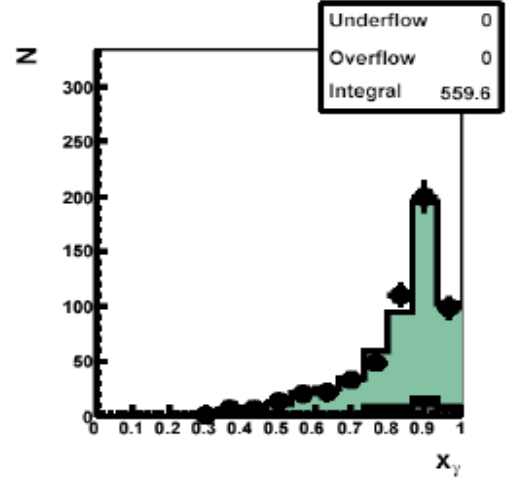
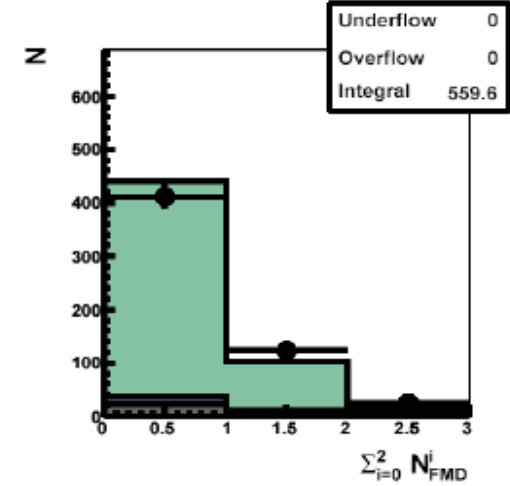
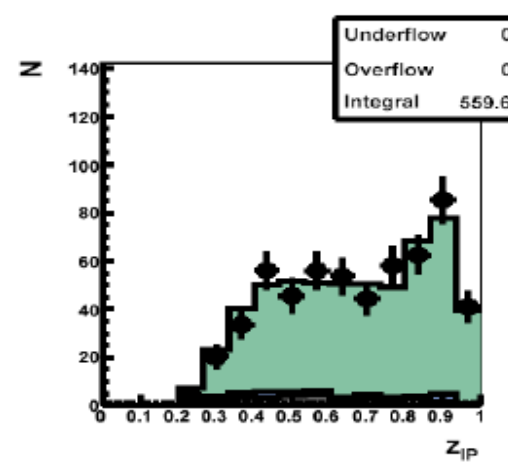
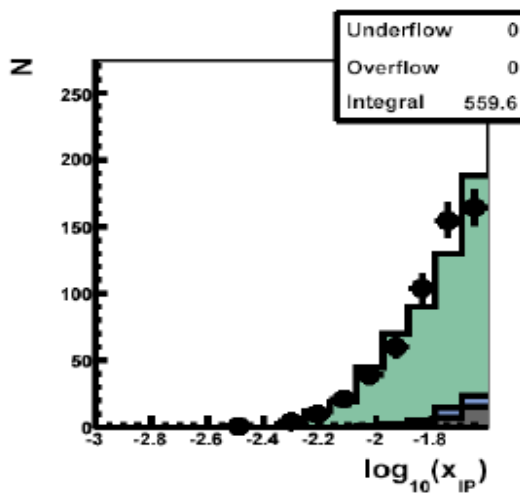
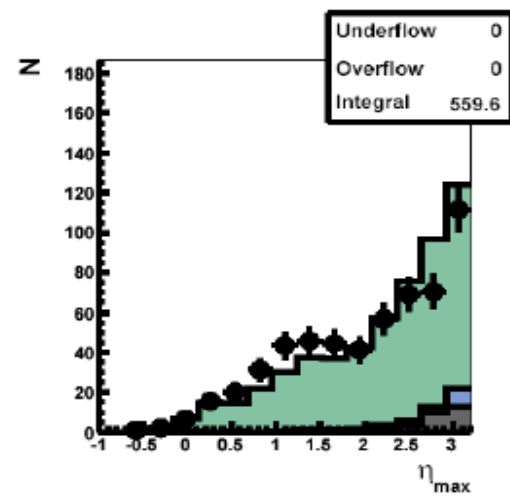
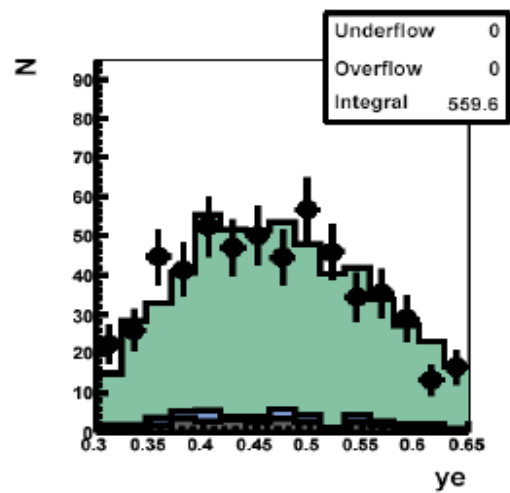
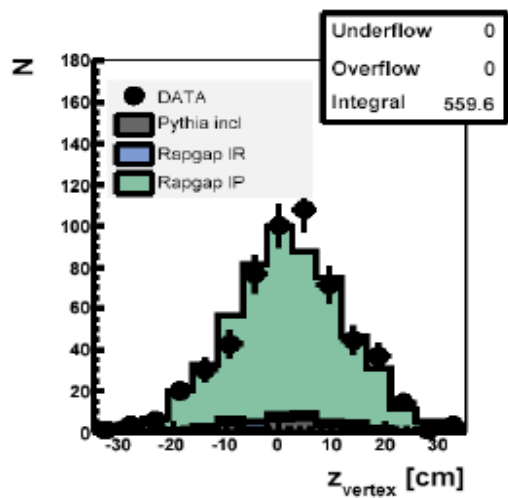




