

Qualitative properties of solutions to the complete fluid systems

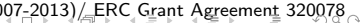
Eduard Feireisl

Institute of Mathematics, Academy of Sciences of the Czech Republic, Prague

with E.Chiodaroli (Lausanne), T.Karper (Trondheim), O.Kreml (Praha)
A.Novotný (Toulon), Y.Sun (Nanjing)

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Navier-Stokes-Fourier system

Mass conservation

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0$$

Momentum balance

$$\begin{aligned} \partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \nabla_x p(\varrho, \vartheta) &= \operatorname{div}_x \mathbb{S}(\nabla_x \mathbf{u}) \\ \operatorname{div}_x \mathbb{S}(\nabla_x \mathbf{u}) &\equiv \mu \Delta \mathbf{u} + \lambda \nabla_x \operatorname{div}_x \mathbf{u}, \quad \mu > 0, \quad \lambda \geq 0 \end{aligned}$$

Internal energy equation

$$\begin{aligned} c_v [\partial_t(\varrho \vartheta) + \operatorname{div}_x(\varrho \vartheta \mathbf{u})] - \operatorname{div}_x(\kappa(\vartheta) \nabla_x \vartheta) \\ = \mathbb{S}(\nabla_x \mathbf{u}) : \nabla_x \mathbf{u} - \vartheta p_\vartheta(\varrho, \vartheta) \operatorname{div}_x \mathbf{u}, \quad \kappa(\vartheta) > 0 \end{aligned}$$

Total energy balance

Pressure

$$p(\varrho, \vartheta) = a\varrho^\gamma + b\varrho + \varrho\vartheta, \quad \gamma > 3, \quad a, b > 0$$

Boundary conditions

$$\mathbf{u}|_{\partial\Omega} = 0, \quad \nabla_x \vartheta \cdot \mathbf{n}|_{\partial\Omega} = 0$$

Total energy balance

$$E(t) = \int \left[\frac{1}{2} \varrho |\mathbf{u}|^2 + c_v \varrho \vartheta + \frac{a}{\gamma - 1} \varrho^\gamma + b \varrho \log(\varrho) \right]$$

$$\frac{d}{dt} E(t) = 0$$

Existence of smooth solutions (classical theory)

Local solutions

A. Valli, W.Zajaczkowski [1982] Existence of classical *local-in-time* solutions in the class:

$$\varrho \in C([0, T_{\max}); W^{3,2}(\Omega)), \vartheta_0 \in C([0, T_{\max}); W^{3,2}(\Omega))$$
$$\mathbf{u} \in C([0, T_{\max}); W^{3,2}(\Omega; R^3))$$

Global solutions

A.Matsumura, T.Nishida [1980,1983] Existence of classical *global-in-time* solutions in the same class for the initial data sufficiently close to a static state

Navier-Stokes-Fourier system (weak form)

Mass conservation

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0$$

Momentum balance

$$\partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \nabla_x p(\varrho, \vartheta) = \operatorname{div}_x \mathbb{S}(\nabla_x \mathbf{u})$$

Thermal energy balance

$$c_v \partial_t(\varrho \vartheta) + c_v \operatorname{div}_x(\varrho \vartheta \mathbf{u}) - \Delta(K(\vartheta)) \\ \geq \mathbb{S}(\nabla_x \mathbf{u}) : \nabla_x \mathbf{u} - p_\vartheta(\varrho, \vartheta) \vartheta \operatorname{div}_x \mathbf{u}$$

Total energy balance

$$\int \left[\frac{1}{2} \varrho |\mathbf{u}|^2 + \frac{a}{\gamma - 1} \varrho^\gamma + b \varrho \log(\varrho) + c_v \varrho \vartheta \right] (\tau, \cdot) \, dx \leq E_0$$

Existence of weak solutions

Existence of weak solutions [E.F.2003]

The Navier-Stokes-Fourier system admits a global-in-time weak solution for any finite energy initial data

