

Dynamical screening effects in X-ray absorption spectra of graphene and BN

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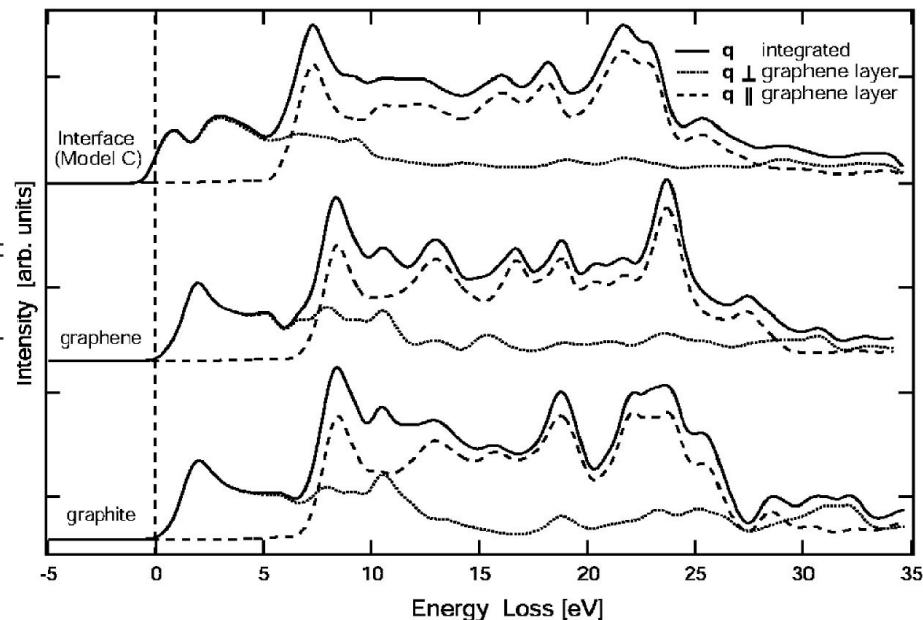
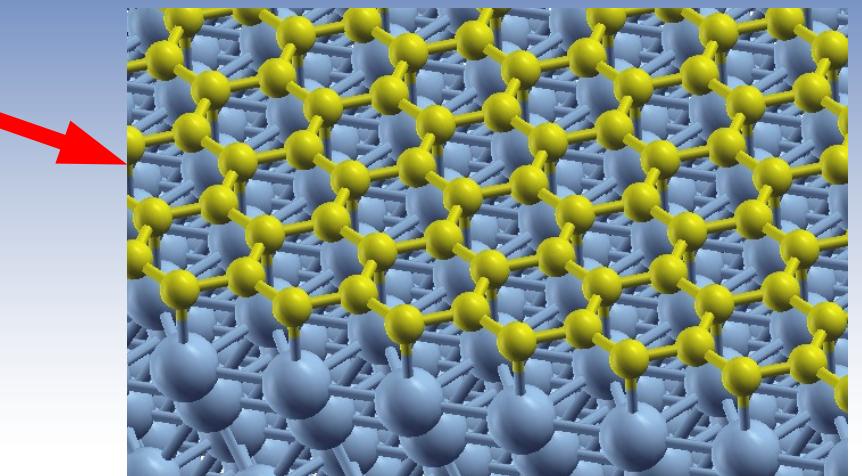
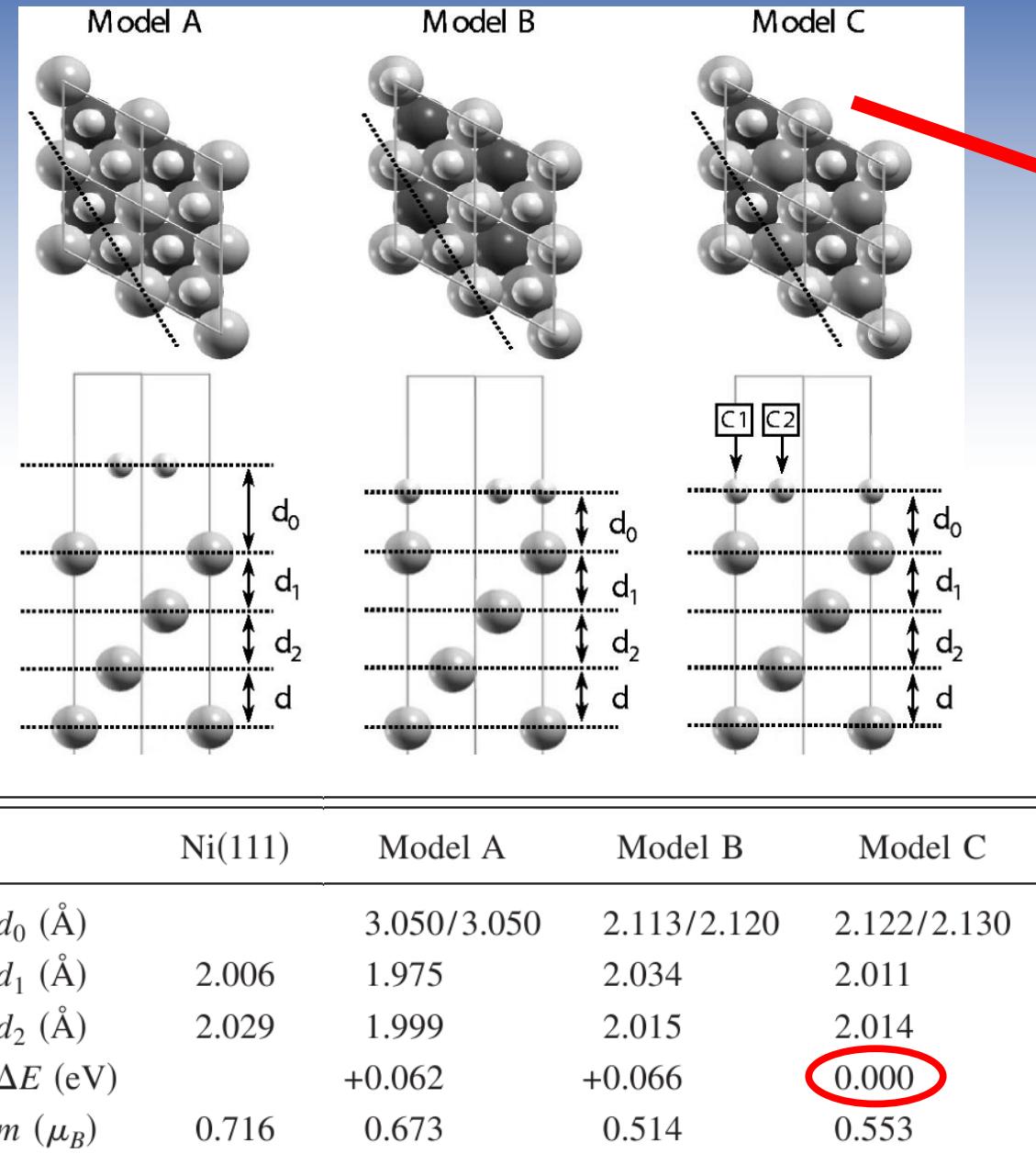
In collaboration with:

O. Wessely, May Ling Ng, N. A. Vinogradov, N. Mårtensson

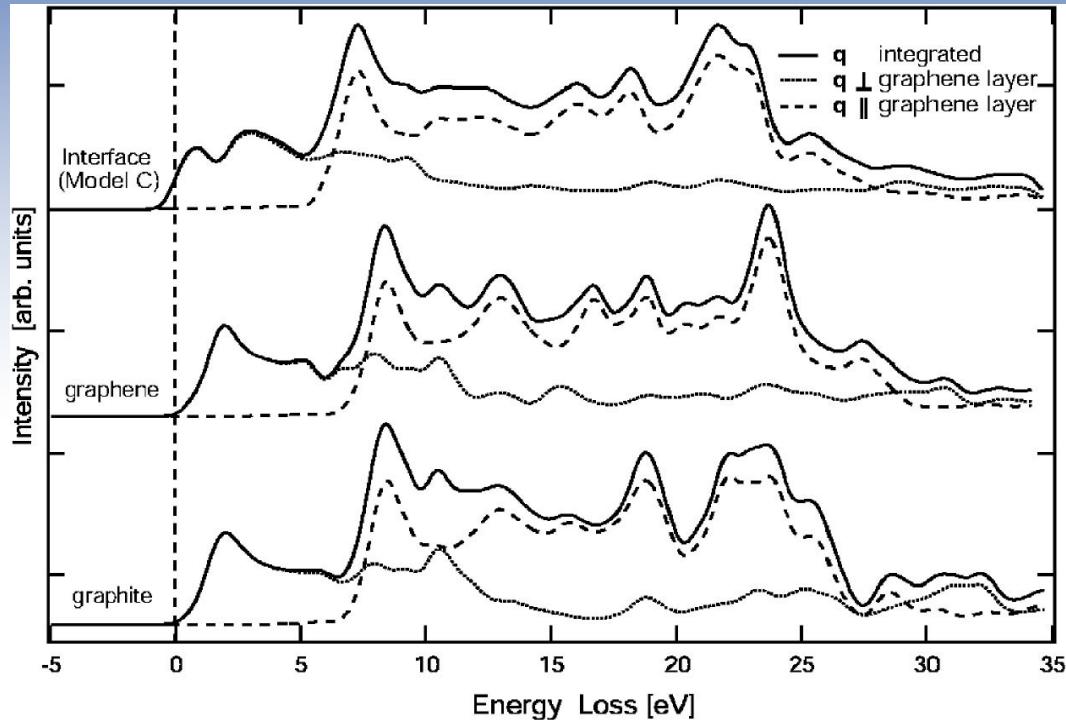
Outline

- ◆ Graphene on Ni(111)
- ◆ BN layer on Ni(111)
- ◆ MND theory
- ◆ XAS calculations using MND
- ◆ Results
 - ◆ MG/Ni(111)
 - ◆ hBN/Ni(111)
- ◆ Conclusions

Graphene on Ni(111)

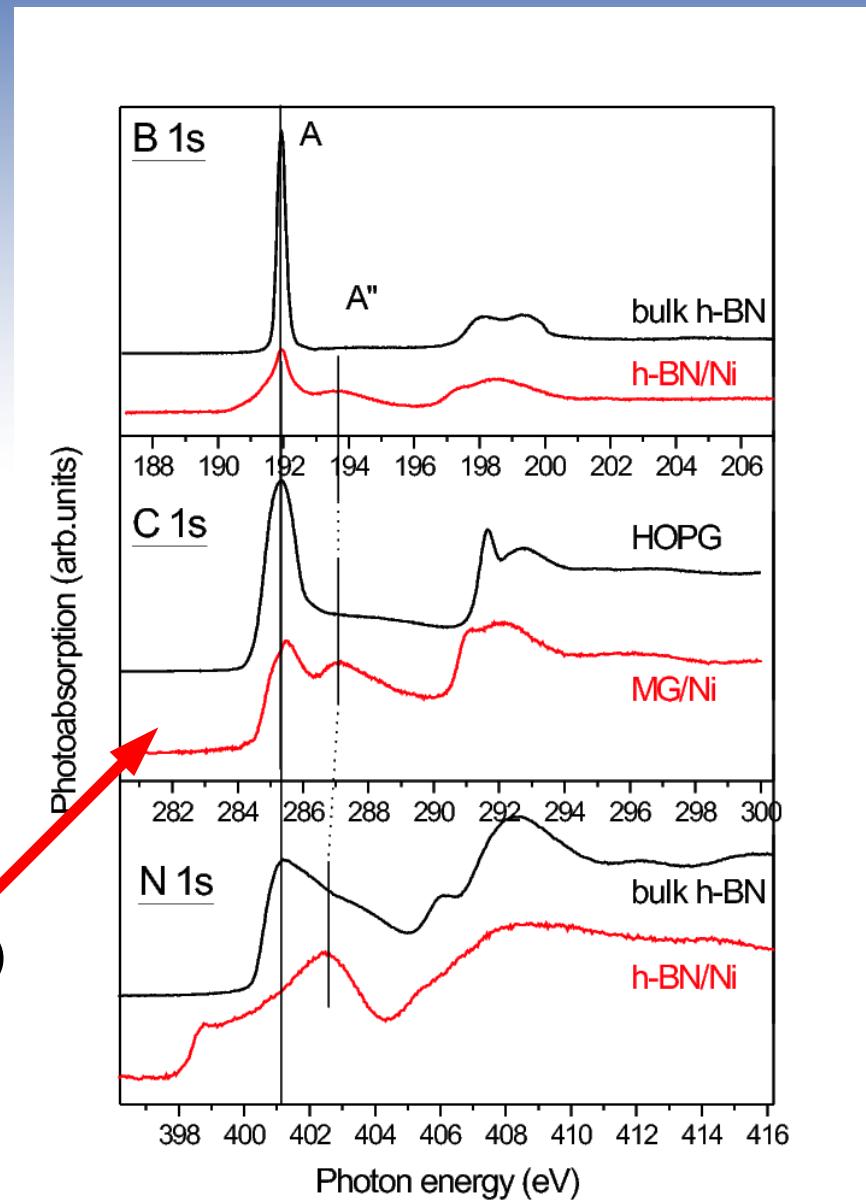


Graphene on Ni(111)