low  $E_T$

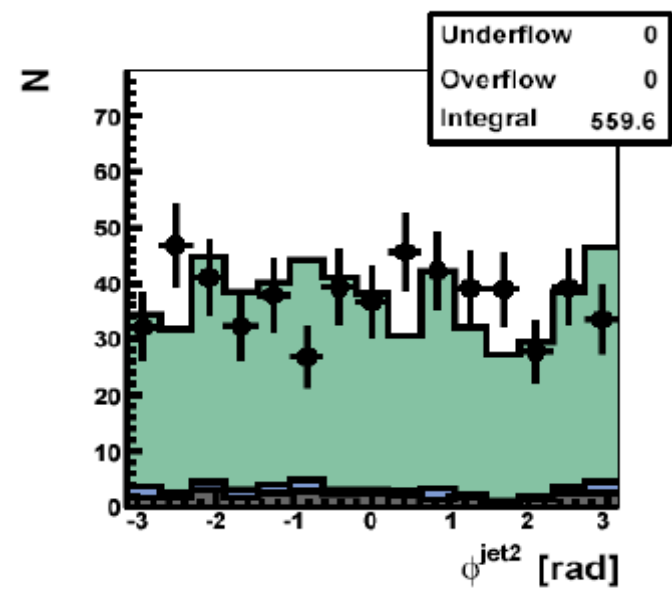
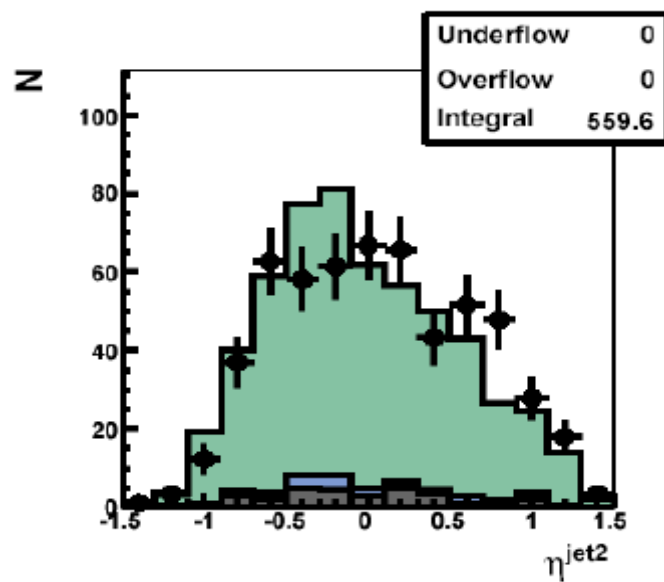
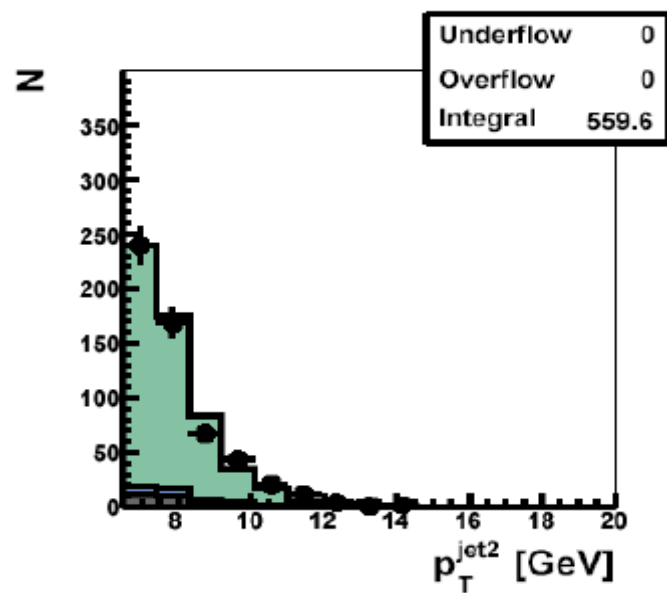
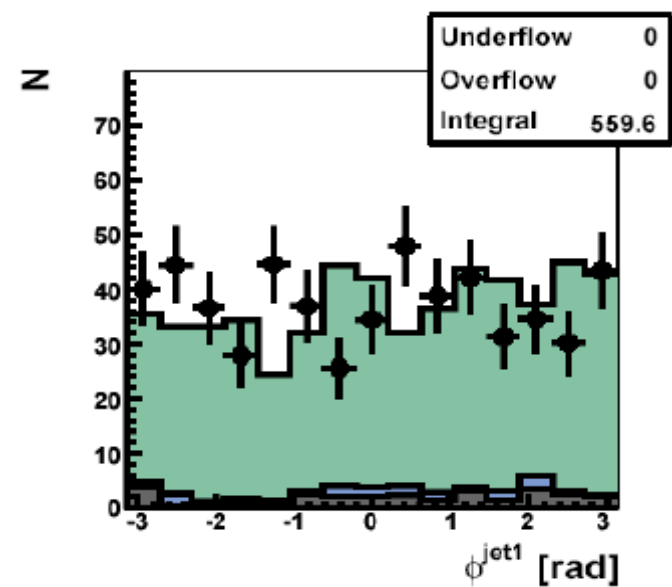
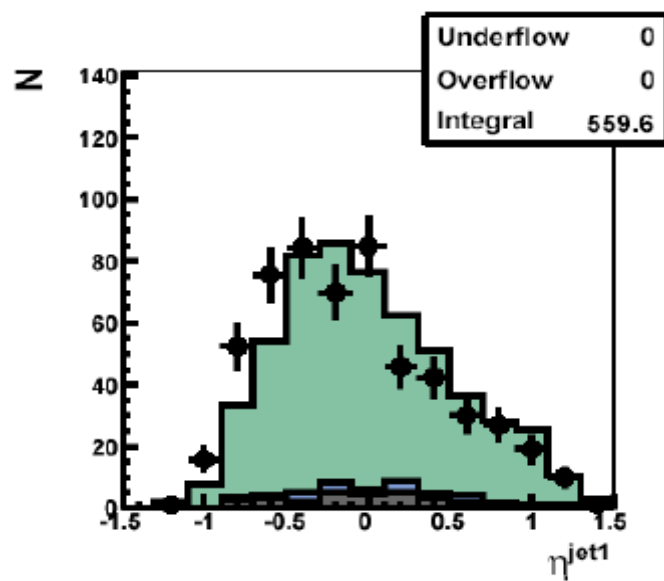
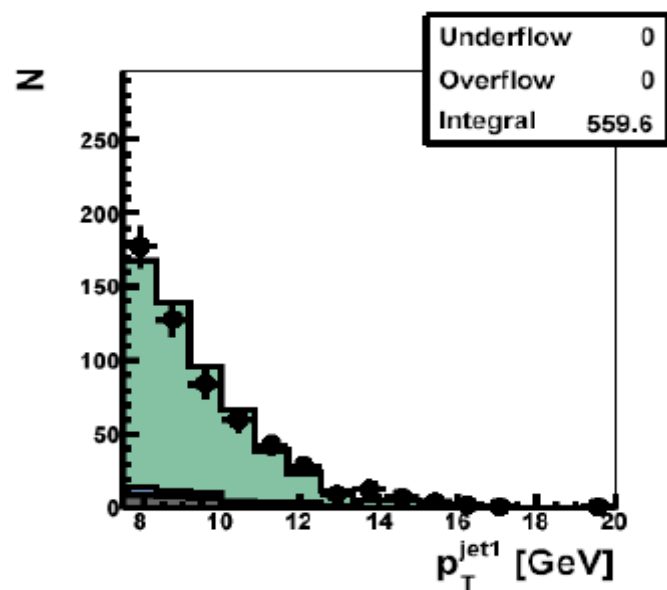
low  $E_T$





high  $E_T$

# high $E_T$





# Cross Section Measurement

- In each bin of given variable cross section is measured by means of following formula:

$$\frac{d\sigma}{dx_i} = \frac{N^{data} / \epsilon^{trigg} - N^{bgd.}}{A \cdot L \cdot \Delta_{x_i bin}}$$