Compatibility property of weak solutions

Any weak solution that possesses necessary smoothness is a classical one

Problems

- density oscillations
- temperature concentrations

Euler-Fourier system

Mass conservation

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0$$

Momentum balance

$$\partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \nabla_x(\varrho \vartheta) = 0$$

Internal energy balance

$$\frac{3}{2} \left[\partial_t(\varrho \vartheta) + \operatorname{div}_x(\varrho \vartheta \mathbf{u}) \right] - \Delta \vartheta = -\varrho \vartheta \operatorname{div}_x \mathbf{u}$$

Existence of weak solutions

Initial data

$$\varrho_0, \vartheta_0, \mathbf{u}_0 \in C^3, \varrho_0 > 0, \vartheta_0 > 0$$

Global existence [E.Chiodaroli, E.F., O.Kreml 2013]

For any (smooth) initial data $\varrho_0, \vartheta_0, \mathbf{u}_0$ the Euler-Fourier system admits infinitely many weak solutions on a given time interval $(0, T)$

Regularity class

$$\varrho \in C^2, \partial_t \vartheta, \nabla_x^2 \vartheta \in L^p \text{ for any } 1 \leq p < \infty$$

$$\mathbf{u} \in C_{\text{weak}}([0, T]; L^2) \cap L^\infty, \operatorname{div}_x \mathbf{u} \in C^1$$

Blow-up criterion

Blow-up of smooth solutions [E.F., Y.Sun 2014]

Suppose that the initial data ϱ_0 , ϑ_0 , and \mathbf{u}_0 are smooth ($W^{2,3}$). Then the Navier-Stokes-Fourier system admits a strong solution defined on a (possibly short) time interval $(0, T)$.

If

$$\sup_{t \in (0, T)} [\|\varrho\|_{L^\infty} + \|\vartheta\|_{L^\infty} + \|\mathbf{u}\|_{L^\infty}] < \infty,$$

then the solution can be extended beyond T .

Regularity criterion

Regularity for weak solutions [E.F., Y.Sun 2014]

Suppose that the initial data ϱ_0 , ϑ_0 , and \mathbf{u}_0 are smooth ($W^{2,3}$). Let $[\varrho, \vartheta, \mathbf{u}]$ be a weak solution of the Navier-Stokes-Fourier system such that

$$\sup_{t \in (0, T)} [\|\varrho\|_{L^\infty} + \|\vartheta\|_{L^\infty} + \|\mathbf{u}\|_{L^\infty} + \|\operatorname{div}_x \mathbf{u}\|_{L^\infty}] < \infty.$$

Then $[\varrho, \vartheta, \mathbf{u}]$ is regular.

Main tools

- Relative energy functional
- Entropy formulation
- Weak strong uniqueness

Numerical solution

FV framework

regular tetrahedral mesh, $Q_h = \{v \mid v = \text{piece-wise constant}\}$

FE framework - Crouzeix - Raviart

$V_h = \left\{ v \mid v = \text{piece-wise affine, } \tilde{v}_\Gamma \text{ continuous on face } \Gamma \right\}$

$$\tilde{v}_\Gamma \equiv \frac{1}{|\Gamma|} \int_\Gamma v \, dS_x$$

Upwind discretization of convective terms

$$\begin{aligned} \langle h\mathbf{u}; \nabla_x \varphi \rangle_E &\approx \sum_\Gamma \int_\Gamma h(\cdot - \tilde{\mathbf{u}}_\Gamma \cdot \mathbf{n}) \tilde{\mathbf{u}}_\Gamma \cdot \mathbf{n} [[\varphi]] \, dS_x \\ &\equiv \sum_\Gamma \int_\Gamma \text{Up}[h, \mathbf{u}][[\varphi]] \, dS_x \end{aligned}$$

