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A thermodynamically consistent phenomenological model for ferroelectric and ferroelastic hysteresis

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We propose a hysteretic model for electromechanical coupling in piezoelectric materials, with the strain and the electric field as inputs and the stress and the polarization as outputs. This constitutive law satisfies the thermodynamic principles and exhibits good agreement with experimental measurements. Moreover, when it is coupled with the mechanical and electromagnetic balance equations, the resulting PDE system is well-posed under the hypothesis that hysteretic effects take place only in one preferred direction. We prove the existence and uniqueness of its global weak solutions for each initial data with prescribed regularity. One of the tools is a new Lipschitz continuity theorem for the inverse Preisach operator with time dependent coefficients.

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1 Introduction

The piezoelectric effect is a coupling between electrical and mechanical fields within certain materials that has numerous applications ranging from ultrasound generation in medical imaging and therapy via acceleration sensors and injection valves in automotive industry to high precision positioning systems. Driven by the increasing demand for devices operating at high field intensities especially in actuator applications, the field of hysteresis modeling for piezoelectric materials is currently one of highly active research. The approaches that have been considered so far can be divided into basically four categories:

- (1) *Thermodynamically consistent models* being based on a macroscopic view to describe microscopic phenomena in such a way that the second law of thermodynamics is satisfied, see for example, Bassiouny and Ghaleb [3], Kamlah and Böhle [17], Landis [25], Schröder and Romanowski [34], Su and Landis [37], Linnemann et al. [26].
- (2) *Micromechanical models* that consider the material on the level of single grains, see, for example, Delibas et al. [9], Fröhlich [10], Huber and Fleck [13], Belov and Kreher [4], Huber [12], McMeeking et al. [28], Smith and Hu [35].
- (3) *Phase field models* that describe the transition between phases (corresponding to the motion of walls between domains with different polarization orientation) using the Ginzburg-Landau equation for some order parameter, see, for example Xu et al. [40], Wang et al. [39].
- (4) *Phenomenological models using hysteresis operators* partly originating from the input-output description of piezoelectric devices for control purposes, see, for example, Hughes and Wen [14], Kuhnen [24], Cima et al. [7], Smith et al. [36], Ball et al. [2], Pasco and Berry [31].

Also multiscale coupling between macro- and microscopic as well as phase field models partly even down to atomistic simulations has been investigated, see, e.g., [29, 33].

For a mathematical analysis of some thermodynamically consistent models we refer to [1, 19, 30].

Whereas most of the so far existing models are designed for the simulation of polarization, depolarization or cycling along the main hysteresis loop, the simulation of actuators requires the accurate simulation of minor loops as well. Moreover, the physical behavior can so far be reproduced only qualitatively, whereas the use of models in actuator simulation (possibly also aiming at simulation based optimization) needs to match measurements precisely. Simulation of a piezoelectric device with a possibly complex geometry requires not only an input-output model but needs to resolve the spatial distribution of

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the crucial electric and mechanical field quantities, which leads to partial differential equations. Therewith, the question of numerical efficiency becomes important.

Preisach operators [5, 18, 22, 27, 38] are phenomenological models for rate independent hysteresis that are capable of reproducing minor loops and can be very well fitted to measurements. Moreover, they allow for a highly efficient evaluation by the application of certain memory deletion rules and the use of so-called Everett or shape functions.

Motivated by these facts, in [11, 15], a model for ferroelectric hysteresis under uniaxial loading using Preisach operators is proposed and studied. Like most of the above mentioned models, it is based on an additive decomposition of the strain and the polarization into reversible and irreversible parts. The reversible quantities follow the linear piezoelectric material law, while the irreversible polarization is represented by a Preisach operator of the imposed electric field and the irreversible strain is a polynomial function of the polarization. Moreover, the piezoelectric coupling coefficient is proportional to the polarization.

However, for the model from [11, 15], it is not clear whether or under which conditions the second law of thermodynamics is satisfied. We will therefore here consider a new material law, which is inspired by the one proposed for magnetostriction [8] whose thermodynamic consistency is based on the use of hysteresis potentials, and which is additionally able to capture ferroelastic effects.

The paper is structured as follows. In Sect. 2, we derive piezoelectric constitutive equations involving a Preisach hysteresis operator from basic thermodynamic principles and show some simulation results with the proposed model. Section 3 is devoted to the proof of a new result on Lipschitz invertibility of the Preisach operator with time dependent coefficients, and in Sect. 4 we construct a unique solution of a system of electromechanical balance equations by Banach contraction principle. We restrict ourselves to the situation of uniaxial loading, hence scalar constitutive relations. A perspective to the situation of electric and mechanic fields depending on three space variables while still considering uniaxial loading, is provided in Sect. 5. The case of vector or actually tensor valued hysteresis will be subject of future research.

2 The model

Differently from most of the above cited approaches, we consider the electric field E and the mechanical strain ε as state variables and the dielectric displacement $D = D(\varepsilon, E)$ as well as the mechanical stress $\sigma = \sigma(\varepsilon, E)$ as state functions. The reason for doing so lies in the fact that the balance equations (namely Newton's law on the mechanical side and Gauss law on the electric side)

$$\rho \ddot{u} - \nabla_s^T \sigma = 0 \quad (1)$$

$$-\nabla \cdot D = 0 \quad (2)$$

(where u is the mechanical displacement, ∇_s the symmetric gradient and ∇_s^T the dyadic divergence) are naturally formulated in terms of these state functions.