Diagram illustrating the formula for cross section measurement, with labels and arrows pointing to the corresponding terms:

- $N^{data}$ : measured events
- $\epsilon^{trigg}$ : trigger eff.
- $N^{bgd.}$ : background events, from MC
- $A$ : acceptance
- $L$ : luminosity
- $\Delta_{x_i bin}$ : bin-width

- Eventually, a factor [ $C^{pdiss} = 0.94 \pm 7\%$ (syst.)] is applied accounting on proton dissociation correction into:  $M_Y < 1.6 \text{ GeV}$  and  $|t| < 1 \text{ GeV}^2$ .

# Systematic Uncertainties

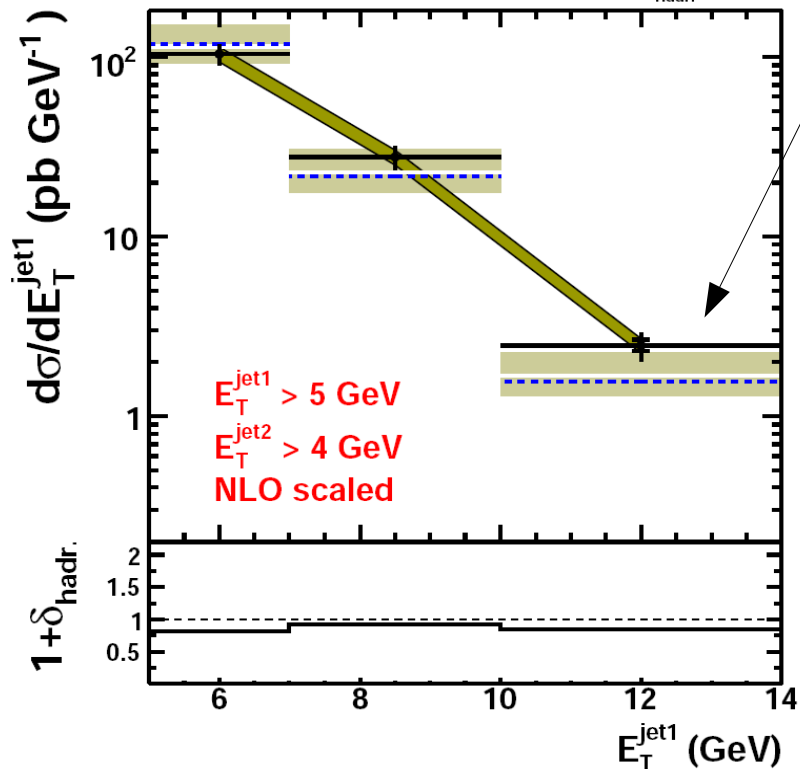
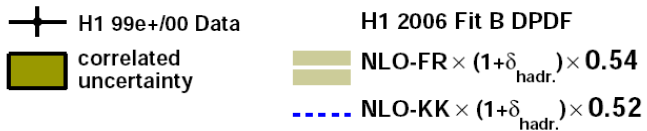
SOURCE	ELEM. SHIFT	$\Delta\sigma$ low $E_T$	$\Delta\sigma$ high $E_T$
HFS	2.5%	6.8%	12.3%
LRG	30.0%	8.1%	6.4%
TE	5.0%	5.0%	5.0%
Lumi	2.0%	2.0%	2.0%
p-diss	7.0%	7.0%	7.0%
FMD noise	0.6%	0.8%	0.8%
etag A	1.5%	1.5%	1.5%
bgd.	50.0%	2.7%	2.9%
$y^{\text{gen}}$ -shape	0.3	0.2%	0.9%
$ t^{\text{gen}} $ -shape	2	4.2%	3.2%
$p_T^{\text{hadr. jet1}}$ -shape	0.4 (0.8 for high $p_T$ )	1.0%	1.3%
$x_{\text{IP}}^{\text{hadr}}$ -shape	0.2 (0.4 for high $p_T$ )	5.0%	9.7%
$z_{\text{IP}}^{\text{hadr}}$ -shape	0.3	3.0%	5.2%
	<b>total =</b>	<b>~16%</b>	<b>~20%</b>

# Results

- Next-to-leading order (NLO) QCD calculations are used to predict the measured cross sections in both analyses.
- NLO program of Frixione et al. (FR) is used.
- NLO program of Kramer-Klasen (KK) is used.
- Both NLOs use the H1 2006 fit B parton distr. functions.
- Hadronization corrections are applied to the NLOs.

$$(1 + \delta_{had})_i = \left( \frac{\sigma_{dijet}^{hadron}}{\sigma_{dijet}^{parton}} \right)_i$$

# H1 PRELIMINARY

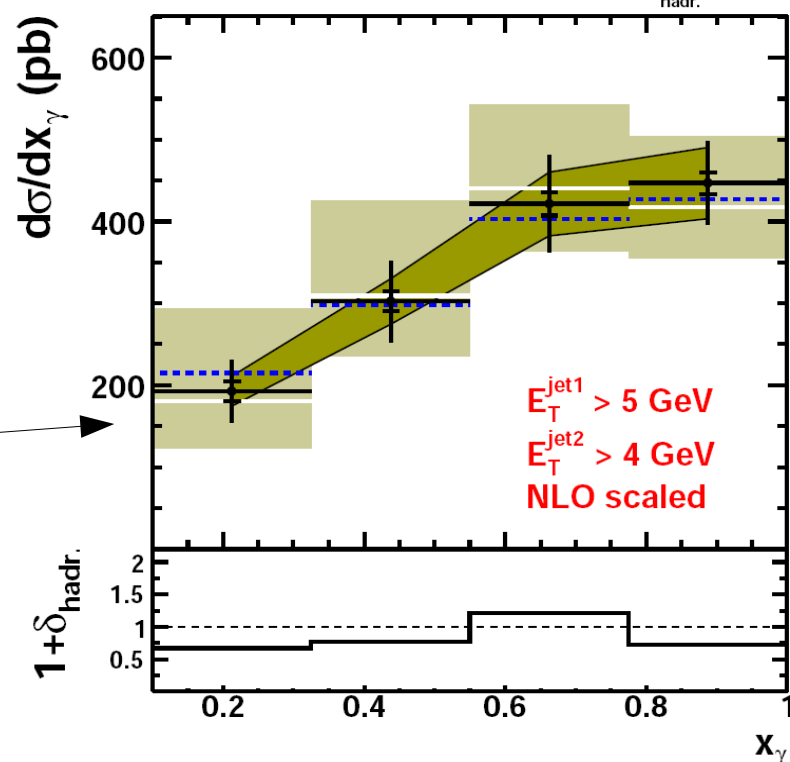
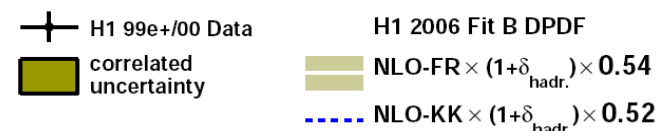


