

SMALL DISTURBANCES WAVES IN BOUNDARY LAYER INDUCED BY ONCOMING WAKES

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Introduction

The flow in the fluid-flow machinery is inherently three-dimensional and unsteady mostly due to wakes from preceding blade rings. It makes earlier inception of the laminar-turbulent transition, what affects the friction on the surface of the blade and heat transfer between flow and blades.

The wake can be identified differently depending on the choice of co-ordinate system: as a velocity defect in the absolute system of co-ordinates or as a jet in the system of co-ordinates fixed with the mean velocity of flow. In the turbomachinery the blades cut the wake into segments. The boundary layer on the blade is affected by the positive jet when the speed of the jet is orientated towards the surface of the blade so the jet impinges the blade and affected by the negative jet if the velocity of the jet is directed from the blade towards the external flow. Meyer [2] using potential flow solution recognized numerically the impinging and suction effect of the wake on the boundary layer. He also considered the motion of vorticity of jets, Fig. 1.

Wiercinski [3] for the first time reported the earlier incipient of laminar-turbulent transition

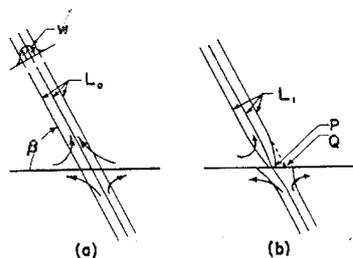


Figure 1: A flat plate and a wake, a) suction and pressure action, b) extra vorticity influence [2]

behind the negative jet than behind the positive one. Later on he [5] reported the differences in measured velocity fluctuations behind both kind of jets. Large and detailed enhancement of this investigation was given by Zabski [6]. It was shown that behind negative jets develop strong disturbances which gradually cover the whole area between wakes. Thus the boundary layer changes its character from laminar one disturbed by oncoming wakes to fully turbulent one.

Though the boundary layer flow is cut by periodically approaching wakes laminar flow persists in between in some circumstances. This is possible for accelerating flow and for early stages of laminar-turbulent transition. It appears that for a flow with adverse pressure gradient some phenomena that are typical for natural transition take place. These are waves that take some features of Tollmien-Schlichting waves. The presence of oncoming wakes do not disturb the rest of flow if only acceleration is large enough. The observed wave disturbances occurred in the mid-area behind the negative jet and before the positive one. There is nothing like that

behind the positive jet. In next paragraphs some characteristics of these waves will be discussed and their significance for transition.

Experimental rig

Measurements were carried out in a subsonic wind tunnel. Turbulence level did not exceed $Tu=0.08\%$, speed of flow in front of plate and wake generator amounted to $U = 15 \text{ m/s}$. Dimensions of the measurement chamber (width, height, length) were: $600 \times 460 \times 1500 \text{ mm}$ and the corners were chamfered. The boundary layer was studied on the upper surface of a flat

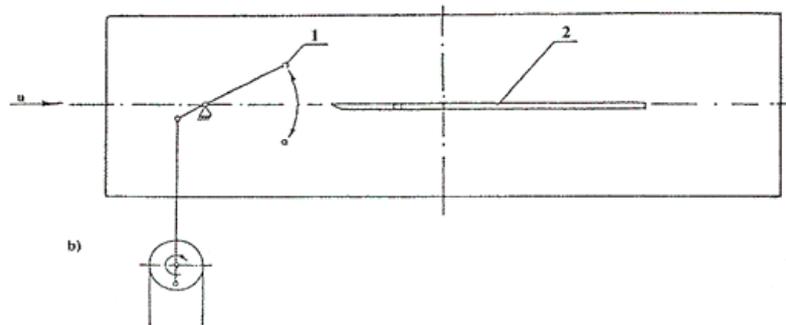


Figure 2: Working section of the wind tunnel: the flat plate and wake generator - a pendulating rod

