



INSTITUTE of MATHEMATICS

ACADEMY of SCIENCES of the CZECH REPUBLIC

**A small note on L^q theory for Stokes
problem with a pressure dependent
viscosity**

Václav Mácha

Preprint No. 63-2014

PRAHA 2014

A small note on L^q theory for Stokes problem with a pressure dependent viscosity

VÁCLAV MÁCHA*

Institute of Mathematics AS CR,
Žitná 25, 115 67 Praha 1, Czech Republic
email: macha@math.cas.cz

Abstract

In this paper we study higher local integrability of a weak solution to the steady Stokes problem. We consider a case of a pressure and shear-rate dependent viscosity, i.e. the elliptic part of the Stokes problem is assumed to be nonlinear and it depends on p and on a symmetric part of a gradient of u , namely, it is represented by a stress tensor $T(Du, p) := \nu(p, |D|^2)D$ which satisfies r -growth condition with $r \in (1, 2]$. In order to get the main result, we use Calderón-Zygmund theory and the method which was presented for example in L. A. Caffarelli, I. Peral: On $W^{1,p}$ Estimates for Elliptic Equations in Divergence form. Comm. on Pure and Appl. Math, Vol. LI, 1–21 (1998).

Keywords: Stokes problem, L^q theory, pressure dependent viscosity

MSC 2010: 35B65, 35Q35, 76D03

1 Introduction

In the celebrated Navier-Stokes system which deals with a flow of Newtonian incompressible fluid, a viscosity is assumed to be constant. However, the constant viscosity cannot explain many interesting physical phenomena such as shear-thinning, shear-thickening, die-swell, etc. A viscosity of non-Newtonian fluids is not generally constant but depends on a shear rate and, as many experimental works show, there are also fluids whose viscosity depends also on a pressure. On the other hand, changes in a density of these liquids are negligible as the pressure grows and thus this fluids can be still treated as incompressible, see c.f. [2]. Also in some situations the pressure grows tremendously and its influence on a viscosity can not be neglected. For example, in works written by Knauf et. al. and Lanzendörfer ([14], [16]) the authors provide numerical simulation of flow of lubricant through a ball bearing (resp. journal bearing). The governing equation for a steady flow of a fluid with a pressure- and shear rate-dependent viscosity has a form

$$\begin{aligned} -\operatorname{div} T(Du, p) + \operatorname{div}(u \otimes u) + \nabla p &= \operatorname{div} T(F, 0) \text{ in } \Omega, \\ \operatorname{div} u &= 0 \text{ in } \Omega, \end{aligned}$$

where $\Omega \subset \mathbb{R}^d$ is a domain, $T : \mathbb{R}^{d^2} \times \mathbb{R} \mapsto \mathbb{R}^{d^2}$ and $F : \Omega \mapsto \mathbb{R}^{d^2}$ are given function, $u : \Omega \mapsto \mathbb{R}^d$ and $p : \Omega \mapsto \mathbb{R}$ are unknowns and Du stands for a symmetric part of a gradient of u , i.e. $2Du = \nabla u + (\nabla u)^T$.

A plenty of works studying this system has been published yet. Regarding the steady case, Gazzola and Secchi in [11] consider an only-pressure-dependent viscosity. The existence of solution for more general viscosity is discussed by Bulíček and Fišerová in [3], by Lanzendörfer and Stebel in [17] and also by Franta et al in [10]. $C^{1,\alpha}$ interior regularity in two dimensions is solved in

*This work was supported by the GAČR project GA13-00522S in the general framework of RVO: 67985840.

Bulíček and Kaplický in [4], the Hölder regularity in the three dimensional case was announced by Mingione, Málek and Stará in [22]. Further, $C^{1,\alpha}$ regularity under Dirichlet boundary conditions in 3 dimensions is solved in [18]. The unsteady case was investigated, for example, by Bulíček, Málek and Rajagopal in [5].

This article is devoted to the interior L^q regularity for a simplified problem which has a following form

$$\begin{aligned} -\operatorname{div} T(Du, p) + \nabla p &= \operatorname{div} T(F, 0) \text{ in } \Omega, \\ \operatorname{div} u &= 0 \text{ in } \Omega. \end{aligned} \quad (P)$$

Since this article concerns only local properties of solution, we do not care about any boundary conditions.

The stress tensor T is considered to be in form $T(Du, p) = \nu(p, |Du|^2)Du$. Moreover, we suppose that following growth conditions are fulfilled.

$$\gamma_1(1 + |D|^2)^{\frac{r-2}{2}} |B|^2 \leq \frac{\partial(\nu(p, |D|^2)D_{ij})}{\partial D_{kl}} B_{ij} B_{kl} \leq \gamma_2(1 + |D|^2)^{\frac{r-2}{2}} |B|^2 \quad (A1)$$

$$\left| \frac{\partial \nu(p, |D|^2)}{\partial p} D \right| \leq \gamma_3(1 + |D|^2)^{\frac{r-2}{4}}. \quad (A2)$$

Furthermore, γ_3 is assumed to be small, specifically

$$\gamma_3 < \frac{\gamma_1}{c_{\operatorname{div}}(\gamma_1 + \gamma_2)}, \quad (A3)$$

where the constant c_{div} comes from Bogovskii inequality – see Lemma 4.