**low  $E_T$**

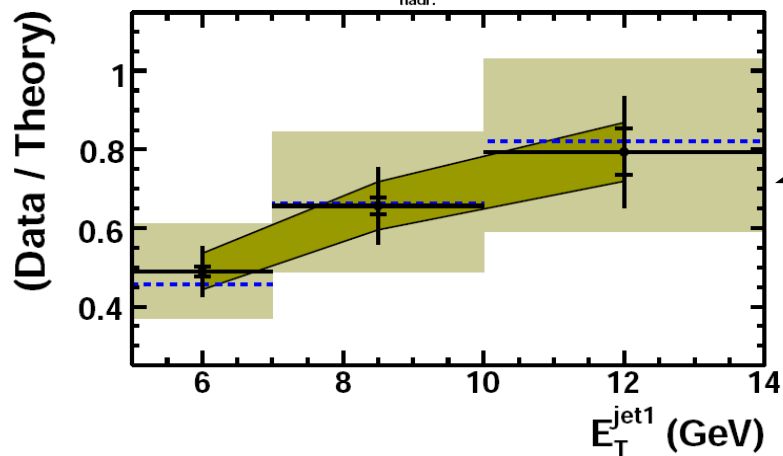
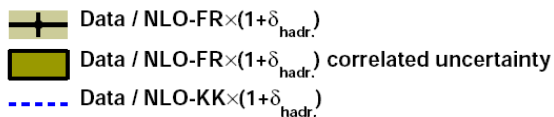
harder  $E_T$   
slope of  
data than  
NLO

good  
descripti  
on of  $x_\gamma$   
shape

# H1 PRELIMINARY



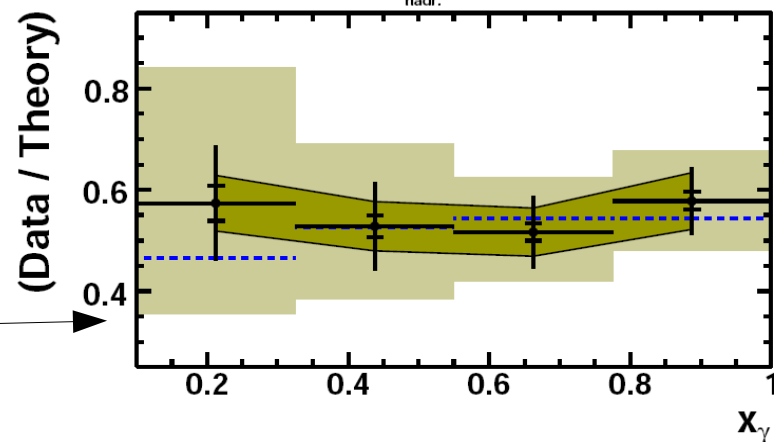
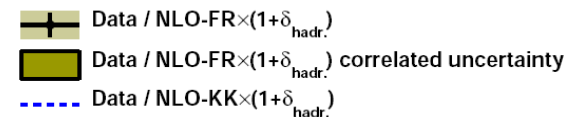
## H1 PRELIMINARY - Data / Theory



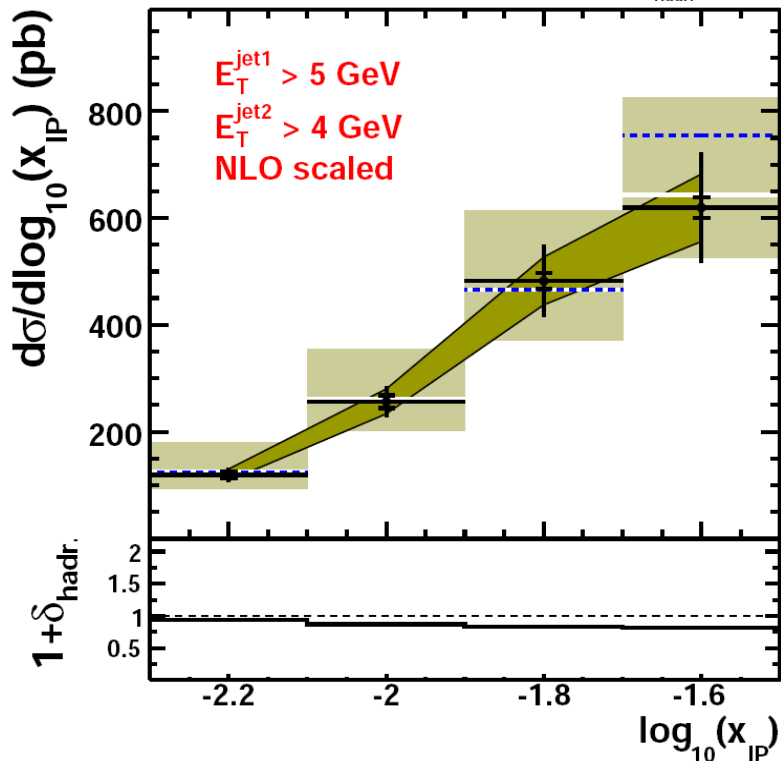
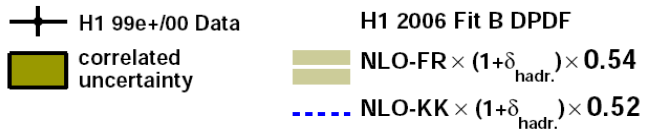
obvious  $E_T$   
dependence