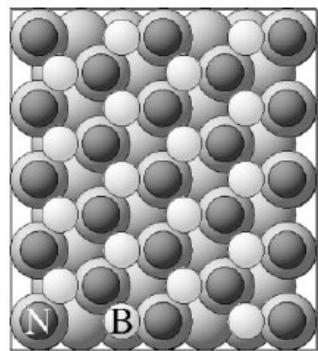


Ground state DFT calculation of ELNES

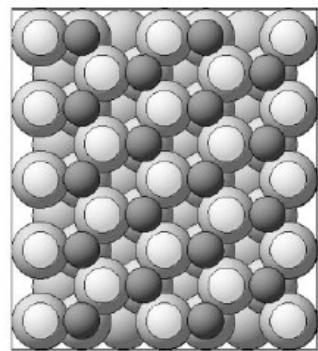
XAS experiment (A. Preobrajenski)



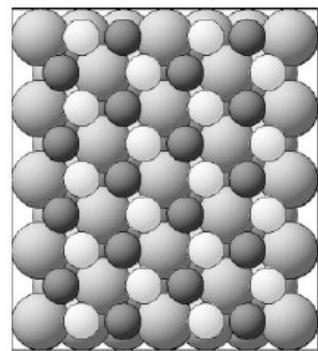
h-BN on Ni(111)



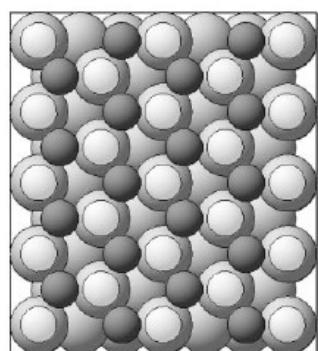
(N,B)=(top,fcc)



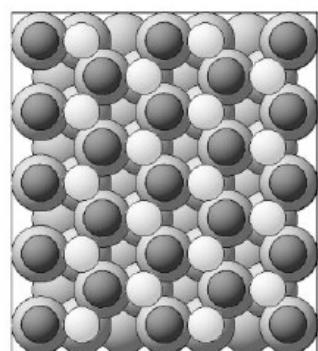
(N,B)=(hcp,top)



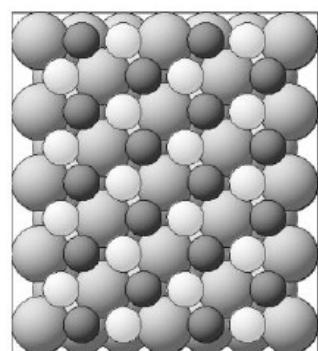
(N,B)=(fcc,hcp)



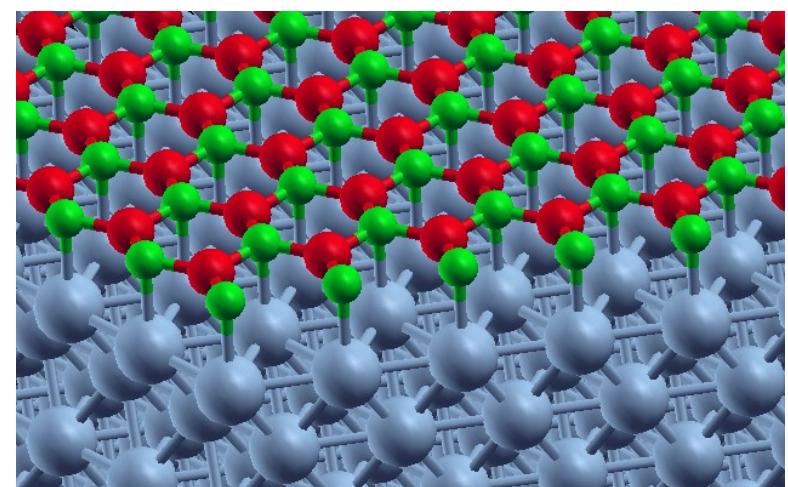
(N,B)=(fcc,top)