plate (position 2 in Fig.2) which dimensions were: $600 \times 700 \times 14 \text{ mm}$ (width, length, thickness). The plate attack angle was equal to $\alpha = -2.44^\circ$ that made flow acceleration over upper surface characterized by coefficient $K = \nu/U^2 \cdot dU/dx = 3.36 \cdot 10^{-7}$. Measurements were made at several stations that traverse most of the plate length from $Re_x = 168843$ ($x = 165 \text{ mm}$) to $Re_x = 686918$ ($x = 650 \text{ mm}$). The round rod of $d = 3 \text{ mm}$ diameter and $l = 600 \text{ mm}$ length served as wake generator (position 1 in Fig.2). It moved up and down and the magnitude of its motion amounted to 200 mm . The frequency of this pendulum device was $f = 4 \text{ Hz}$. When the rod is on equal height with the leading edge of the plate in horizontal position phase mark is recorded for further data processing. Velocity and its fluctuations were measured by means of DANTEC StreamLine 90N10 Frame with CTA Module 90C10 with SteamWare software. The probe was DANTEC 55P15 and the data was collected by National Instruments 6040E acquisition computer plug-in. Sampling frequency was $f = 5 \text{ kHz}$ and sample length $t = 10 \text{ s}$. The final distance of the last point in traverse was determined by means of so-called „hydraulic-zero method”. The same method yielded also the slope of velocity curve at the wall and hence friction velocity and drag coefficient. The use of so-called phase-averaging procedure was the farther step of investigation. Single periods of velocity ensemble were put one on another and the averaging along matching points was performed. These result was a single period of velocity time-trace with suppressed random disorders. The time duration of the averaged period was equal to $t = 0.25 \text{ s}$. A screenshot of the procedure in MATLAB with illustrative lines is given in Fig.3. Taking into account ensemble duration, sampling frequency, frequency of rod movement the obtained period is a result of averaging 40 periods. More details was included by Zabski [6].

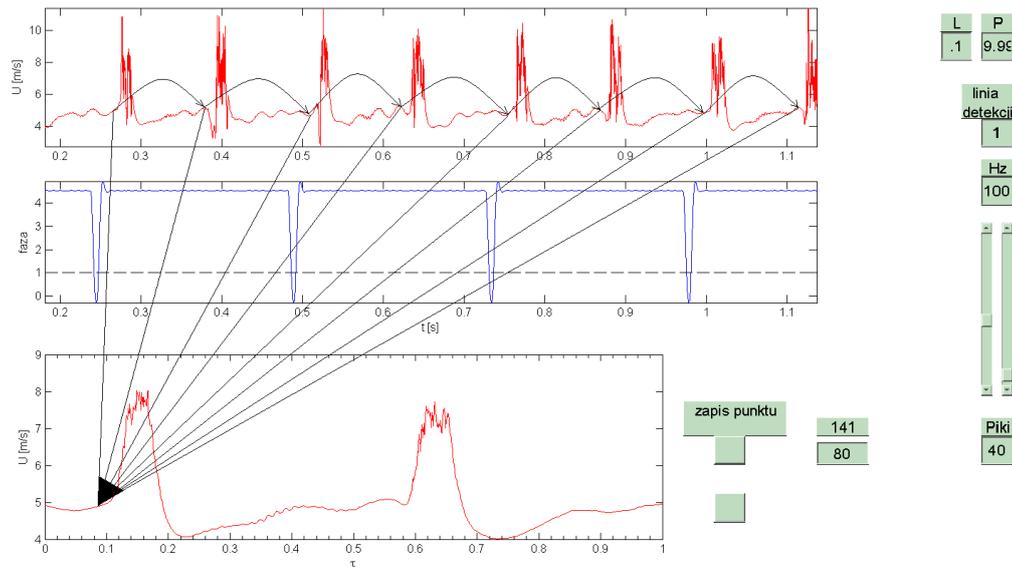


Figure 3: Phase-averaging process

α	K
0.29°	$8.66 \cdot 10^{-8}$
1.05°	$1.55 \cdot 10^{-7}$
2.44°	$3.36 \cdot 10^{-7}$