flat in  $x_\gamma$

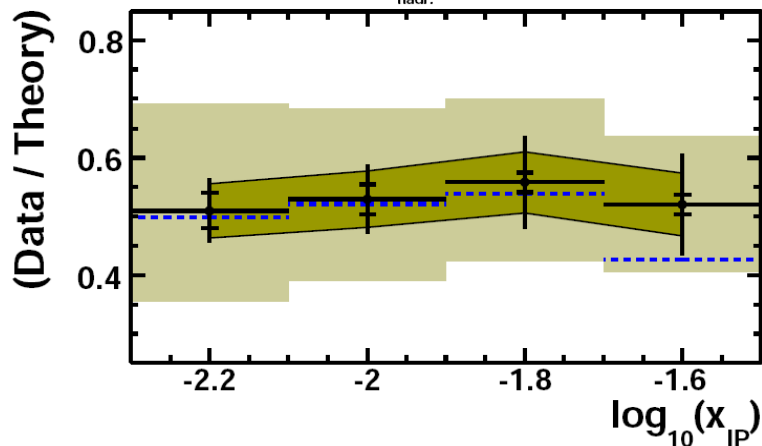
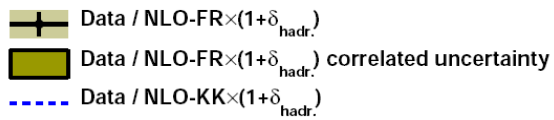
## H1 PRELIMINARY - Data / Theory



# H1 PRELIMINARY

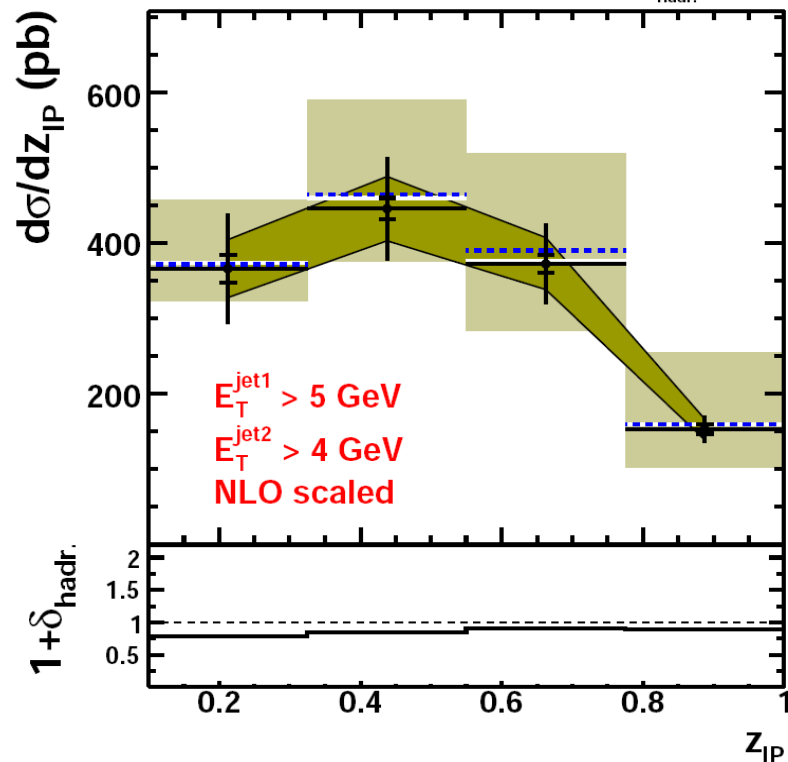
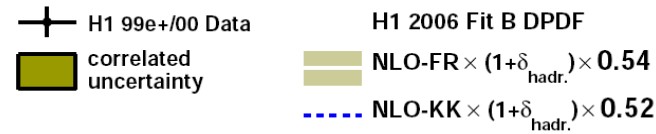


## H1 PRELIMINARY - Data / Theory

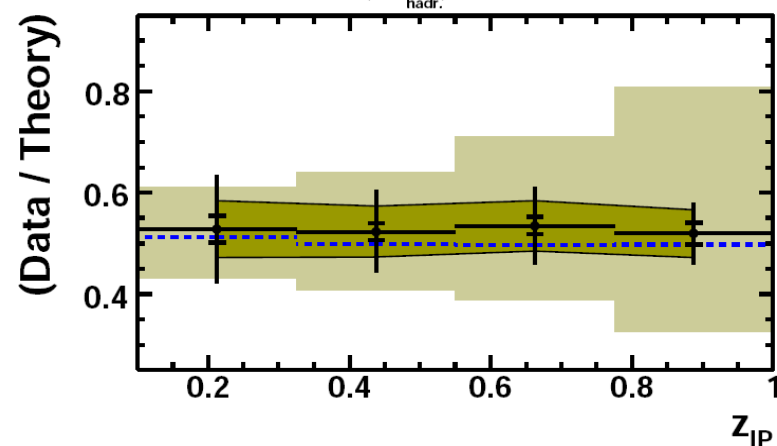
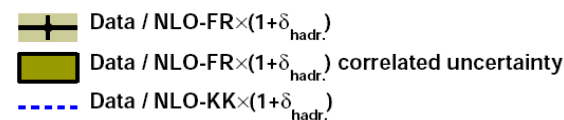