For a physically relevant viscosity which fulfills (A1), (A2) and (A3) we refer reader to [16], [20] or [23].

We point out that the form of the right hand side of $(P)_1$ is not restrictive since for every $G : \Omega \mapsto \mathbb{R}^{d^2}$ there exists F such that $G = T(F, 0)$. The chosen form of the right hand side allows us to write the main result in a clear way.

In what follows, usual Lebesgue spaces of functions defined on a set Ω are denoted by $L^s(\Omega)$, Sobolev spaces are denoted by $W^{1,s}(\Omega)$. Lebesgue (resp. Sobolev) spaces of vector-valued functions are denoted by $L^s(\Omega)^d$ (resp. $W^{1,s}(\Omega)^d$). However, we will use $L^s(\Omega)$ (resp. $W^{1,s}(\Omega)$) even for vector-valued functions in case we avoid any misunderstandings. Norms in these spaces will be denoted by $\|\cdot\|_s$ (resp. $\|\cdot\|_{1,s}$). In case we want to emphasize the supporting set, we write $\|\cdot\|_{s,\Omega}$ (resp. $\|\cdot\|_{1,s,\Omega}$), i.e. $\|f\|_{s,\Omega} = \left(\int_{\Omega} |f|^s\right)^{\frac{1}{s}}$. Further, $L_{\operatorname{loc}}^s(\Omega)$ (resp. $W_{\operatorname{loc}}^{1,s}(\Omega)$) stands for local version of Lebesgue (resp. Sobolev) space. A space of all functions from $W^{1,s}(\Omega)$ with compact support is denoted by $W_c^{1,s}(\Omega)$. We define also an integral average as follows

$$\int_{\Omega} f := \frac{1}{|\Omega|} \int_{\Omega} f.$$

For $\alpha > 0$ and a cube $Q \subset \mathbb{R}^d$ centered at $a \in \mathbb{R}^d$ we define a cube αQ as

$$\alpha Q := \left\{ x \in \mathbb{R}^d, \frac{x-a}{\alpha} \in Q \right\}.$$

A weak solution of (P) is a pair $(u, p) \in W_{\operatorname{loc}}^{1,r}(\Omega) \times L_{\operatorname{loc}}^{r'}(\Omega)$, $\operatorname{div} u = 0$ fulfilling

$$\int_{\Omega} T(Du, p) \nabla \varphi - \int_{\Omega} p \operatorname{div} \varphi = \int_{\Omega} T(F, 0) \nabla \varphi$$

for every $\varphi \in W_c^{1,r}(\Omega)$. The existence of a weak solutions to (P) is solved by Theorem 3.11 in [15]. We also refer to Theorem 2 in [16] and references given there. In the rest of this paper we always assume that (A1), (A2) and (A3) hold.

The main result of this paper is summed up in the following theorem.

1 Theorem. Let $r \in (1, 2]$, $q \in \left(1, \frac{d}{d-2}\right)$ (resp. $q \in (1, \infty)$ in case $d = 2$) and let assumptions (A1)–(A3) be fulfilled. Let $f \in L^{qr}(\Omega)$, (u, p) be a weak solution to (P) and let $Q \subset 4Q \subset \Omega$ be a sufficiently small cube. Then¹

$$\int_Q \left((1 + |Du|^2)^{\frac{r-2}{2}} |Du|^2 \right)^q \leq c \left(1 + \int_{4Q} F^{qr} + \left(\int_{4Q} (1 + |Du|^2)^{\frac{r-2}{2}} |Du|^2 \right)^q \right), \quad (1.1)$$

and, moreover,

$$\int_Q \left| p - \int_Q p \right|^{qr'} \leq c \left(1 + \int_{4Q} F^{qr} + \left(\int_{4Q} (1 + |Du|^2)^{\frac{r-2}{2}} |Du|^2 \right)^q \right), \quad (1.2)$$

where c is a constant independent of f, u, p and Q .

The method of the proof is based on local comparison with the solution to the problem with zero right hand side. Among lots of papers based on the comparison technique we would like to mention [13] where Iwaniec showed L^q theory result for linear problem, [6] where Caffarelli and Peral proved L^q estimates for elliptic equations in divergence form and [8] where Diening and Kaplický used the method to derive L^q estimates of Stokes system with general growth. Our approach is based on [6, Theorem A]. However, since our system is slightly more complicated than the one in [6], we have to use a more general version which is presented here as Lemma 5.

2 Preliminaries

We define a function $V : \mathbb{R}^{d^2} \mapsto \mathbb{R}^{d^2}$ as $V(D) := \sqrt{\nu(0, |D|^2)} D$ and a function $\Phi'(s) = \nu(0, |s|^2) s$. We emphasize that $\Phi(x) := \int_0^x \Phi'(s) ds$ is an N-function satisfying Δ_2 condition and, moreover, $\Phi'(s) \sim (1 + s^2)^{\frac{r-2}{2}} s$. By Φ^* we denote an N-function which is complementary to Φ . Hereinafter, we suppose that $r \leq 2$.