For the interrelation between these quantities, we assume, similarly to [8], that hysteresis effects are due to one single Preisach operator \mathcal{P} with potential \mathcal{U} acting on an auxiliary state function $q = q(\varepsilon, E)$. In the scalar case, this leads to the assumption that the stress σ , the dielectric displacement D , and the free energy $F = F(\varepsilon, E)$ are of the form

$$\sigma = \underline{c}\varepsilon - eE + a\mathcal{P}[q] + b\mathcal{U}[q], \quad (3)$$

$$D = e\varepsilon + \kappa E + c\mathcal{P}[q] + d\mathcal{U}[q], \quad (4)$$

$$F = \frac{c}{2}\varepsilon^2 + \frac{\kappa}{2}E^2 + \xi\mathcal{P}[q] + \eta\mathcal{U}[q], \quad (5)$$

where the coefficients $a = a(\varepsilon, E)$, $b = b(\varepsilon, E)$, $c = c(\varepsilon, E)$, $d = d(\varepsilon, E)$, $\xi = \xi(\varepsilon, E)$, $\eta = \eta(\varepsilon, E)$ as well as the function $q = q(\varepsilon, E)$, are to be determined in agreement with the principles of thermodynamics. The elastic, dielectric, and piezoelectric coupling coefficients $\underline{c} > 0$, $e \in \mathbb{R}$, and $\kappa > 0$ are given constants. It will be shown in Sect. 3 that the Preisach hysteresis operator with nonnegative density satisfies the energy inequality

$$q(t) \frac{d}{dt} \mathcal{P}[q](t) - \frac{d}{dt} \mathcal{U}[q](t) \geq 0 \quad \text{a.e.} \quad \forall q \in W^{1,1}(0, T). \quad (6)$$

Similarly to Sect. 3.4 in [8], we now derive conditions to ensure thermodynamic admissibility, namely the requirement that

$$\dot{\varepsilon}\sigma + \dot{D}E - \dot{F} \geq 0 \quad (7)$$

has to hold for all processes $\varepsilon, \sigma, D, E$ that obey the constitutive laws (3), (4). We obtain

$$0 \leq (a\dot{\varepsilon} + E\dot{c} - \dot{\xi})\mathcal{P}[q] + (b\dot{\varepsilon} + E\dot{d} - \dot{\eta})\mathcal{U}[q] + (Ec - \xi)\frac{d}{dt}\mathcal{P}[q] + (Ed - \eta)\frac{d}{dt}\mathcal{U}[q]. \quad (8)$$

Both $\mathcal{P}[q]$ and $\mathcal{U}[q]$ are hysteresis operators and may take arbitrary values independent of each other, so that we have to demand

$$\begin{aligned} 0 &= a\dot{\varepsilon} + E\dot{c} - \dot{\xi} = \dot{\varepsilon} \left(a + E \frac{\partial c}{\partial \varepsilon} - \frac{\partial \xi}{\partial \varepsilon} \right) + \dot{E} \left(E \frac{\partial c}{\partial E} - \frac{\partial \xi}{\partial E} \right), \\ 0 &= b\dot{\varepsilon} + E\dot{d} - \dot{\eta} = \dot{\varepsilon} \left(b + E \frac{\partial d}{\partial \varepsilon} - \frac{\partial \eta}{\partial \varepsilon} \right) + \dot{E} \left(E \frac{\partial d}{\partial E} - \frac{\partial \eta}{\partial E} \right). \end{aligned} \quad (9)$$

These relations have to hold for all processes, hence the coefficients of $\dot{\varepsilon}$ and \dot{E} have to vanish identically. In other words, if we set $f_1 = \eta - Ed$, $f_2 = Ec - \xi$, we must have

$$a = -\frac{\partial f_2}{\partial \varepsilon}, \quad b = \frac{\partial f_1}{\partial \varepsilon}, \quad c = \frac{\partial f_2}{\partial E}, \quad d = -\frac{\partial f_1}{\partial E}. \quad (10)$$

Inequality (8) then becomes

$$0 \leq f_2 \frac{d}{dt}\mathcal{P}[q] - f_1 \frac{d}{dt}\mathcal{U}[q], \quad (11)$$

and will be satisfied, by virtue of (6), provided we choose $q := f_2/f_1$ with $f_1 > 0$.

We see that the whole model depends on the choice of two state functions $f_1 = f_1(\varepsilon, E) > 0$ and $f_2 = f_2(\varepsilon, E)$ which characterize the material properties. A canonical choice, which is sufficient in many situations, consists in putting $f_2(\varepsilon, E) = E$, $f_1(\varepsilon, E) = f(\varepsilon)$ with a suitable function f of one variable. Then the constitutive law (3)–(5) becomes

$$\sigma = \underline{c}\varepsilon - eE + f'(\varepsilon)\mathcal{U}\left[\frac{E}{f(\varepsilon)}\right], \quad (12)$$

$$D = e\varepsilon + \kappa E + \mathcal{P}\left[\frac{E}{f(\varepsilon)}\right], \quad (13)$$

$$F = \frac{c}{2}\varepsilon^2 + \frac{\kappa}{2}E^2 + f(\varepsilon)\mathcal{U}\left[\frac{E}{f(\varepsilon)}\right]. \quad (14)$$