low  $E_T$

# H1 PRELIMINARY

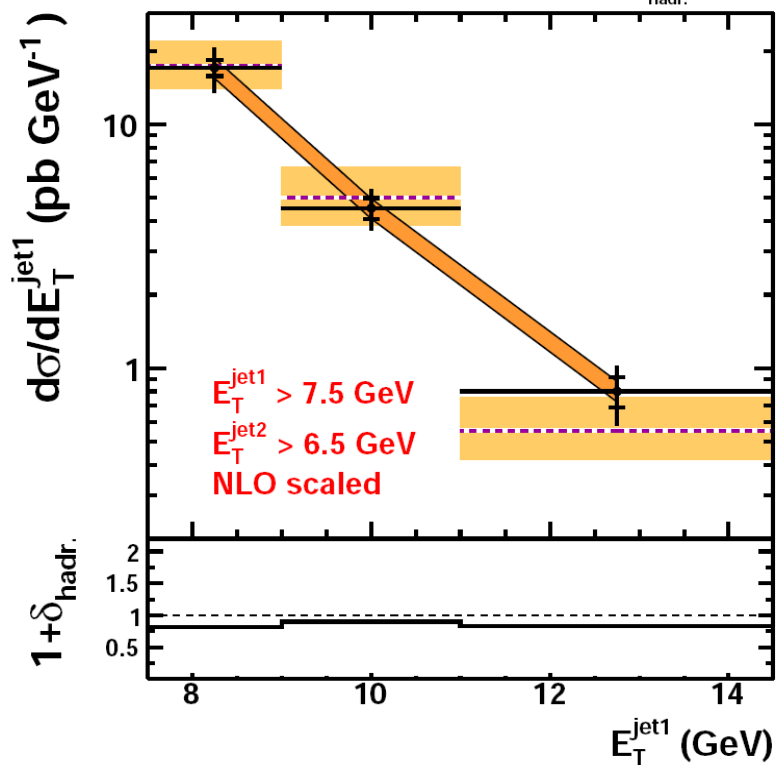


## H1 PRELIMINARY - Data / Theory



# H1 PRELIMINARY

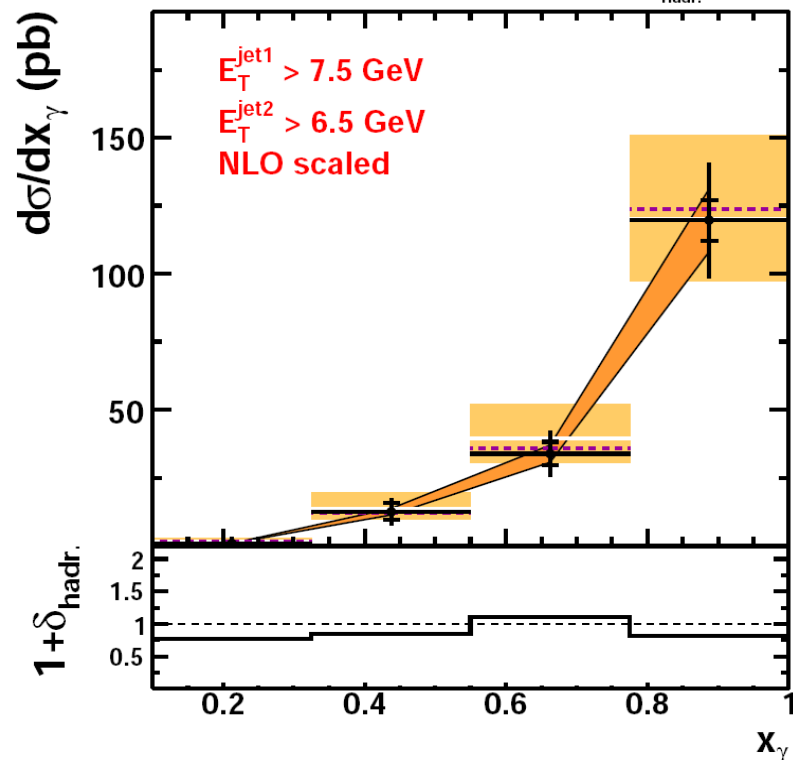
+ H1 99e+/00 Data  
 H1 2006 Fit B DPDF  
 correlated uncertainty  
 NLO-FR  $\times (1 + \delta_{\text{hadr.}}) \times 0.62$   
 NLO-KK  $\times (1 + \delta_{\text{hadr.}}) \times 0.63$



high  $E_T$

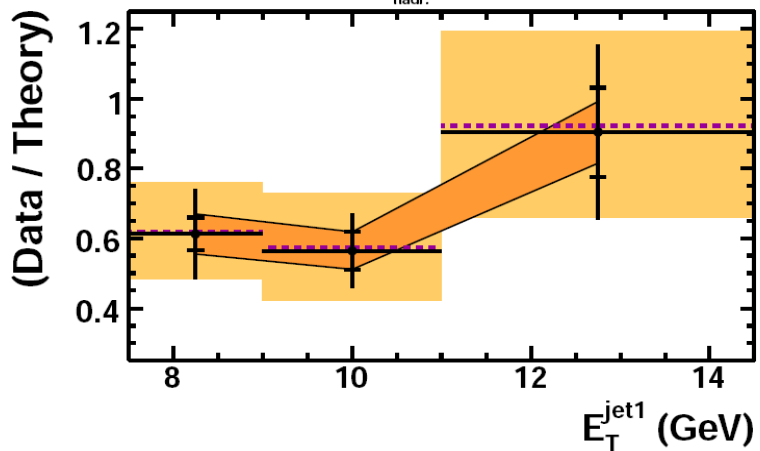
# H1 PRELIMINARY

+ H1 99e+/00 Data  
 H1 2006 Fit B DPDF  
 correlated uncertainty  
 NLO-FR  $\times (1 + \delta_{\text{hadr.}}) \times 0.62$   
 NLO-KK  $\times (1 + \delta_{\text{hadr.}}) \times 0.63$



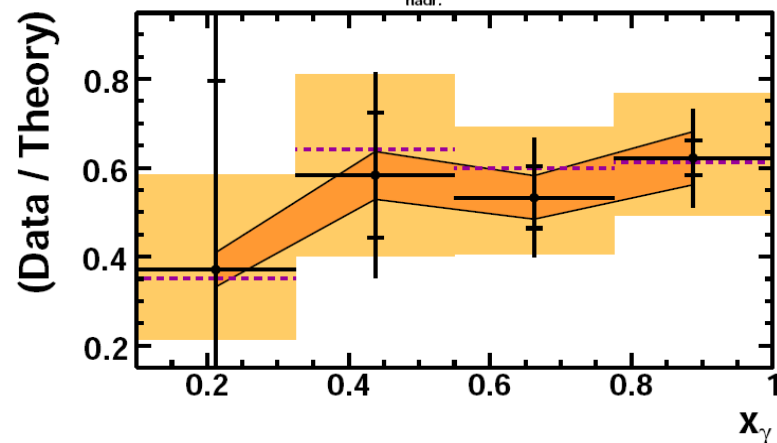
## H1 PRELIMINARY - Data / Theory

+ Data / NLO-FR  $\times (1 + \delta_{\text{hadr.}})$   
 Data / NLO-FR  $\times (1 + \delta_{\text{hadr.}})$  correlated uncertainty  
 Data / NLO-KK  $\times (1 + \delta_{\text{hadr.}})$