2 Lemma. For all cubes Q , $u, v \in W^{1,r}(Q)$ and $p, \pi \in L^{r'}(Q)$ it holds

- $\Phi'(|Du|) \sim |T(Du, p)|$ uniformly in p ,
- $|V(Du)|^2 \sim T(Du, p) Du \sim (1 + Du)^{\frac{r-2}{2}} |Du|^2 \sim \Phi(|Du|) \sim \Phi^*(\Phi'(|Du|))$, uniformly in p ,
-

$$\begin{aligned} I_{(u,v)} &:= \int_0^1 (1 + |Dv + s(Du - Dv)|^2)^{\frac{r-2}{2}} ds |Du - Dv|^2 \\ &\leq \frac{2}{\gamma_1} (T(Du, p) - T(Dv, \pi))(Du - Dv) + \frac{\gamma_3^2}{\gamma_1^2} |p - \pi|^2, \end{aligned} \quad (2.1)$$

- $I_{(u,v,Q)} := \int_Q I_{u,v} \sim \int_Q |V(Du) - V(Dv)|^2$.
- For all $\varphi \in W^{1,2}(Q')$ it holds

$$\int_Q (T(Du, p) - T(Dv, \pi))(D\varphi) \leq \gamma_2 \sqrt{I_{u,v,Q}} \|\nabla \varphi\|_{2,Q} + \gamma_3 \|p - \pi\|_2 \|\nabla \varphi\|_2.$$

Proof. Since $T(0, p) = 0$, it holds

$$\begin{aligned} |T(Du, p)| |Du| &= \int_0^1 \frac{\partial}{\partial s} |T(s Du, p)| ds |Du| \sim \int_0^1 (1 + |s Du|^2)^{\frac{r-2}{2}} ds |Du|^2 \\ &\sim (1 + |Du|^2)^{\frac{r-2}{2}} |Du|^2 \sim \Phi'(|Du|) |Du| \end{aligned}$$

¹Hereinafter we use a letter c for a constant which may vary from line to line however, it is always independent of solution and right hand side.

where we use [7, Lemma 19]. The first estimate follows easily.

The second estimate is an easy consequence of the first estimate and of the following sequence of inequalities (which can be found e.g. as [7, Lemma 3]):

$$|V(Du)|^2 \sim \Phi(|Du|) \sim \Phi'(|Du|)|Du| \sim \Phi^*(\Phi'(|Du|)).$$

For the third inequality, we refer reader to [10, Lemma 3.3].

In order to prove the fourth inequality it is enough to see that

$$\begin{aligned} I_{(u,v,Q)} &\sim \int_Q \int_0^1 \frac{\partial}{\partial s} T((s(Du - Dv) + Dv), 0) ds (Du - Dv) \\ &= \int_Q (T(Du, 0) - T(Dv, 0))(Du - Dv) \sim \int_Q |V(Du) - V(Dv)|^2 \end{aligned}$$

where the last estimate comes from [7, Lemma 3].

It remains to prove the fifth estimate. It holds that

$$\begin{aligned} \int_Q (T(Du, p) - T(Dv, \pi)) D\varphi &= \int_Q \left(\int_0^1 \frac{\partial}{\partial s} T(Dv + s(Du - Dv), \pi + s(p - \pi)) ds \right) D\varphi \\ &= \int_Q \int_0^1 \frac{\partial T}{\partial D} (Dv + s(Du - Dv), \pi + s(p - \pi)) (Du - Dv) ds D\varphi \\ &\quad + \int_Q \int_0^1 \frac{\partial T}{\partial p} (Dv + s(Du - Dv), \pi + s(p - \pi)) (p - \pi) ds D\varphi \\ &\leq \gamma_2 \int_Q \int_0^1 (1 + |Dv + s(Du - Dv)|^2)^{\frac{r-2}{2}} |Du - Dv| ds |D\varphi| \\ &\quad + \gamma_3 \left(\int_Q |p - \pi|^2 \right)^{1/2} \left(\int_Q |\nabla \varphi|^2 \right)^{1/2} \end{aligned}$$