The full 1D system for unknown functions $u(x, t)$, $E(x, t)$, $(x, t) \in (0, \ell) \times (0, T)$, describing longitudinal oscillations of a piezoelectric beam, then reads

$$\begin{aligned} \rho u_{tt} - \sigma_x &= 0, \\ D_x &= 0, \\ \varepsilon &= u_x. \end{aligned} \quad (15)$$

The equation $D_x = 0$ means that D is a function of t only, say, $D(x, t) = r(t)$, that is,

$$e\varepsilon + \kappa E + \mathcal{P}\left[\frac{E}{f(\varepsilon)}\right] = r(t), \quad (16)$$

where $r(t)$ is a function which is known from the boundary condition $D(0, t) = D(\ell, t) = r(t)$, corresponding to an impressed (or measured) boundary current. Furthermore, we complement the mechanical constitutive law (12) with a viscosity term $\nu\varepsilon_t$, where $\nu > 0$ is the viscosity coefficient, that is,

$$\sigma = \nu\varepsilon_t + \underline{c}\varepsilon - eE + f'(\varepsilon)\mathcal{U}\left[\frac{E}{f(\varepsilon)}\right] \quad (17)$$

instead of (12). We prescribe some boundary conditions, for example $u = 0$ on $x = 0$, and $\sigma = s(t)$ on $x = \ell$, which correspond to the experimental setting of a beam which is clamped at the left tip, with an impressed (or measured) force at the right beam tip. We resume the analysis of this system in Sect. 4 and before, in Sect. 3, we establish some new properties of the Preisach operator which will enable us to eliminate E from the system by solving Eq. (16) independently with respect to $q = E/f(\varepsilon)$.

Remark 2.1. We wish to mention that independence of the dielectric saturation value on the stress (which is the case e.g. for the setting $c = 1$, $d = 0$ considered in Sect. 3.4.1 of [8]) physically makes sense also here: Note that saturation of the polarization, i.e., of $P = \mathcal{P}[q]$ taking its maximal absolute value, corresponds to the situation of all c-axes (and

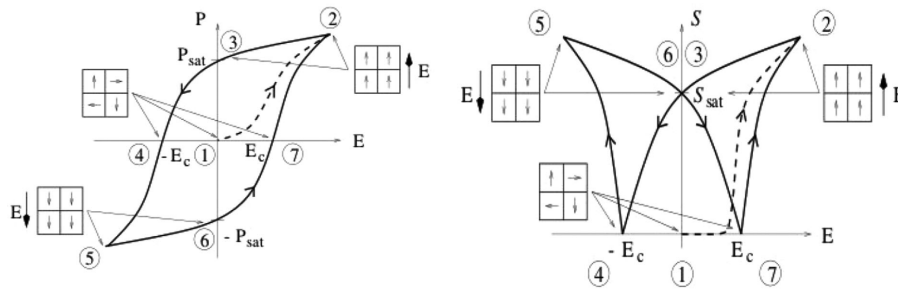


Fig. 1 Polarization hysteresis (left) and strain (butterfly) hysteresis (right) under a bipolar electric field excitation; taken from [16].

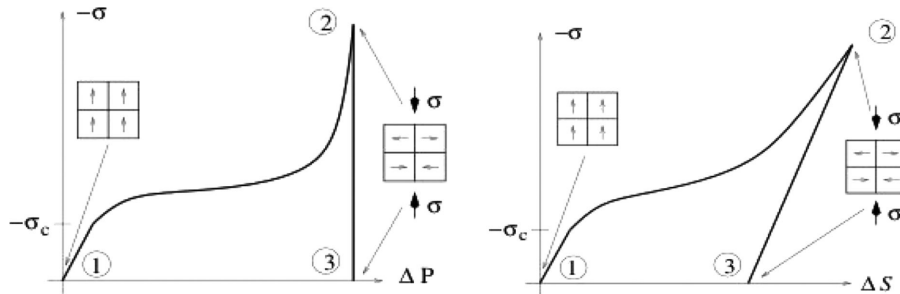


Fig. 2 Mechanical depolarization (left) and stress-strain relation (right) under compressive stress load; taken from [16].

therewith all elementary dipoles) of the single crystals being aligned as much as possible to the load axis, under the given geometric constraints (note that each grain has its own coordinate system of preferred directions). This maximal alignment gives the same polarization, independent of the imposed stress. In other words, no matter how large the imposed stress is, there exists a sufficiently large impressed electric field (in load axis direction) that will bring the polarization to its maximal value by aligning all elementary dipoles (and therewith all c-axes) as much as possible in load direction.

We conclude this section with some simulation results for a very simple choice of the Preisach operator \mathcal{P} and the function f , see Figs. 3, 4, 5, 6, in order to illustrate that the expected qualitative behavior of ferroelectric and ferroelastic hysteresis (see Figs. 1, 2 taken from [16]) can indeed be recovered. Here we use the functions

$$g(r, s) = \begin{cases} \text{proj}_{[-1+r, 1-r]}(s) = \max\{-1+r, \min\{1-r, s\}\} & \text{if } r \leq 1 \\ 0 & \text{else} \end{cases} \quad (18)$$

in the Definition 3.2 of \mathcal{P} , as well as

$$f(x) = 1.1 - x \quad -1 \leq x \leq 1 \quad (19)$$

or

$$f(x) = \frac{1}{2} + \frac{(x-1)^4}{4} \quad -1 \leq x \leq 1 \quad (20)$$

(in both cases the extension to $\mathbb{R} \setminus [-1, 1]$ is done such that Hypothesis 3.3 is satisfied) $c = 1$, $e = 0$, $\kappa = 0.01$ in (14), and normalize all input quantities to the unit interval $[-1, 1]$. In Figures 3, 4, 5, 6, increasing line thickness indicates proceeding time in order to show that the curves are traversed in the right direction. Of course, also quantitative agreement with measurement can be expected to be achievable by appropriate fitting methods.