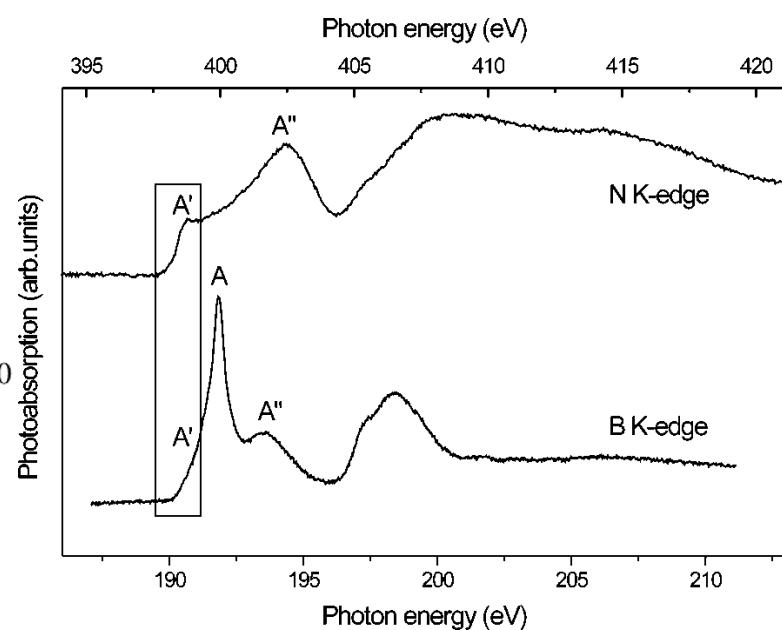
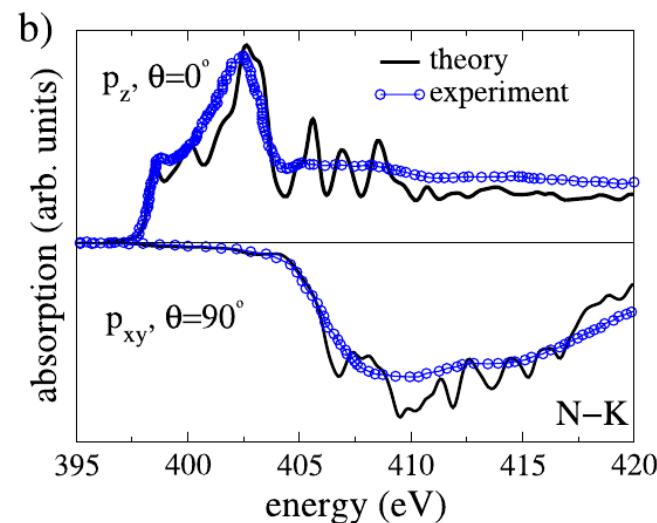
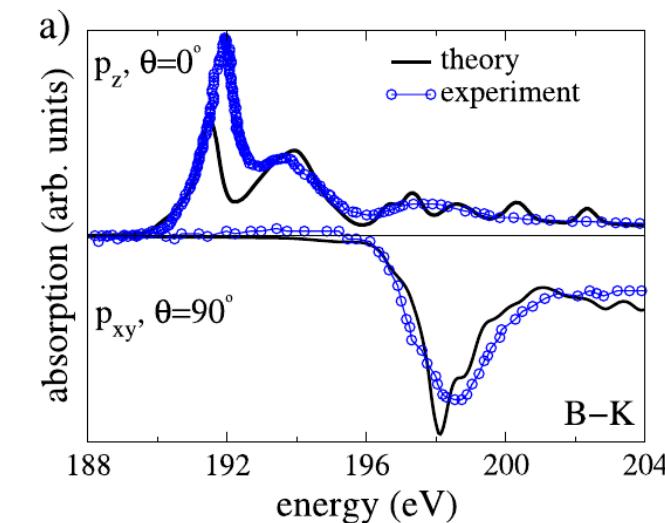
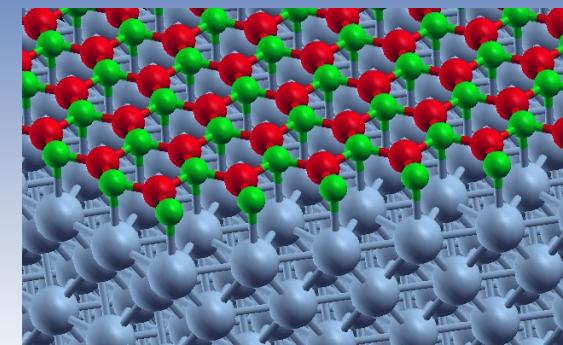
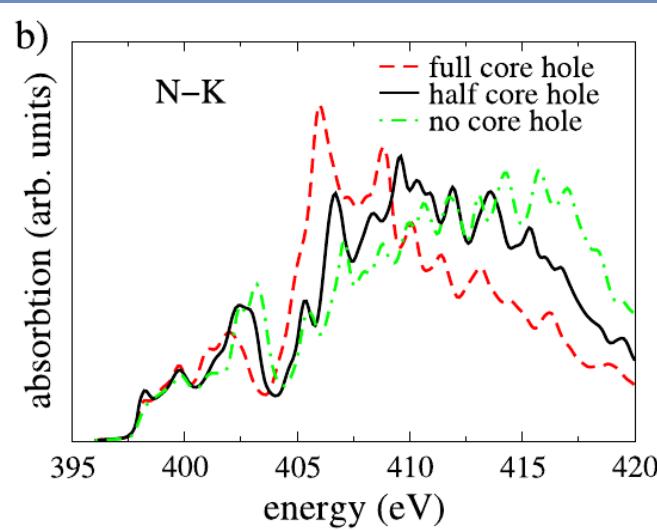
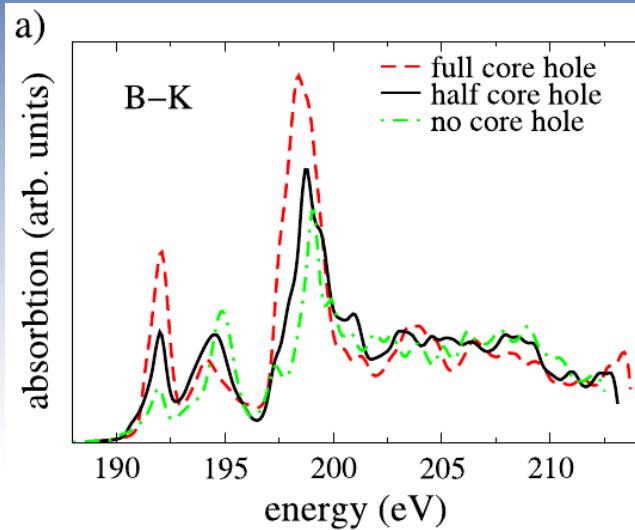
(N,B)=(top,hcp)