Further, since $r < 2$, the Hölder inequality yields

$$\begin{aligned} \int_Q \int_0^1 (1 + |Dv + s(Du - Dv)|^2)^{\frac{r-2}{2}} |Du - Dv| ds |D\varphi| \\ \leq \left(\int_Q \int_0^1 (1 + |Dv + s(Du - Dv)|^2)^{\frac{r-2}{2}} |Du - Dv|^2 \right)^{1/2} \left(\int_Q |\nabla \varphi|^2 \right)^{1/2} \end{aligned}$$

and the demanded estimate follows. \square

3 Lemma (Theorem 6.10 in [9]). *Let $\varphi \in W_0^{1,\phi}(Q)$. There exists a constant c independent of Q and φ such that*

$$\int_Q \Phi(|\nabla \varphi|) \leq c \int_Q \Phi(|D\varphi|)$$

4 Lemma (Lemma 3.3 in [1]). *Let $s \in (1, \infty)$. There is a constant c such that for a cube Q and for all $f \in L^{s'}(Q)$ there exists a function $\varphi \in W^{1,s}(Q)$ such that*

$$\begin{aligned} \operatorname{div} \varphi &= |f|^{s'-2} f - \int_Q |f|^{s'-2} f \text{ in } Q \\ \varphi|_{\partial Q} &= 0 \end{aligned}$$

and,

$$\|\varphi\|_{1,s,Q} \leq c \|f\|_{s',Q}^{1/(s-1)}.$$

For the case $s = 2$ we write c_{div} instead of c .

3 Proof of the main theorem

The proof of the main theorem relies on the following lemma

5 Lemma (Lemma 2.7 from [19]). *Let $\mathcal{O} \subset \mathbb{R}^n$, $1 \leq p < q < s < \infty$, $f \in L^{q/p}(\mathcal{O})$, $g \in L^{q/p}(\mathcal{O})$ and $w \in L^p(\mathcal{O})^n$. Further, let $Q \subset \mathcal{O}$ be a cube and Q_k be dyadic cubes obtained from Q . Then there exists $\varepsilon > 0$ independent of Q and \mathcal{O} such that the following implication holds:*

If for every dyadic sub-cube $Q_k \subset Q$ there exists $w_a \in L^p(4\tilde{Q}_k \cap \mathcal{O})^n$ with following properties:

$$\begin{aligned} \left(\int_{2\tilde{Q}_k \cap \mathcal{O}} |w_a|^s \right)^{\frac{1}{s}} &\leq \frac{C}{2} \left(\int_{4\tilde{Q}_k \cap \mathcal{O}} |w_a|^p \right)^{\frac{1}{p}}, \\ \int_{4\tilde{Q}_k \cap \mathcal{O}} |w_a|^p &\leq C \int_{4\tilde{Q}_k \cap \mathcal{O}} |w|^p + C \int_{4\tilde{Q}_k \cap \mathcal{O}} |g|, \\ \int_{4\tilde{Q}_k \cap \mathcal{O}} |w - w_a|^p &\leq \varepsilon \int_{4\tilde{Q}_k \cap \mathcal{O}} |w|^p + C \int_{4\tilde{Q}_k \cap \mathcal{O}} |f|, \end{aligned}$$

then $w \in L^q(Q)^n$. Positive constants C and ε are independent of Q_k , w_a and w . Furthermore, there exists a positive constant c independent of f , g and w such that

$$\int_Q |w|^q \leq c \left(\int_{4Q \cap \mathcal{O}} |f|^{\frac{q}{p}} + \int_{4Q \cap \mathcal{O}} |g|^{\frac{q}{p}} + \left(\int_{4Q \cap \mathcal{O}} |w|^p \right)^{\frac{q}{p}} \right). \quad (3.1)$$

Let $\frac{1}{4}Q'$ be an arbitrary dyadic cube obtained from Q . We consider a reference system

$$\begin{aligned} -\operatorname{div} T(Dv, \pi) + \nabla \pi &= 0 \text{ in } Q' \\ \operatorname{div} v &= 0 \text{ in } Q' \\ v|_{\partial Q'} &= u, \end{aligned} \quad (R)$$

where we assume that $\int_{Q'} \pi = \int_{Q'} p$. The existence of such solution is due to monotone operator method, see c.f. [24], chapter 2.

6 Lemma. *Let v, π be a solution to (R) on a domain Q' . Let $Q \subset \bar{Q} \subset Q'$ be a cube. Let $\alpha \in (0, 1)$ and $\xi \in C^\infty(Q)$ be such that $\xi|_{\alpha Q} = 1$, $\operatorname{supp} \xi \subset Q$ and $\xi(x) \in [0, 1]$ for every $x \in Q$. Then there exists $\delta < 1$ such that for all $i \in \{1, \dots, d\}$*

$$\int_Q \partial_i T(Dv, \pi)(\xi^2 \partial_i \nabla v) \geq \lim_{h \rightarrow 0} \frac{1}{h^2} I_{(v(x+he_i), v(x), \alpha Q)} - \delta \lim_{h \rightarrow 0} \frac{1}{h^2} I_{(v(x+he_i), v(x), Q)}$$

Proof. We start with the term

$$J_h := \int_Q (T(Dv(x+he_i), p(x+he_i)) - T(Dv(x), p(x)))(Dv(x+he_i) - Dv(x))\xi^2.$$

Note that T is symmetric and thus $\int_Q \partial_i T(Dv, \pi)(\xi^2 \partial_i \nabla v) = \lim_{h \rightarrow 0} \frac{J_h}{h^2}$. We integrate (2.1) with $u = v(x+he_i)$ over Q and we get

$$I_{(v(x+he_i), v(x), \alpha Q)} \leq \frac{2}{\gamma_1} J_h + \frac{\gamma_3^2}{\gamma_1^2} \int_Q |\pi(x+he_i) - \pi(x)|^2. \quad (3.2)$$