3 Inversion of hysteresis operators with time dependent coefficients

In this section, we prove Lipschitz continuity statements for mappings which with given functions $w, b \in C[0, T]$ associate the solution q of the equation

$$q(t) + b(t)\mathcal{P}[q](t) = w(t) \quad \forall t \in [0, T], \quad (21)$$

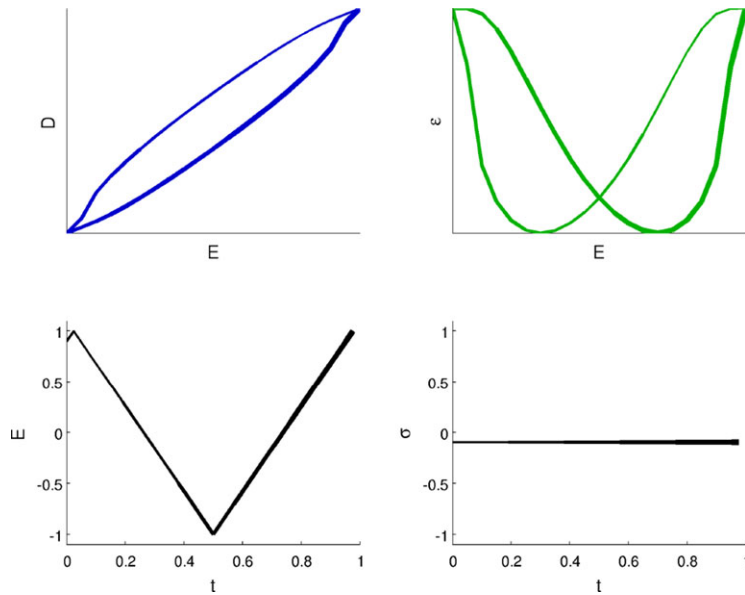


Fig. 3 Polarization hysteresis (top row left) and strain (butterfly) hysteresis (top row right) under a bipolar electric field excitation (bottom row); simulations with model (14) where f is chosen according to (19), and \mathcal{P} according to Definition 3.2 with (18).

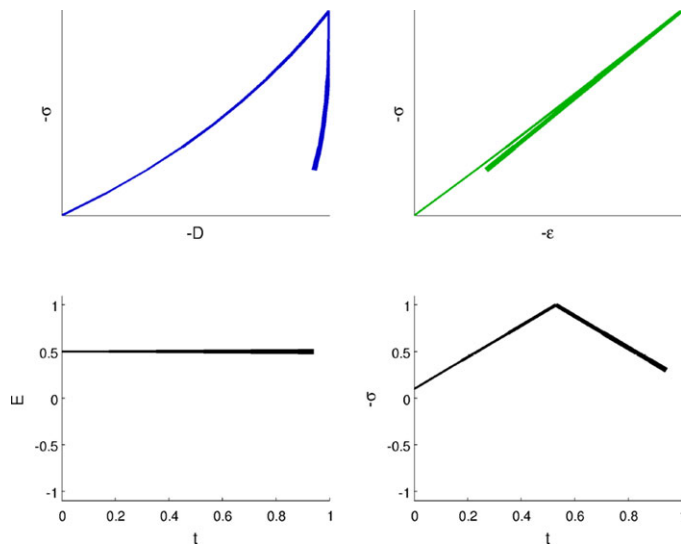


Fig. 4 Mechanical depolarization (left) and stress-strain relation (right) under compressive stress load (bottom row); simulations with model (14) where f is chosen according to (19), and \mathcal{P} according to Definition 3.2 with (18).

where $\mathcal{P} : C[0, T] \rightarrow C[0, T]$ is a Lipschitz continuous operator with respect to the sup norm

$$\|v\| = \max_{\tau \in [0, T]} |v(\tau)|,$$

more precisely in the following time causal sense

$$|\mathcal{P}[q_1](t) - \mathcal{P}[q_2](t)| \leq M \max_{\tau \in [0, t]} |q_1(\tau) - q_2(\tau)| \tag{22}$$

for some $M > 0$ and for all $q_1, q_2 \in C[0, T]$.

It follows from (22) that \mathcal{P} has the Volterra property, i.e., for all $t \in [0, T]$

$$\forall q_1, q_2 \in C[0, T] : q_1 = q_2 \text{ on } [0, t] \text{ implies } \mathcal{P}[q_1] = \mathcal{P}[q_2] \text{ on } [0, t].$$

We first specify sufficient conditions under which is the mapping $(w, b) \mapsto q$ well defined.