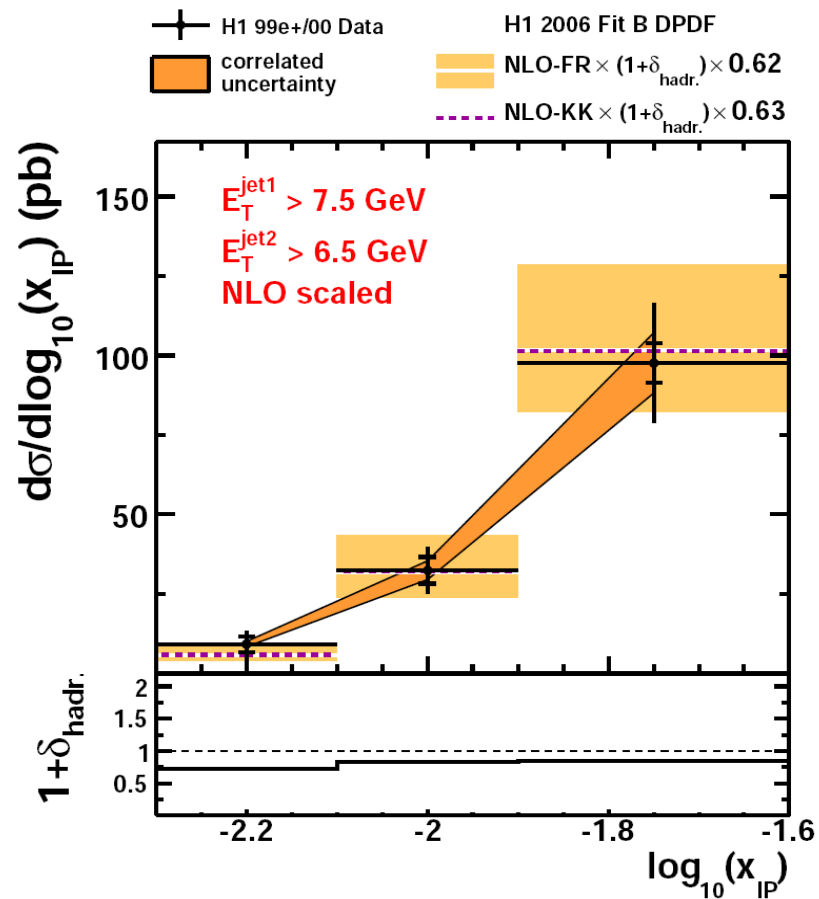


## H1 PRELIMINARY - Data / Theory

+ Data / NLO-FR  $\times (1 + \delta_{\text{hadr.}})$   
 Data / NLO-FR  $\times (1 + \delta_{\text{hadr.}})$  correlated uncertainty  
 Data / NLO-KK  $\times (1 + \delta_{\text{hadr.}})$