For a general function f we define operator $\delta_{-h,i}$ as $\delta_{-h,i} f = f(x-he_i) - f(x)$. Furthermore, we define a test function φ as

$$\begin{aligned} \operatorname{div} \varphi &= \pi(x+he_i) - \pi(x) - \int_Q (\pi(x+he_i) - \pi(x)) \text{ on } Q \\ \varphi &= 0 \text{ on } \partial Q. \end{aligned}$$

We test (R) by $\delta_{-h,i}\varphi$ and, as far as $\delta_{-h,i}K = 0$ for every constant K , we get according to Lemma 2

$$\begin{aligned} & \int_Q |\pi(x + he_i) - \pi(x)|^2 \\ & \leq \gamma_2 c_{\text{div}} \sqrt{I_{(v(x+he_i), v(x), Q)}} \|\pi(x + he_i) - \pi(x)\|_{2,Q} + \gamma_3 c_{\text{div}} \int_Q |\pi(x + he_i) - \pi(x)|^2, \end{aligned}$$

consequently,

$$\frac{\gamma_3^2}{\gamma_1^2} \int_Q |\pi(x + he_i) - \pi(x)|^2 \leq \left(\frac{\gamma_2 \gamma_3 c_{\text{div}}}{\gamma_1 (1 - \gamma_3 c_{\text{div}})} \right)^2 I_{(v(x+he_i), v(x), Q)}.$$

We would like to emphasize that due to (A3) it holds that

$$\left(\frac{\gamma_2 \gamma_3 c_{\text{div}}}{\gamma_1 (1 - \gamma_3 c_{\text{div}})} \right)^2 < 1.$$

In order to get the claim it suffices to put this estimate into (3.2), divide by h^2 and let $h \rightarrow 0$. \square

The first and second estimate from 5 is verified in the following lemma.

7 Lemma. *There exists a constant c independent of v , π , u , p , Q' and Q such that*

$$\left(\int_{\frac{1}{2}Q'} |V(Dv)|^q \right)^{\frac{1}{q}} \leq c \left(\int_{Q'} |V(Dv)|^2 \right)^{\frac{1}{2}}, \quad (3.3)$$

where $q \in (1, \frac{2n}{n-2}]$ (resp. $q \in (1, \infty)$ in case $d = 2$).

Furthermore,

$$\int_{Q'} |V(Dv)|^2 \leq c \int_{Q'} |V(Du)|^2. \quad (3.4)$$

Proof. For simplicity we assume that $d = 3$. This assumption will be commented later. Using standard bootstrap argument one may also easily show that $v \in W_{\text{loc}}^{2,p}(Q')$ and $\nabla V(Dv) \in L_{\text{loc}}^2(Q')$. See cf. [21].

Let $q : \mathbb{R}^3 \mapsto \mathbb{R}^3$ be a linear function with $\nabla q = \langle \nabla v \rangle_{Q'}$. Let α and α' be such that $\frac{1}{2} \leq \alpha < \alpha' \leq \frac{3}{4}$ and let $\xi \in C^\infty(Q)$ be a smooth test function such that $\xi|_{\alpha Q'} = 1$, $\text{supp } \xi \subset \alpha' Q'$, $\xi(x) \in [0, 1]$ for $x \in \alpha' Q'$ and $|\nabla^j \xi| \leq \frac{C}{(\alpha' - \alpha)^j}$. We test (R) by $\varphi = \text{curl}(\xi^2 \text{curl}(v - q))$.

We would like to point out that $\text{curl } v = (\partial_3 v_2 - \partial_2 v_3, \partial_1 v_3 - \partial_3 v_1, \partial_2 v_1 - \partial_1 v_2)$. It follows that φ is not a suitable test function. Instead one has to take $\psi = \text{curl}(\xi^2 \text{curl}_h(v - q))$ where curl_h is defined as $\text{curl}_h(v) = (\delta_{h,3} v_2 - \delta_{h,2} v_3, \delta_{h,1} v_3 - \delta_{h,3} v_1, \delta_{h,2} v_1 - \delta_{h,1} v_2)$. Note that $\text{div } \varphi = \text{div } \psi = 0$. For sake of lucidity we work with φ instead of ψ .