(N,B)=(hcp,fcc)



h-BN on Ni(111)



Preobrajenski et al., PRB70, 165404 (2004)
 Laskowski et al., JPCM21, 104210 (2009)

MND theory

- Model

$$H = \sum_{RR'} \varepsilon_{RR'} a_R^\dagger a_{R'} + E_0 b^\dagger b - u a_0^\dagger a_0 b_0^\dagger b_0^\dagger$$

- XAS spectrum

$$I(\omega) \sim |W|^2 \operatorname{Re} \int_{-\infty}^{\infty} dt K_{00}(t, t) e^{i\omega t}$$

- 2-particle Green's function

$$K_{RR'}(\tau, t) = \frac{1}{i\omega} \langle 0 | T a_R(\tau) b_0^\dagger(t) a_{R'}^\dagger(0) b_0(0) | 0 \rangle$$

Mahan, Phys. Rev. 163, 612 (1967)

Nozieres & de Dominicis, Phys. Rev. 178, 1097 (1969)

Grebennikov et al., Phys. Stat. Sol. (b) 79, 423 (1977)

MND implementation

- Calculate initial and final state Green's func.

- Solve

$$\tilde{\varphi}(\epsilon, t) = \mathbf{1} + \mathbf{V} \int_{\mu}^{\infty} \frac{d\epsilon_1}{\pi} \operatorname{Im} \mathbf{g}^R(\epsilon_1) \mathbf{L}^+(\epsilon, \epsilon_1, t) \tilde{\varphi}(\epsilon_1, t),$$

$$\mathbf{L}^+(\epsilon, \epsilon_1, t) = \frac{\mathbf{I}^*(\epsilon_1, t) - e^{i(\epsilon-\epsilon_1)t} \mathbf{I}^*(\epsilon, t)}{\epsilon_1 - \epsilon},$$

$$\mathbf{I}(\epsilon, t) = \mathbf{1} + \mathbf{V} \int_{-\infty}^{\infty} \frac{d\epsilon_1}{\pi} \frac{e^{i(\epsilon-\epsilon_1)t} - 1}{\epsilon - \epsilon_1} \operatorname{Im} \mathbf{G}^R(\epsilon_1).$$