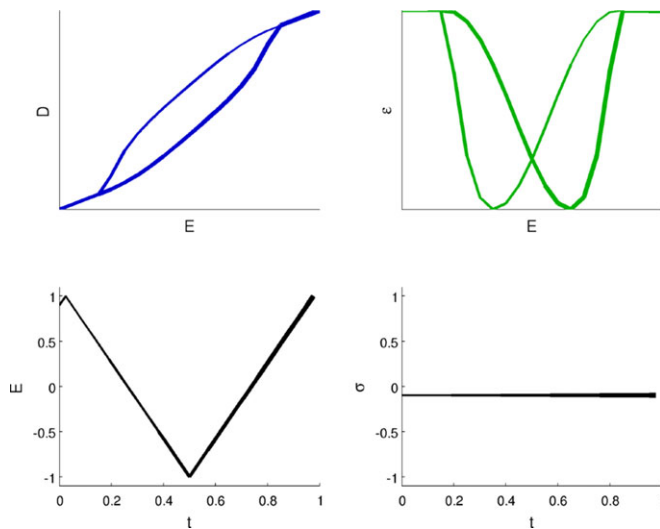


Fig. 5 Polarization hysteresis (top row left) and strain (butterfly) hysteresis (top row right) under a bipolar electric field excitation (bottom row); simulations with model (14) where f is chosen according to (20), and \mathcal{P} according to Definition 3.2 with (18).

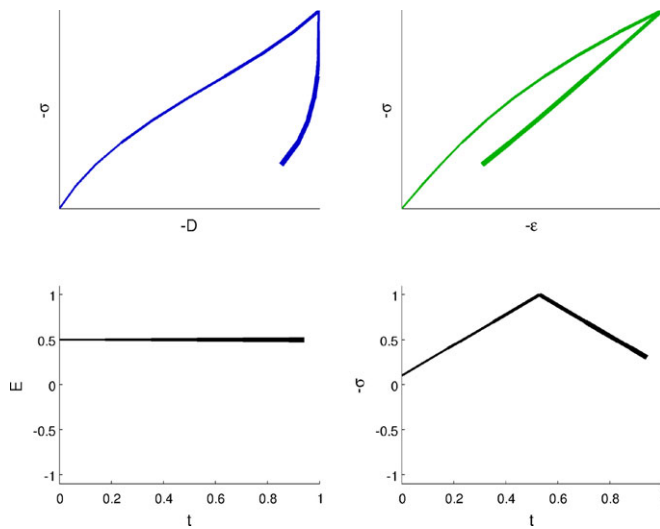


Fig. 6 Mechanical depolarization (left) and stress-strain relation (right) under compressive stress load (bottom row); simulations with model (14) where f is chosen according to (20), and \mathcal{P} according to Definition 3.2 with (18).

Lemma 3.1. *Let \mathcal{P} be as in (22), and such that $I + c\mathcal{P}$ admits for every $c \geq 0$ a Lipschitz continuous inverse in the sense that the inequality*

$$|(I + c\mathcal{P})^{-1}[w_1](t) - (I + c\mathcal{P})^{-1}[w_2](t)| \leq L \max_{\tau \in [0,t]} |w_1(\tau) - w_2(\tau)| \tag{23}$$

holds for all $q_1, q_2 \in C[0, T]$ with Lipschitz constant $L > 0$ independent of c . Let $w, b \in C[0, T]$ be given functions, and let $b(t) \geq 0$ for all $t \in [0, T]$. Then there exists a unique $q \in C[0, T]$ such that (21) holds.

The Lipschitz continuity of $(I + c\mathcal{P})^{-1}$ independently of the constant c is fulfilled in many typical situations arising in the theory of hysteresis operators. If \mathcal{P} is the Preisach operator defined below in Definition 3.2, it was shown in [6] that (23) is satisfied with $L = 2$. If \mathcal{P} belongs to the subclass of Prandtl-Ishlinskii operators, then $(I + c\mathcal{P})^{-1}$ admits in addition an explicit representation, see [20].

Proof. We choose $\gamma < 1/(ML)$ and a division $0 = t_0 < t_1 < \dots < t_m = T$ of $[0, T]$ such that $|b(t) - b(t_{k-1})| < \gamma$ for $t \in [t_{k-1}, t_k], k = 1, \dots, m$. We prove by induction that

$$\forall k \in \{1, \dots, m\} \exists! q \in C[0, t_k] : q \text{ satisfies (21).}$$

In $[0, t_1]$, we rewrite (21) as

$$q(t) + b(0)\mathcal{P}[q](t) = w(t) - (b(t) - b(0))\mathcal{P}[q](t),$$

or, equivalently,

$$q(t) = (I + b(0)\mathcal{P})^{-1}[w - (b(\cdot) - b(0))\mathcal{P}[q]](t). \quad (24)$$

By our assumptions of Lipschitz continuity and the choice of γ , the mapping on the right hand side of (24) is a contraction with respect to q , hence it admits a unique fixed point.

Assume now that we have constructed a unique $q(t)$ satisfying (21) on $[0, t_{k-1}]$. We introduce the set

$$V_k = \{\hat{q} \in C[0, t_k] : \hat{q}(t) = q(t) \text{ on } [0, t_{k-1}]\}.$$

We define a mapping $R_k : V_k \rightarrow V_k$ which with each $\hat{q} \in V_k$ associates the solution q of the equation

$$q(t) + b(t_{k-1})\mathcal{P}[q](t) = w(t) - (b(t) - b(t_{k-1}))\mathcal{P}[\hat{q}](t) \quad (25)$$

for $t \in [0, t_k]$. The mapping R_k is again a contraction on V_k (by our assumptions of Lipschitz continuity, the choice of γ , as well as the Volterra property) and therefore admits a unique fixed point. The induction argument over k then completes the proof. \square

Here, we focus on the case of Preisach operators. We use the following definition which is shown in [21] to be equivalent to the original Preisach construction in [32].