After some cumbersome manipulation we get

$$\begin{aligned} & \int_{\alpha' Q'} \nabla T(Dv, \pi) \xi^2 \nabla^2 v \\ & = \int_{\alpha' Q'} T(Dv, \pi) (\nabla(\nabla(\xi^2) \times \text{curl}(v - q)) + \text{div}(\nabla(\xi^2) \otimes \nabla(v - q)) - \nabla(\nabla(\xi^2) \nabla(v - q))) \end{aligned}$$

In the same way as in the proof of [8, Lemma 3.5] we derive

$$\int_{\alpha' Q'} \nabla T(Dv, \pi) \xi^2 \nabla^2 v \leq \frac{c(\varepsilon)}{(\alpha' - \alpha)^2} \int_{\alpha' Q'} |V(Dv)|^2 + \varepsilon \int_{\alpha' Q'} |\nabla V(Dv)|^2,$$

where ε comes from a Young inequality. Further, since

$$\sum_{i=1}^d \lim_{h \rightarrow 0} \frac{1}{h^2} I_{(v(x+he_i), v(x), \alpha' Q')} \sim \int_{\alpha' Q'} |\nabla V(Dv)|^2,$$

Lemma 6 gives

$$\sum_{i=1}^d \lim_{h \rightarrow 0} \frac{1}{h^2} I_{(v(x+he_i), v(x), \alpha Q')} \leq (\delta + \varepsilon) \sum_{i=1}^d \lim_{h \rightarrow 0} \frac{1}{h^2} I_{(v(x+he_i), v(x), \alpha' Q')} + \frac{c(\varepsilon)}{(\alpha' - \alpha)^2} \int_{\alpha Q'} |V(Dv)|^2.$$

As far as ε could be chosen such that $\delta + \varepsilon < 1$, we may use algebraic lemma for example [12, Lemma 6.1]) in order to get

$$\int_{1/2 Q'} |\nabla V(Dv)|^2 \sim \sum_{i=1}^d \lim_{h \rightarrow 0} \frac{1}{h^2} I_{(v(x+he_i), v(x), 1/2 Q')} \leq \frac{c(\varepsilon)}{(\text{diam } Q')^2} \int_{\frac{3}{4} Q'} |V(Dv)|^2$$

The estimate (3.3) follows due to Poincaré inequality.

The case of general dimension is just a matter of a suitable test function φ . See cf. [19, Proof of Lemma 3.3] for more details.

To prove the second estimate we test (R) by $\varphi = u - v$. We get

$$\int_{Q'} T(Dv, \pi)(Du - Dv) = 0.$$

Consequently, due to Lemma 2

$$\begin{aligned} \int_{Q'} |V(Dv)|^2 &\sim \int_{Q'} T(Dv, \pi) Dv = \int_{Q'} T(Dv, \pi) Du \\ &\leq \delta \int_{Q'} \Phi^*(|T(Dv, \pi)|) + c_\delta \int_{Q'} \Phi(|Du|) \leq \delta \int_{Q'} |V(Dv)|^2 + c_\delta \int_{Q'} |V(Du)|^2, \end{aligned}$$

where δ comes from Young inequality. The estimate (3.4) follows immediately. \square

Proof of Theorem 1. In what follows, we focus on the third estimate of Lemma 5. We subtract (R) from (P). Thus, for all $\varphi \in W_0^{1,r}(Q')$, it holds that

$$\int_{Q'} (T(Du, p) - T(Dv, \pi)) D\varphi - \int_{Q'} (p - \pi) \text{div } \varphi = \int_{Q'} F \nabla \varphi \quad (3.5)$$

Taking $\varphi = u - v$ in (3.5), we get

$$\begin{aligned} I_{(u,v,Q')} &\leq \frac{2}{\gamma_1} \int_{Q'} (T(Du, p) - T(Dv, \pi)) (Du - Dv) + \frac{\gamma_3^2}{\gamma_1^2} \int_{Q'} |p - \pi|^2 \\ &= c \int_{Q'} T(F, 0) (\nabla(u - v)) + \frac{\gamma_3^2}{\gamma_1^2} \int_{Q'} |p - \pi|^2, \quad (3.6) \end{aligned}$$

where we use Lemma 2. Let φ be a solution to

$$\begin{aligned} \text{div } \varphi &= p - \pi \text{ in } Q' \\ \varphi|_{\partial Q'} &= 0. \end{aligned}$$

We use this φ as a test function in (3.5). Lemma 2 yields

$$\begin{aligned} \|p - \pi\|_{2,Q'}^2 &\leq \int_{Q'} (T(Du, p) - T(Dv, \pi)) D\varphi - \int_{Q'} T(F, 0) \nabla \varphi \\ &\leq \gamma_2 \sqrt{I_{(u,v,Q')}} \|\nabla \varphi\|_{2,Q'} + \gamma_3 \|p - \pi\|_{2,Q'} \|\nabla \varphi\|_{2,Q'} + \int_{Q'} T(F, 0) \nabla \varphi \end{aligned}$$

and, due to Lemma 4 and Hölder inequality,

$$(1 - c_{\text{div}}\gamma_3)\|p - \pi\|_{2,Q'} \leq c_{\text{div}}\gamma_2\sqrt{I_{(u,v,Q')}} + \|T(F,0)\|_{2,Q'}.$$

Consequently

$$\|p - \pi\|_{2,Q'}^2 \leq \left(\left(\frac{c_{\text{div}}\gamma_2}{1 - c_{\text{div}}\gamma_3} \right)^2 + \varepsilon \right) I_{(u,v,Q')} + c(\varepsilon)\|T(F,0)\|_{2,Q'}^2.$$