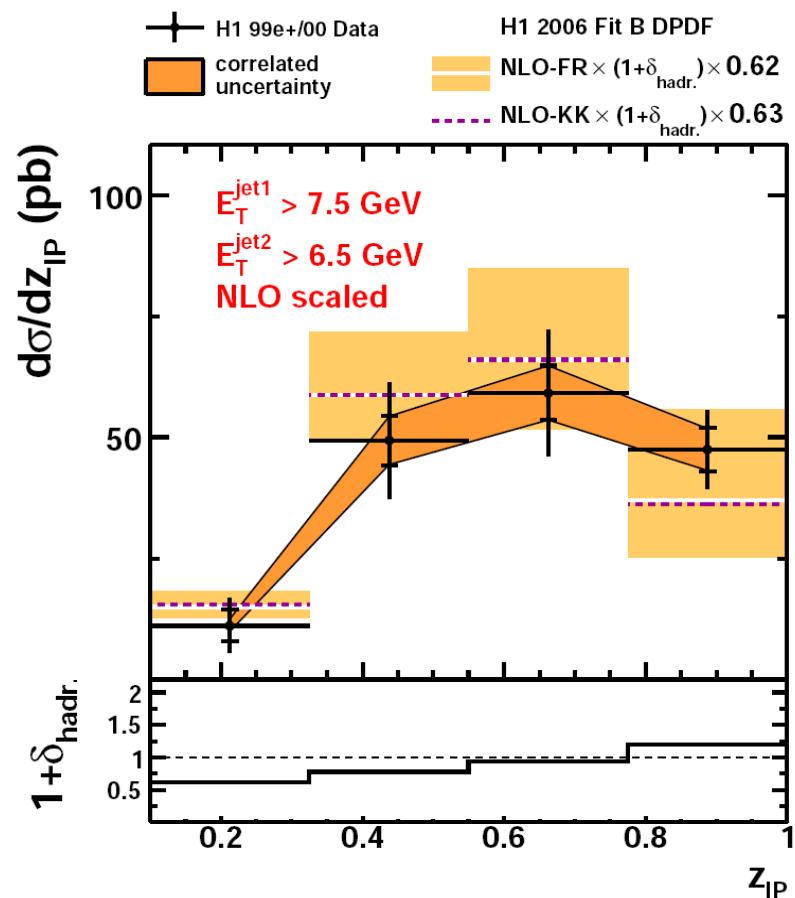


# H1 PRELIMINARY

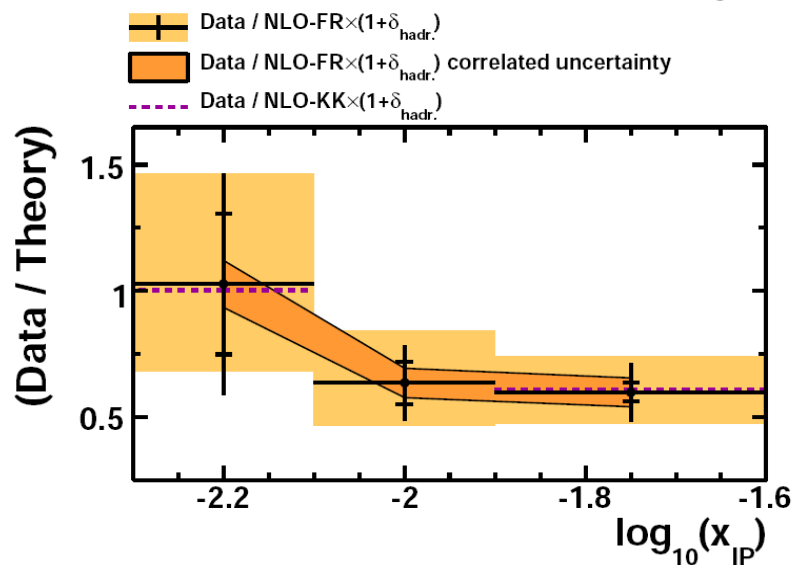


high  $E_T$

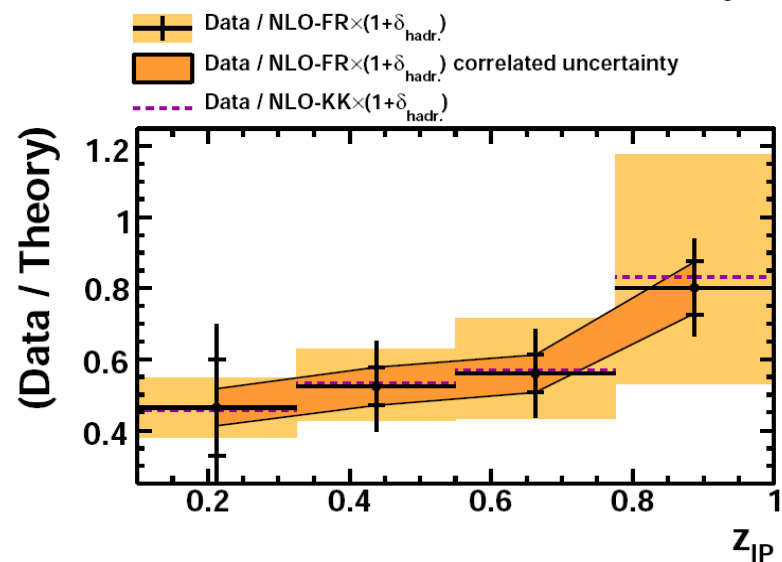
# H1 PRELIMINARY



## H1 PRELIMINARY - Data / Theory



## H1 PRELIMINARY - Data / Theory



- The suppression factors applied to NLO predictions are calculated from the total data and NLO cross sections.

**low  $E_T$**

$$f^{FR} = 0.54 \pm 0.01 (stat.) \pm 0.09 (syst.) {}^{+0.14}_{-0.13} (scale)$$
$$f^{KK} = 0.52 \pm 0.01 (stat.) \pm 0.08 (syst.)$$

**high  $E_T$**

$$f^{FR} = 0.62 \pm 0.04 (stat.) \pm 0.13 (syst.) {}^{+0.16}_{-0.14} (scale)$$
$$f^{KK} = 0.63 \pm 0.03 (stat.) \pm 0.13 (syst.)$$



# Conclusions

- The **NLO predictions overestimate measured cross section** in both analyses – stronger suppression in the low  $E_T$  analysis.
- **$E_T$  dependence** of the suppression **observed** in the low  $E_T$  analysis.
- With respect to the larger uncertainties the  **$E_T$  dependence** of the suppression in the **high  $E_T$**  analysis is **not excluded**.
- New analyses might illuminate the issue:
  - Photoproduction:  $\sigma_{dijets}^{diff.} / \sigma_{dijets}^{incl.}$
  - Photoproduction & DIS:  $\frac{\sigma_{dijets}^{Php, data} / \sigma_{dijets}^{Php, NLO}}{\sigma_{dijets}^{DIS, data} / \sigma_{dijets}^{DIS, NLO}}$

# Backup – cuts

!! the analysis is carried out in 2 cut schemes !!

- 1) cuts identical with previous H1 measurement
- 2) ZEUS-like analysis cuts (harder jets,...)

## detector level cuts

	<b>low <math>E_T</math></b>	<b>high <math>E_T</math></b>
scattered electron cuts	$\left\{ \begin{array}{l} 0.3 < y_e < 0.65 \\  X_{clus}^{etag}  < 6.5 \text{ cm} \\ E_{photon det.} < 2 \text{ GeV} \end{array} \right.$	
jet cuts	$\left\{ \begin{array}{l} p_T^{jet1} > 5 \text{ GeV} \\ p_T^{jet2} > 4 \text{ GeV} \\ -1 < \eta^{jet1 \text{ and } jet2} < 2 \end{array} \right.$	$\left\{ \begin{array}{l} p_T^{jet1} > 7.5 \text{ GeV} \\ p_T^{jet2} > 6.5 \text{ GeV} \\ -1.5 < \eta^{jet1 \text{ and } jet2} < 1.5 \end{array} \right.$
diffr. cuts	$\left\{ \begin{array}{l} \eta^{max} < 3.2 \\ x_{IP} < 0.03 \\ \sum_{i=1}^5 N_i^{PRT} < 1 \\ \sum_{i=1}^3 N_i^{FMD} < 3 \wedge \sum_{i=1}^2 N_i^{FMD} < 2 \end{array} \right.$	$\left\{ \begin{array}{l} x_{IP} < 0.025 \end{array} \right.$

# Backup – MC samples

- Rapgap 3.1 IP and IR exchange with H1 2006 fit b for DPDFs:

	<u>IP</u>	<u>IR</u>
BGF-c	870k	110k
BGF-uds	1.82M	190k
QCD Compton-uds	330k	140k
resolved-uds	12M	2.5M

- *GRV-G LO* photon structure function
- Pythia 6.2 inclusive photoproduction sample for background

# Backup – reconstruction of kinematics

<u>rec. level</u>	<u>hadron. level</u> (GKI stable hadrons)	<u>parton. level</u> (GKI partons)
$y_e = 1 - \frac{E_{etag}}{E_e^{beam}}$	$y_e = 1 - \frac{E_e^{gen}}{E_e^{beam}}$	-//-
$W = \sqrt{s \cdot y_e}$	-//-	-//-
$x_y = \frac{\sum_{jets} E_i - p_{z,i}}{\sum_{HFS} E_i - p_{z,i}}$	-//-	-//-
$x_{IP} = \frac{\sum_{HFS} E_i + p_{z,i}}{2 \cdot E_p^{beam}}$	$x_{IP} = \frac{\sum_{hadrons}^{no\,elas.P} E_i + p_{z,i}}{2 \cdot E_p^{beam}}$	$x_{IP} = x_{IP}^{gen}$
$z_{IP} = \frac{\sum_{jets} E_i + p_{z,i}}{2 \cdot x_{IP} \cdot E_p^{beam}}$	-//-	-//-
$M_{12} = \sqrt{J_1^\mu \cdot J_{2,\nu}}$	-//-	-//-
$M_X = \sqrt{s \cdot y_e \cdot x_{IP}}$	$M_X = \sqrt{X^\mu \cdot X_\mu} \dots X^\mu = \sum_{hadrons}^{no\,elas.P} P_i^\mu$	$M_X = \sqrt{X^\mu \cdot X_\mu} \dots X^\mu = \sum_{partons} P_i^\mu$

# Backup – FR NLO settings

## resolved, IP

5 ! # of flavours  
4 1 ! Hadron type  
1001 ! PDF set for the hadron, negative  
for  
DPDFLib - 1150 = fit2  
5003 ! PDF set for 1. particle  
0.228E+ 00 ! Lambda\_5, 0 for default  
'MS' ! Scheme for hadron

## direct, IP:

5 ! # of flavours  
1 ! Hadron type  
1001 ! PDF set for the hadron, negative  
for  
DPDFLib - 1150 = fit2  
0.228E+ 00 ! Lambda\_5, 0 for default  
'MS' ! Scheme for hadron  
'DI' ! Scheme for photon