Definition 3.2. Let $g : (0, \infty) \times \mathbb{R} \rightarrow \mathbb{R}$ be a measurable function which is Lipschitz continuous in the second variable and such that $g(r, 0) = 0$ for a. e. $r > 0$. For a given input $q \in W^{1,1}(0, T)$, we define the output $\mathcal{P}[q] \in W^{1,1}(0, T)$ of a Preisach operator \mathcal{P} by the integral

$$\mathcal{P}[q](t) = \int_0^\infty g(r, \xi_r(t)) \, dr \quad \text{for } t \in [0, T], \quad (26)$$

where $\xi_r \in W^{1,1}(0, T)$ is the unique solution of the variational inequality

$$\begin{cases} |q(t) - \xi_r(t)| \leq r & \forall t \in [0, T], \\ \dot{\xi}_r(t)(q(t) - \xi_r(t) - rz) \geq 0 & \text{a. e. } \forall |z| \leq 1, \\ \xi_r(0) = \max\{q(0) - r, \min\{0, q(0) + r\}\}. \end{cases} \quad (27)$$

The Preisach operator \mathcal{P} is called a Prandtl-Ishlinskii operator if g is linear in v , that is, $g(r, v) = \mu(r)v$ for some $\mu \in L^1_{loc}(0, \infty)$.

We easily check from (27) that $\xi_r(t) = 0$ for $r \geq \|q\|$, where $\|\cdot\|$ denotes the sup-norm in $C[0, T]$, so that the integral in (26) is meaningful.

The parameter r is the memory variable, and the mapping $q \mapsto \xi_r$ introduced in [18] is called the play operator. Note that its extension to $C[0, T]$ is Lipschitz continuous with Lipschitz constant 1, that is,

$$|\xi_r^1(t) - \xi_r^2(t)| \leq \max_{\tau \in [0, t]} |q_1(\tau) - q_2(\tau)|. \quad (28)$$

For our purposes, it is convenient to reduce the set of admissible functions g , and we adopt the following hypothesis.

Hypothesis 3.3. We assume that $g(r, 0) = 0$ for a. e. $r > 0$, and

- (i) $0 \leq \frac{\partial g}{\partial v}(r, v) \leq \mu(r)$ a. e., where $\mu \in L^1(0, \infty)$, $\int_0^\infty \mu(r) \, dr = M$;
- (ii) There exists $\infty > M_1 > 0$ such that

$$\int_0^\infty \int_{-\infty}^\infty \frac{\partial g}{\partial v}(r, v) \, dv \, dr = M_1. \quad (29)$$

Hypothesis 3.3 (i) means that the Preisach operator \mathcal{P} is dominated by a Prandtl-Ishlinskii operator. Hypothesis 3.3 (ii) implies that $\|\mathcal{P}[q]\| \leq M_1$ for all $q \in C[0, T]$. It is easy to check that even without Hypothesis 3.3 (i), (ii), the Preisach operator satisfies the energy inequality (6) with the choice

$$\mathcal{U}[q](t) = \int_0^\infty G(r, \xi_r(t)) \, dr, \quad (30)$$

with

$$G(r, v) = \int_0^v v' \frac{\partial g}{\partial v}(r, v') \, dv'. \tag{31}$$

Moreover, it satisfies the Volterra property by virtue of (28). Still by (28), the operator \mathcal{P} can be extended to $C[0, T]$, and if Hypothesis 3.3 (i) is fulfilled, then the Lipschitz property (22) holds.

The main result of this section reads as follows.

Theorem 3.4. *Let Hypothesis 3.3 (i) hold, and let $b, w_1, w_2 \in C[0, T]$ be given, $0 \leq b(t) \leq \bar{b}$ for all $t \in [0, T]$. Let $q_1, q_2 \in C[0, T]$ be such that*

$$q_i(t) + b(t)\mathcal{P}[q_i](t) = w_i(t) \quad \forall t \in [0, T], \quad i = 1, 2. \tag{32}$$

Then we have

$$\|q_1 - q_2\| \leq e^{\bar{b}M} \|w_1 - w_2\|. \tag{33}$$

The proof of Theorem 3.4 is divided into several steps. Indeed, as $W^{1,1}(0, T)$ is dense in $C[0, T]$, it is enough to prove (33) for $q_i \in W^{1,1}(0, T)$, and we repeatedly use the variational inequality (27). The following observation is due to Brokate, see [5], and is related to the Brokate identity in more general situations, cf. [23]. We give a full proof here, also because it uses techniques that will play a substantial role in the proof of Theorem 3.4.