We get, according to Lemma 3, Lemma 2, Young inequality and (3.4), that

$$\begin{aligned} \int_{Q'} T(F,0)(\nabla(u-v)) &\leq c(\varepsilon_1) \int_{Q'} \Phi^*(\Phi'(|F|)) + \varepsilon_1 \int_{Q'} \Phi(|Du - Dv|) \\ &\leq c(\varepsilon_1) \int_{Q'} \Phi(|F|) + \varepsilon_1 \int_{Q'} \Phi(|Du|) + \varepsilon_1 \int_{Q'} \Phi(|Dv|) \\ &\leq c(\varepsilon_1) \int_{Q'} \Phi(|F|) + \varepsilon \int_{Q'} |V(Du)|^2 \end{aligned}$$

We put all these estimate into (3.6) and we get

$$I_{(u,v,Q')} \leq \left(\left(\frac{\gamma_3 c_{\text{div}}\gamma_2}{\gamma_1(1 - c_{\text{div}}\gamma_3)} \right)^2 + \varepsilon \right) I_{(u,v,Q')} + c(\varepsilon) \int_{Q'} \Phi(|F|) + c \int_{Q'} |T(F,0)|^2 + \varepsilon \int_{Q'} |V(Du)|^2$$

Due to (A3) there exists ε such that it holds that $\left(\frac{\gamma_3 c_{\text{div}}\gamma_2}{\gamma_1(1 - c_{\text{div}}\gamma_3)} \right)^2 + \varepsilon < 1$ and thus

$$I_{(u,v,Q')} \leq c \int_{Q'} (\Phi(|F|) + |T(F,0)|^2) + \varepsilon \int_{Q'} |V(Du)|^2.$$

According to Lemma 2

$$\int_{Q'} |V(Du) - V(Dv)|^2 \leq \varepsilon \int_{Q'} |V(Dv)|^2 + c \int_{Q'} (\Phi(|F|) + |T(F,0)|^2).$$

Further, since $\Phi(|F|) \sim (1 + |F|^2)^{\frac{r-2}{2}}|F|^2 \leq |F|^r$ and $|T(F,0)|^2 \sim (1 + |F|^2)^{r-2}|F|^2 \leq (1 + |F|^2)^{\frac{r-2}{2}} \leq |F|^r$ we get

$$\int_{Q'} |V(Du) - V(Dv)|^2 \leq \varepsilon \int_{Q'} |V(Dv)|^2 + c \int_{Q'} |F|^r.$$

All assumptiuons of Lemma 5 are met and the estimate (1.1) follows directly from (3.1).

Let $\bar{p} = p - \int_Q p$. We consider a test function φ fulfilling

$$\begin{aligned} \text{div } \varphi &= |\bar{p}|^{s'-2}\bar{p} - \int_Q |\bar{p}|^{s'-2}\bar{p} \text{ in } Q \\ \varphi|_{\partial Q} &= 0, \end{aligned}$$

where s is such that $s' = qr'$. Lemma 4 yields

$$\|\nabla\varphi\|_{s,Q} \leq c\|\bar{p}\|_{s',Q}^{1/(s-1)}. \quad (3.7)$$

We multiply $(P)_1$ by this φ in order to get

$$\|\bar{p}\|_{s',Q}^{s'} \leq (\|T(Du,p)\|_{s',Q} + \|T(F,0)\|_{s',Q}) \|\nabla\varphi\|_{s,Q},$$

and, together with (3.7), we obtain

$$\|\bar{p}\|_{s',Q} \leq c (\|T(Du, p)\|_{s',Q} + \|T(F, 0)\|_{s',Q}) \quad (3.8)$$

Further, since $t^{r'} \leq c\Phi^*(t)$, we have

$$|T(Du, p)|^{qr'} \leq c (\Phi^*(\Phi'(|Du|)))^q \sim \left((1 + |Du|^2)^{\frac{r-2}{2}} |Du|^2 \right)^q.$$

This together with (3.8) yield

$$\int_Q |\bar{p}|^{qr'} \leq c \int_Q \left((1 + |Du|^2)^{\frac{r-2}{2}} |Du|^2 \right)^q + c \int_Q \left((1 + |F|^2)^{\frac{r-2}{2}} |F|^2 \right)^q.$$