- Evaluate

$$\Delta(t) = i \operatorname{Tr} \mathbf{V} \int_0^{-t} dt_1 \int_{-\infty}^{\mu} \frac{d\epsilon}{\pi} \operatorname{Im} \mathbf{g}^R(\epsilon) \tilde{\varphi}(\epsilon, t_1)$$

- And finally

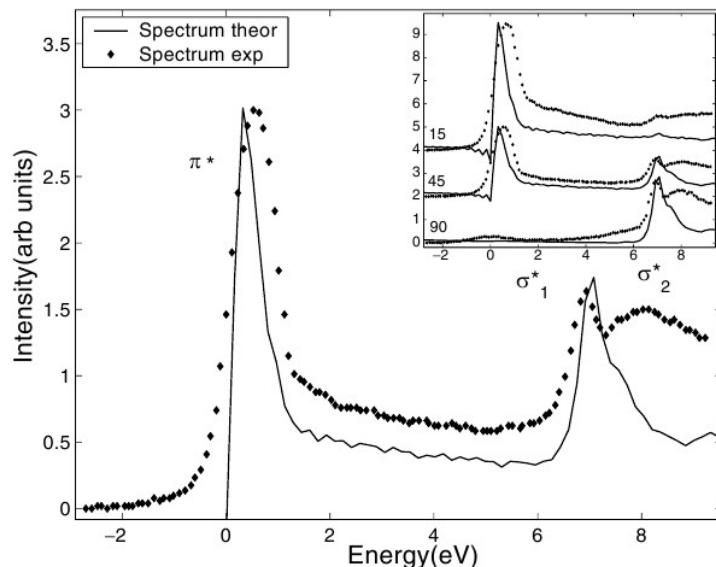
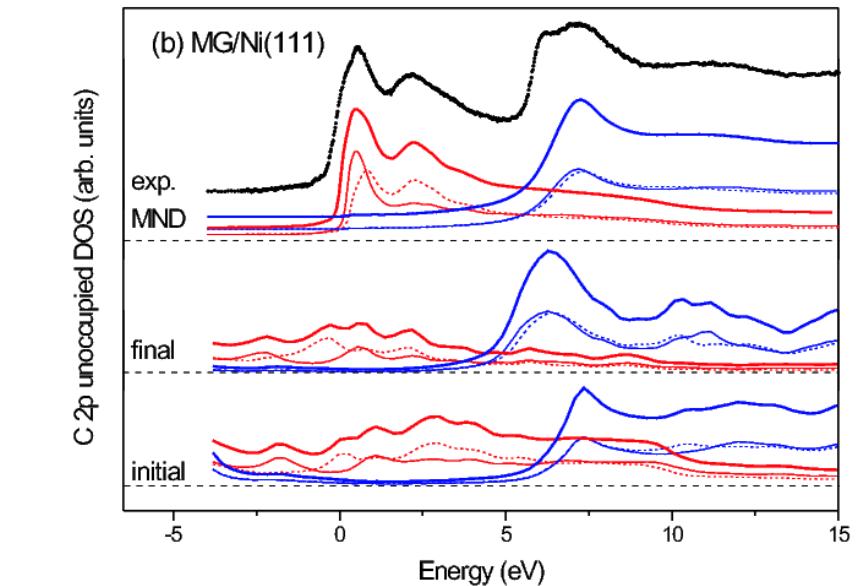
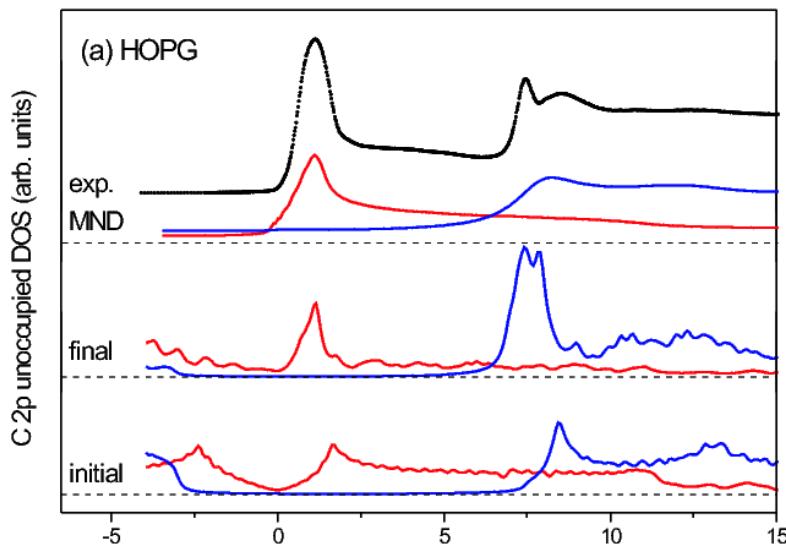
$$I(\omega) \sim \operatorname{Re} \int_0^{\infty} dt \int_{\mu}^{\infty} \frac{d\epsilon}{\pi} e^{it(\omega-\epsilon+\epsilon_c)+\Delta(-t)}$$