Lemma 3.5. *Let $r_2 > r_1 > 0$ and $q \in W^{1,1}(0, T)$ be given, and let $\xi_{r_i}, i = 1, 2$, be the solution of (27). Then we have*

$$\begin{cases} |\dot{\xi}_{r_1}(t) - \dot{\xi}_{r_2}(t)| \leq r_2 - r_1 & \forall t \in [0, T], \\ \dot{\xi}_{r_2}(t)(\xi_{r_1}(t) - \xi_{r_2}(t) - (r_2 - r_1)z) \geq 0 & \text{a. e. } \forall |z| \leq 1. \end{cases} \tag{34}$$

Proof. Let $\eta \in W^{1,1}(0, T)$ be the solution of the variational inequality

$$\begin{cases} |\dot{\xi}_{r_1}(t) - \eta(t)| \leq r_2 - r_1 & \forall t \in [0, T], \\ \dot{\eta}(t)(\xi_{r_1}(t) - \eta(t) - (r_2 - r_1)z) \geq 0 & \text{a. e. } \forall |z| \leq 1 \\ \eta(0) = \max\{\xi_{r_1}(0) - (r_2 - r_1), \min\{0, \xi_{r_1}(0) + (r_2 - r_1)\}\}. \end{cases} \tag{35}$$

Putting in the above inequality $z = \frac{\xi_{r_1}(t \pm \delta) - \eta(t \pm \delta)}{r_2 - r_1}$ and letting δ tend to $0+$ we obtain

$$\dot{\eta}(t)(\dot{\xi}_{r_1}(t) - \dot{\eta}(t)) = 0 \quad \text{a. e.} \tag{36}$$

We have by (27) for $r = r_1$ that $\dot{\xi}_{r_1}(t)(q(t) - \xi_{r_1}(t) - r_1z) \geq 0$ a. e. for all $|z| \leq 1$, and (36) implies the implication $\dot{\eta}(t) \neq 0 \implies \dot{\eta}(t) = \dot{\xi}_{r_1}(t)$. Hence,

$$\dot{\eta}(t)(q(t) - \xi_{r_1}(t) - r_1z) \geq 0 \text{ a. e. } \forall |z| \leq 1. \tag{37}$$

We have $|q(t) - \eta(t)| \leq r_2$, and adding (37) to (35) we obtain

$$\dot{\eta}(t)(q(t) - \eta(t) - r_2z) \geq 0 \text{ a. e. } \forall |z| \leq 1. \tag{38}$$

By (27) for $r = r_2$ we get that

$$\dot{\xi}_{r_2}(t)(q(t) - \xi_{r_2}(t) - r_2z) \geq 0 \text{ a. e. } \forall |z| \leq 1. \tag{39}$$

It is easy to check that $\eta(0) = \xi_{r_2}(0)$, hence, comparing (39) with (38), we conclude that $\eta(t) = \xi_{r_2}(t)$ for all t , which we wanted to prove. \square

Next, we consider *discrete Preisach operators*

$$\mathcal{P}_k[q] = \sum_{j=1}^k g_j(\xi_{r_j}) \quad \text{for } k \in \mathbb{N}, \tag{40}$$

corresponding to a sequence $0 = r_0 < r_1 < \dots$ and an associated sequence $\{g_j\}_{j \in \mathbb{N}}$ of nondecreasing functions.

Proposition 3.6. *Let $b \in C[0, T]$ be as in Theorem 3.4, and let $\{g_j\}_{j \in \mathbb{N}}$ be a sequence of Lipschitz continuous functions in \mathbb{R} such that $g_j(0) = 0$ and $0 \leq g'_j(v) \leq \mu_j$ for all $j \in \mathbb{N}$ and a. e. $v \in \mathbb{R}$, where $\mu_j \geq 0$ are constants. Let $q_1, q_2, w_1, w_2 \in C[0, T]$ be such that*

$$q_i(t) + b(t)\mathcal{P}_k[q_i](t) = w_i(t) \quad \forall t \in [0, T], \quad i = 1, 2 \tag{41}$$

for some $k \in \mathbb{N}$. Then

$$\|q_1 - q_2\| \leq \rho_k \|w_1 - w_2\|, \quad (42)$$

where $\rho_k = \prod_{j=1}^k (1 + \bar{b}\mu_j)$.

Proof. As mentioned above, the proof will be carried out for $q_1, q_2 \in W^{1,1}(0, T)$. We proceed by induction over k . For $k = 1$, Eq. (41) reads

$$q_i(t) + b(t)g_1(\xi_{r_1}^i)(t) = w_i(t) \quad \forall t \in [0, T], \quad i = 1, 2, \quad (43)$$

where $\xi_{r_1}^i$ is the solution of (27) with input $q = q_i$ and memory level $r = r_1$. In (27) we choose $z = (q_2 - \xi_{r_1}^2)/r_1$ for $i = 1$ and $z = (q_1 - \xi_{r_1}^1)/r_1$ for $i = 2$, and obtain

$$\dot{\xi}_{r_1}^1(t)(q_1(t) - \xi_{r_1}^1(t) - q_2(t) + \xi_{r_1}^2(t)) \geq 0 \quad \text{a.e.}, \quad (44)$$

$$\dot{\xi}_{r_1}^2(t)(q_2(t) - \xi_{r_1}^2(t) - q_1(t) + \xi_{r_1}^1(t)) \geq 0 \quad \text{a.e.} \quad (45)$$

The inequalities remain to hold when we multiply (44) by $g_1'(\xi_{r_1}^1(t))$ and (45) by $g_1'(\xi_{r_1}^2(t))$. We sum up the two and obtain

$$\frac{d}{dt}(g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2))(t)((q_1 - q_2) - (\xi_{r_1}^1 - \xi_{r_1}^2))(t) \geq 0 \quad \text{a.e.} \quad (46)$$