Using (1.1) one may easily conclude (1.2). □

References

- [1] AMROUCHE, C., GIRAULT, V.: *Decomposition of vector spaces and application to the Stokes problem in arbitrary dimension* Czechoslovak Math. J. 44(119) (1994), no. 1, 109–140. Zbl 0823.35140, MR1257940
- [2] BRIDGMAN, P. W.: *The physics of high pressure*, MacMillan, New York, (1931)
- [3] BULÍČEK, M., FIŠEROVÁ, V.: *Existence theory for steady flows of fluids with pressure and shear rate dependent viscosity, for low values of the power-law index* Z. Anal. Anwend. 28 (2009), no. 3, 349–371. Zbl 1198.35174, MR2506365
- [4] BULÍČEK, M., KAPLICKÝ, P.: *Incompressible fluids with shear rate and pressure dependent viscosity: regularity of steady planar flows* Discrete Contin. Dyn. Syst. Ser. S1 (2008), no. 1, 41–50. Zbl 1153.35314, MR2375580
- [5] BULÍČEK, M., MÁLEK J., RAJAGOPAL, K., R.: *Analysis of the flows of incompressible fluids with pressure dependent viscosity fulfilling $\nu(p, \cdot) \rightarrow \infty$ as $p \rightarrow \infty$* Czech. Math. J. 59 (134) (2009) 503–528. Zbl 1224.35311, MR2532387
- [6] CAFFARELLI, PERAL *On $W^{1,p}$ estimates for elliptic equations in divergence form* Comm. Pure Appl. Math. 51 (1998), no. 1, 1–21. Zbl. 0906.35030, MR1486629
- [7] DIENING, ETTWEIN *Fractional estimates for non-differentiable elliptic systems with general growth* Forum Math. 20 (2008), no. 3, 523–556. Zbl 1188.35069, MR2418205
- [8] DIENING, L., KAPLICKÝ, P.: *L^q theory for a generalized Stokes system* Manuscripta Math. 141 (2013), no. 1–2, 333–361. Zbl 1263.35175, MR3042692
- [9] DIENING, L., RŮŽIČKA, M., SCHUMACHER, K.: *A decomposition technique for John domains* Ann. Acad. Sci. Fenn. Math. 35 (2010), no. 1, 87–114. MR2643399
- [10] FRANTA, M., MÁLEK, J., RAJAGOPAL, K., R.: *On steady flows of fluids with pressure- and shear-dependent viscosities*, Proc. R. Soc. A (2005) **461**, 651–670. Zbl 1145.76311, MR2121929
- [11] GAZZOLA F., SECCHI P.: *Some results about stationary Navier-Stokes equations with a pressure-dependent viscosity. Navier-Stokes equations: theory and numerical methods (Varenna, 1997)*, 31–37, Pitman Res. Notes Math. Ser., 388, Longman, Harlow, 1998. Zbl 0940.35156, MR1773582
- [12] GIUSTI, E.: *Metodi diretti nel calcolo delle variazioni*, Unione Matematica Italiana, Bologna, 1994

- [13] IWANIEC, T.: *On L^p -integrability in PDE's and quasiregular mappings for large exponents*, Ann. Acad. Sci. Fenn. Ser. A I Math **7** (1982), no. 2, 301–322. Zbl 0505.30011, MR0686647
- [14] KNAUF, S., FREI, S., RICHTER, T., RANNACHER, R.: *Towards a complete numerical description of lubricant film dynamics in ball bearings*, Comput. Mech. **53** (2014), no. 2, 239–255. MR3158820
- [15] LANZENDÖRFER, M.: *Numerical simulations of the flow in the journal bearing*, Master's Thesis, Charles University i Prague, Faculty of Mathematics and Physics (2003)
- [16] LANZENDÖRFER, M.: *On steady inner flows of an incompressible fluid with the viscosity depending on the pressure and the shear rate* Nonlinear Analysis: Real World Applications **10** (2009) 1943–1954. Zbl 1163.76335, MR2508405
- [17] LANZENDÖRFER, M., STEBEL, J.: *On pressure boundary conditions for steady flows of incompressible fluids with pressure and shear rate dependent viscosities* Applications of mathematics, **56** (2011) No. 3 , 265–285. Zbl 1224.35347, MR2800578
- [18] MÁCHA, V.: *Partial regularity of solution to generalized Navier-Stokes problem* Cent. Eur. J. Math. **12** (2014), no. 10, 1460–1483. MR3224012
- [19] MÁCHA, V., TICHÝ, J.: *Higher integrability of solution to generalized Stokes system under perfect slip boundary conditions* J. Math. Fluid Mech. **16** (2014), no. 4, 823–845. MR3267551
- [20] MÁLEK, J.: *Mathematical properties of flows of incompressible power-law-like fluids that are described by implicit constitutive relations* Electronic Transactions on Numerical Analysis, Vol. **31** (2008), 110–125. Zbl 1182.35182, MR2569596
- [21] MÁLEK, J., MINGIONE, G., STARÁ, J.: *Fluids With Pressure Dependent Viscosity, Partial Regularity of Steady flows* to appear.
- [22] MÁLEK, J., MINGIONE, G., STARÁ, J.: *Fluids With Pressure Dependent Viscosity, Partial Regularity of Steady Flows*, Equadiff 2003, Proceedings of the International Conference of Differential Equations, 380-385 (2005). MR2185057
- [23] MÁLEK, J., NEČAS, J., RAJAGOPAL, K., R.: *Global analysis of the flows of fluids with pressure-dependent viscosities* Arch. Rational Mech. Anal. **165** (2002) 243–269. Zbl 1022.76011, MR1941479
- [24] ROUBÍČEK, T.: *Nonlinear Partial Differential Equations with Applications* International Series of Numerical Mathematics, **153**. Birkhäuser Verlag, Basel 2005. Zbl 1087.35002, MR2176645