Table 1: Mean values of acceleration coefficient K

Results

Measurements were carried out for three attack angles of the plate: the aforementioned $\alpha = 2.44^\circ$, $\alpha = 1.05^\circ$ and $\alpha = 0.29^\circ$. These angles refer to the acceleration coefficient values given in Tab.1. The most interesting results were obtained for $\alpha = 2.44^\circ$. Some time averaged results are shown in Fig.4 and Fig.5. The former produces mean drag coefficient along the plate, the latter mean velocity profile at $x = 365 \text{ mm}$ ($Re_x = 346604$). The flow seems rather laminar with increased turbulence level. The evidence of that is drag coefficient distribution which is parallel to laminar curve in Fig.4 ($C_f = 0.664 \cdot Re^{-0.5}$) but above it. Last three points go slightly upwards that indicates an incipience of transition to turbulence. The red marked point indicates position of velocity profile produced in Fig.5. Mean velocity profile seems also laminar, however it is situated mostly above the Blasius curve. There is an inflexion point in the profile.

More information carries phase-averaged chart, Fig.6. These are time traces of velocity in a station corresponding to $Re_x = 346604$. Time on the graph is made dimensionless t/T and T denotes the period equal to $T = 0.25 \text{ s}$ as earlier described. The upmost trace is located at the edge of boundary layer at $y/\delta = 1$, $y^+ = 82.97$. On the other hand the most lower put trace corresponds to $y/\delta = 0.03$, $y^+ = 2.80$. There is a gap in velocity traces at the middle of chart, between $U = 10 \text{ m/s}$ and $U = 8 \text{ m/s}$. The first trace of strongly concentrated traces refers to $U = 8.57 \text{ m/s}$, $y/\delta = 0.27$ and $y^+ = 22.24$. Two strong velocity defects considering from left to right are negative and positive jets respectively. However this is not the main objective of this

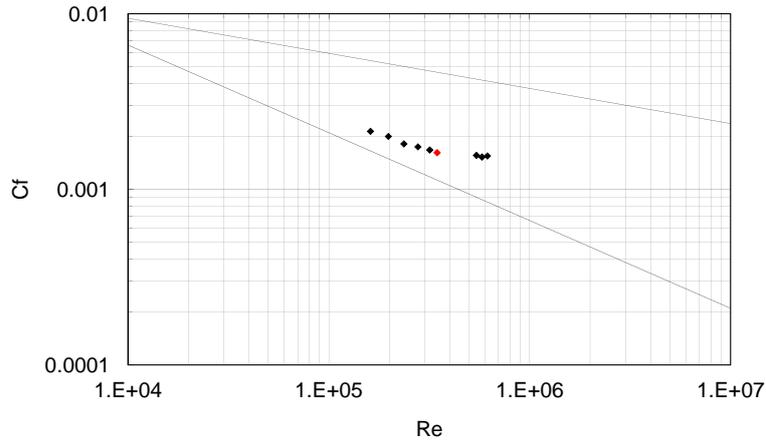


Figure 4: Drag coefficient C_f longitudinal distribution, red point - $Re_x = 346604$, $x = 365 \text{ mm}$

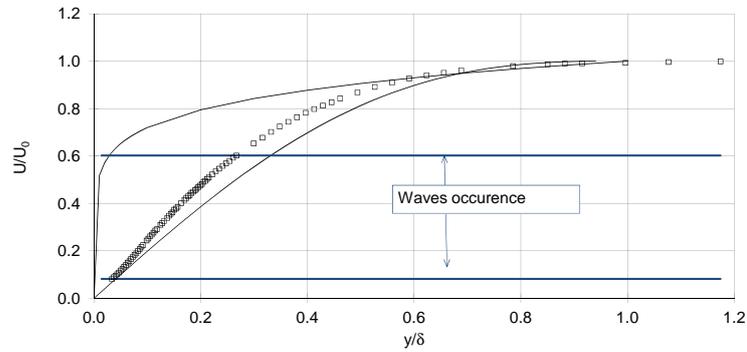


Figure 5: Mean velocity profile

paper.