Numerical scheme [Karlsen-Karper], I

Equation of continuity

$$\int_{\Omega} D_t \varrho_h^k \varphi_h \, dx \equiv \int_{\Omega} \frac{\varrho_h^k - \varrho_h^{k-1}}{\Delta t} \varphi_h \, dx = \sum_{\Gamma \in \Gamma_h} \int_{\Gamma} \text{Up}[\varrho_h^k, u_h^k] [[\varphi_h]] \, dS_x$$

$$-h^\alpha \left[\sum_{\Gamma \in \Gamma_h} \int_{\Gamma} [[[\varrho_h^k]]][[\varphi_h]] \, dS_x \right] \text{ for all } \varphi_h \in Q_h$$

Numerical scheme [Karlsen-Karper], II

Momentum equation

$$\begin{aligned} & \int_{\Omega} D_t (\varrho_h^k \widehat{\mathbf{u}}_h^k) \cdot \varphi_h \, dx - \sum_{\Gamma \in \Gamma_h} \int_{\Gamma} \text{Up}[\varrho_h^k \widehat{\mathbf{u}}_h^k, \mathbf{u}_h^k] \cdot [[\widehat{\varphi}_h]] \, dS_x \\ & - \int_{\Omega} p(\varrho_h^k, \vartheta_h^k) \text{div}_x \varphi_h \, dx - h^\alpha \boxed{\sum_{\Gamma \in \Gamma_h} \int_{\Gamma} [[\varrho_h^k]] [[\widehat{\varphi}_h]] \{ \widehat{\mathbf{u}}_h^k \} \, dS_x} \\ & = - \int_{\Omega} (\mu \nabla_h \mathbf{u}_h : \nabla_h \varphi + \lambda \text{div}_h \mathbf{u}_h^k \text{div}_h \varphi_h) \, dx \\ & \text{for all } \varphi_h \in V_h(\Omega) \end{aligned}$$

Numerical scheme [Karlsen-Karper], III

Energy equation

$$\begin{aligned} & \int_{\Omega} D_t (\varrho_h^k \vartheta_h^k) \varphi_h \, dx - \sum_{\Gamma \in \Gamma_h} \int_{\Gamma} \text{Up}[\varrho_h^k \vartheta_h^k, \mathbf{u}_h^k] [[\varphi_h]] \, dS_x \\ & \quad + \sum_{\Gamma \in \Gamma_h} \int_{\Gamma} \frac{1}{d_h} [[K(\vartheta_h^k)]] [[\varphi_h]] \, dS_x \\ & = \int_{\Omega} (\mu |\nabla_h \mathbf{u}_h^k|^2 + \lambda |\text{div}_h \mathbf{u}_h^k|^2) \varphi_h \, dx \\ & \quad - \int_{\Omega} \vartheta_h^k \partial_{\vartheta} p(\varrho_h^k, \vartheta_h^k) \text{div}_x \mathbf{u}_h^k \varphi_h \, dx \\ & \quad \text{for all } \varphi_h \in Q_h(\Omega) \end{aligned}$$

Convergence of the numerical scheme

Convergence of the numerical scheme [E.F., T.Karper, A.Novotný 2014]

The numerical solutions converge, up to a subsequence, to a weak solution of the Navier-Stokes-Fourier system

Synergy analysis - numerics, assumptions

Numerical solutions with regular initial data

Suppose that $[\varrho_h, \vartheta_h, \mathbf{u}_h]$ is a sequence of numerical solutions for regular initial data

Boundedness

Suppose that

$$\varrho_h^k, \vartheta_h^k, \mathbf{u}_h^k, \operatorname{div}_h \mathbf{u}_h^k$$

are bounded independently of the order of discretization h .

Synergy analysis - numerics, conclusion

Conclusion

The numerical solutions converge to a weak solution with

$$\sup_{t \in (0, T)} [\|\varrho\|_{L^\infty} + \|\vartheta\|_{L^\infty} + \|\mathbf{u}\|_{L^\infty} + \|\operatorname{div}_x \mathbf{u}\|_{L^\infty}] < \infty.$$

Consequently:

- the limit solution is smooth
- the limit solution is unique
- the numerical scheme converges unconditionally
- error estimates (?)