$$\times \sum_{ij} t_i [\operatorname{Im} \mathbf{g}^R(\epsilon) \tilde{\varphi}(\epsilon, t)]_{ij} t_j^*,$$

Computational approach

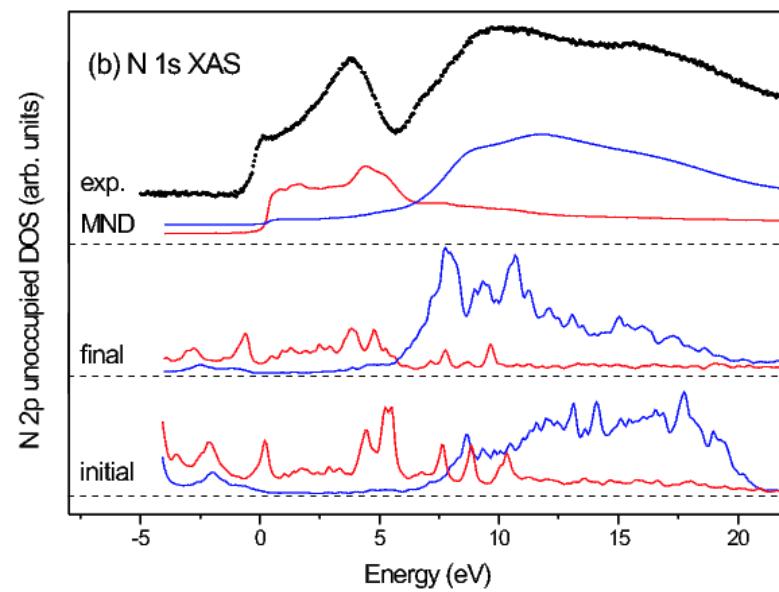
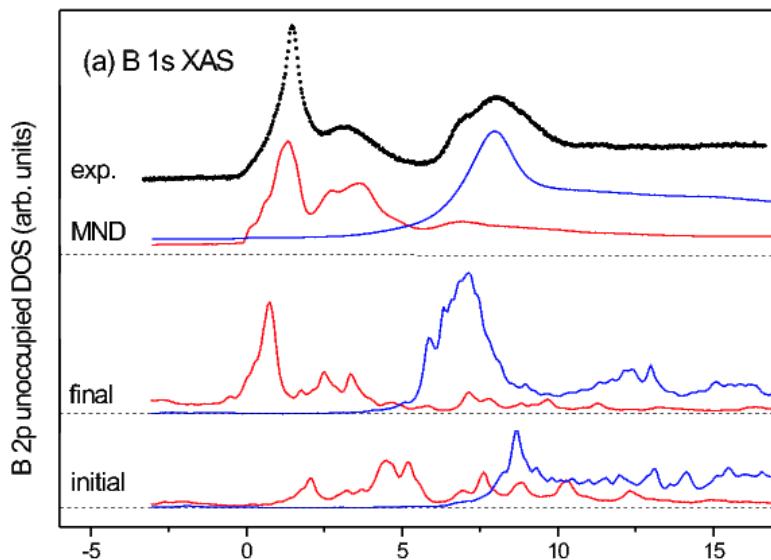
- Optimization of atom positions with VASP
- 6 layers of Ni, vacuum layer of 18Å
- WIEN2k calculations of the electronic structure
 - initial state
 - final state – supercell 3x3x1 with a core hole
- Evaluation of the local Green's function matrices
- OpenMP implementation of MND equations in fortran

MND spectra of graphene



Wessely et al., PRL94, 167401 (2005)

MND spectra of h-BN



Conclusions

- Parameter-free calculations of the XAS
- Importance of the dynamical core-hole screening effects
- Good agreement with experiment for all B, C and N K-edge spectra
- Improvement over initial state, final state or Slater transition state XAS calculations