A red rectangle marked with section B-B encloses a zone in centre between the negative and positive jet with small, smooth disturbances. A thin chart on the right side produces them zoomed. These disturbances look like periodic waves and their duration amounts to about $t/T = 0.1$ of period, that is $t = 0.025 \text{ s}$. The flow around seems undisturbed and laminar up to jets. Above $y/\delta = 0.27$, $U = 8.57 \text{ m/s}$ waves vanish. To calculate frequency some of them surrounded with small rectangle were chosen. It is equal to $f = 337 \text{ Hz}$. This value is convergent to most amplified Tollmien-Schlichting waves according to Walker [4]

$$f = \frac{3.2U_0^2 Re_{\delta^*}^{-\frac{3}{2}}}{2\pi\nu} = 334 \text{ Hz}. \quad (1)$$

Further investigation reveals that the waves are put up in a region of amplified frequencies of TS waves, Fig.7. Neutral curves were taken from Levecenko [1] for three acceleration coefficients Λ

$$\Lambda = \frac{\delta^2}{\nu} \cdot \frac{dU}{dx}. \quad (2)$$

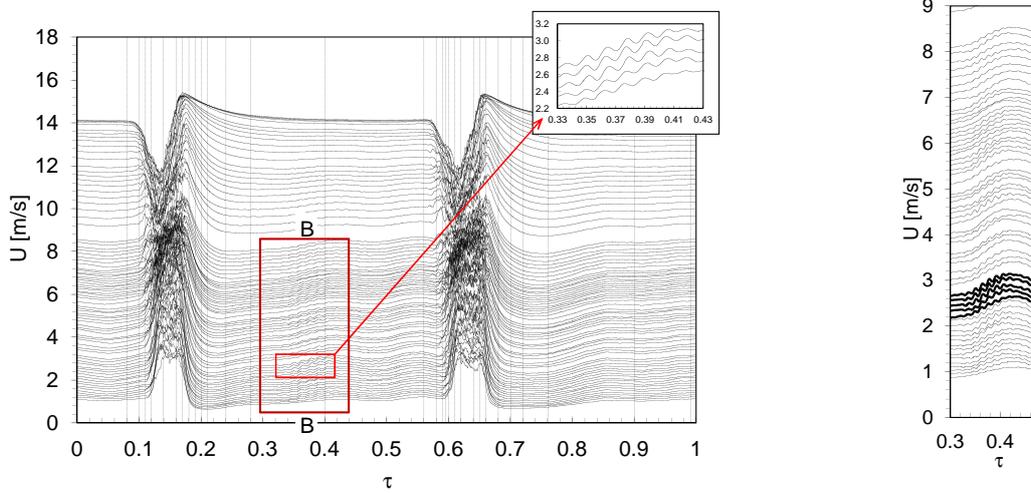


Figure 6: Phase-averaged velocity traces at $Re_x = 346604$

The neutral curves used in Fig.7 seem rather rough. The investigated frequencies lie in the unstable region for $\Lambda = 0$ and on the Walker curve. However, they are not enclosed in the area of $\Lambda = 1$ and $\Lambda = 2$. It is worth reminding that mean external flow acceleration was equal to $K = 3.36 \cdot 10^{-7}$ and local acceleration coefficient Λ referring to boundary layer thickness δ at the investigated distance from the leading edge equaled to $\Lambda = 2.34$. More precise neutral curves would be an advantage. Though wave disturbances occurred at the distance from the leading edge of $Re_x = 346604$, $x = 365 \text{ mm}$ there was no evidence of them in other sites. That refers also to two neighbouring stations $Re_x = 317720$, ($x = 335 \text{ mm}$) and $Re_x = 542542$, ($x = 560 \text{ mm}$), Fig.4. For nearly horizontally positioned plate ($\alpha = -0.29^\circ$) there was no

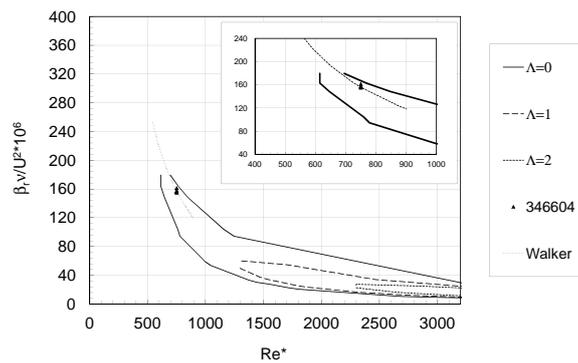


Figure 7: Curves of neutral stability of TS waves

evidence of waves for all traverses investigated ($Re_x = 168843 \div 686918$, $x = 165 \div 650 \text{ mm}$) in phase-averaged charts. It is intriguing, however, that in the middle of a region behind a negative jet and in front of positive one there are strong harsh disturbances, Fig.8, section B-B. Two subsequent measured traverses in Fig.8 reveal laminar-turbulent transition induced by widening

sharp velocity distortion past negative jet, section A-A. The phenomenon of negative jet influence on turbulization was touched at the beginning of this paper. Hier one should notice the coincidence of TS-like waves in accelerated flow locus, Fig.6, and present disturbances locus, Fig.8. Taking into consideration both charts it seems that velocity distortions in section B-B are excited by disturbances shedding from the negative jet. However, the appearance of the region between wakes suggests some instability in B-B and its amplification by the negative jet. No evidence of distortions in section B-B was observed for the first station measured for this case ($\alpha = -0.29^\circ$, $Re_x = 168843$).

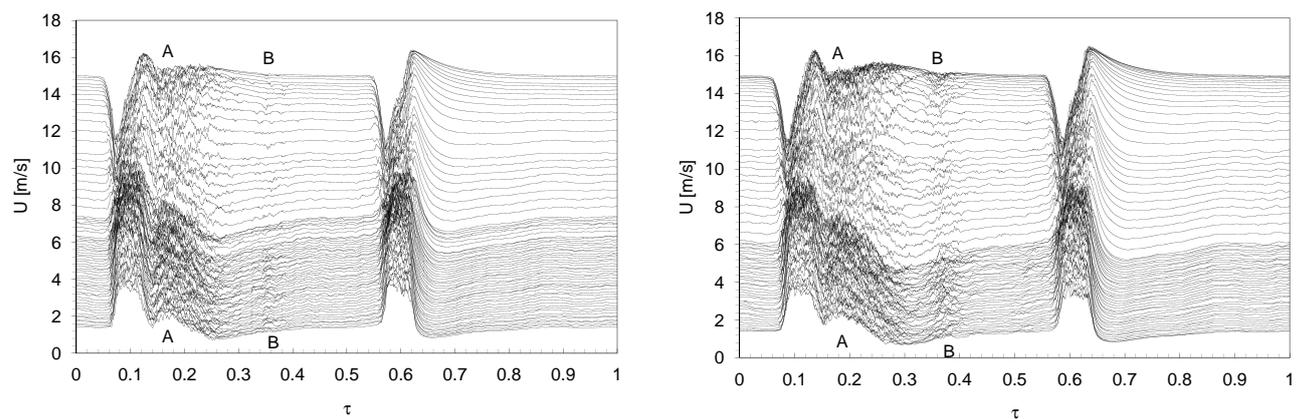


Figure 8: Phase-averaged velocity traces for constant velocity at $Re_x = 216979$ and $Re_x = 257412$

Conclusions

Unexpected disturbances that take some features of TS waves were observed in the boundary layer strongly disturbed by oncoming wakes. The flow was accelerating and the wave locus was at the centre of the laminar area behind a negative jet and in front of positive one. On the other hand there was no such evidence past a positive jet and behind a negative one. For a constant velocity flow there were observed sharp velocity distortions at the same site as TS waves for accelerating case. They had the same appearance as disturbances shedding from the negative jet which drove to turbulent flow. Further investigation seems necessary to get to know the phenomena which take place behind negative jets which is supposed to be a source of disturbances.

References

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