For $t \in [0, T]$ put

$$V_1(t) = \max \left\{ (g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2))^2(t), (\mu_1 \|w_1 - w_2\|)^2 \right\}. \quad (47)$$

We claim that

$$\dot{V}_1(t) \leq 0 \quad \text{a.e.} \quad (48)$$

To prove (48), we proceed by contradiction. Assume that there exists a Lebesgue point $t \in (0, T)$ of both $\xi_{r_1}^1, \xi_{r_1}^2$ such that $\dot{V}_1(t) > 0$. Then we have

$$|g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2)|(t) > \mu_1 \|w_1 - w_2\|, \quad (49)$$

$$\frac{d}{dt}(g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2))(t)(g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2))(t) > 0. \quad (50)$$

It follows from (43), (46), and (50) that

$$(g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2))(t)((w_1 - w_2)(t) - b(t)(g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2))(t) - (\xi_{r_1}^1 - \xi_{r_1}^2)(t)) \geq 0, \quad (51)$$

and since g_1 is monotone and $\mu_1 |\xi_{r_1}^1 - \xi_{r_1}^2| \geq |g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2)|$, the estimate

$$\begin{aligned} (g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2))(t)(\xi_{r_1}^1 - \xi_{r_1}^2)(t) &= |g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2)|(t) |\xi_{r_1}^1 - \xi_{r_1}^2|(t) \\ &\geq \frac{1}{\mu_1} (g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2))^2(t) \end{aligned}$$

holds, hence we obtain

$$|g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2)|(t) \leq \frac{\mu_1}{1 + b(t)\mu_1} \|w_1 - w_2\|(t) \leq \mu_1 \|w_1 - w_2\|, \quad (52)$$

in contradiction with (49). Hence, (48) holds, $V_1(t)$ is nonincreasing, and we have by (47) for every $t \in [0, T]$ that

$$|g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2)|(t) \leq \max \left\{ |g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2)|(0), \mu_1 \|w_1 - w_2\| \right\}. \quad (53)$$

The monotonicity of g_1 and of the initial value mapping in (27) now yield

$$\|w_1 - w_2\|(0) = |q_1 - q_2|(0) + b(0)|g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2)|(0),$$

hence

$$|g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2)|(0) \leq \mu_1 |\xi_{r_1}^1 - \xi_{r_1}^2|(0) \leq \mu_1 |q_1 - q_2|(0) \leq \mu_1 \|w_1 - w_2\|(0),$$

and we conclude from (43) and (53) that

$$\|q_1 - q_2\| \leq \bar{b} \|g_1(\xi_{r_1}^1) - g_1(\xi_{r_1}^2)\| + \|w_1 - w_2\| \leq (1 + \bar{b}\mu_1) \|w_1 - w_2\|, \quad (54)$$

and the first induction step for $k = 1$ is done.

Let now $k > 1$ be arbitrary and assume that (42) holds for $k - 1$ in place of k . Eq. (41) can be written as

$$q_i(t) + b(t)\mathcal{P}_{k-1}[q_i](t) = w_i(t) - b(t)g_k(\xi_{r_k}^i(t)) \quad \forall t \in [0, T], \quad i = 1, 2, \tag{55}$$

and by induction hypothesis we have

$$\|q_1 - q_2\| \leq \rho_{k-1} (\|w_1 - w_2\| + \bar{b}\|g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2)\|). \tag{56}$$

By the choices $z = \frac{q_i(t) - \xi_{r_k}^i(t)}{r_k}$, $i = 1, 2$ in (27) and $z = \frac{\xi_{r_j}^i(t) - \xi_{r_k}^i(t)}{r_k - r_j}$, $i = 1, 2$ in Lemma 3.5, we obtain similarly as in (44)–(45)

$$\dot{\xi}_{r_k}^1(t) ((q_1(t) - q_2(t)) - (\xi_{r_k}^1(t) - \xi_{r_k}^2(t))) \geq 0 \quad \text{a.e.}, \tag{57}$$

$$\dot{\xi}_{r_k}^2(t) ((q_2(t) - q_1(t)) - (\xi_{r_k}^2(t) - \xi_{r_k}^1(t))) \geq 0 \quad \text{a.e.}, \tag{58}$$

$$\dot{\xi}_{r_k}^1(t) ((\xi_{r_j}^1(t) - \xi_{r_j}^2(t)) - (\xi_{r_k}^1(t) - \xi_{r_k}^2(t))) \geq 0 \quad \text{a.e. for } j = 1, \dots, k - 1, \tag{59}$$

$$\dot{\xi}_{r_k}^2(t) ((\xi_{r_j}^2(t) - \xi_{r_j}^1(t)) - (\xi_{r_k}^2(t) - \xi_{r_k}^1(t))) \geq 0 \quad \text{a.e. for } j = 1, \dots, k - 1. \tag{60}$$

The same argument as in the transition from (44)–(45) to (46) yields

$$\frac{d}{dt} (g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) ((q_1 - q_2) - (\xi_{r_k}^1 - \xi_{r_k}^2))(t) \geq 0 \quad \text{a.e.}, \tag{61}$$

$$\frac{d}{dt} (g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) ((\xi_{r_j}^1 - \xi_{r_j}^2) - (\xi_{r_k}^1 - \xi_{r_k}^2))(t) \geq 0 \quad \text{a.e.} \tag{62}$$

for every $j = 1, \dots, k - 1$ in (62). We continue as above and define the function

$$V_k(t) = \max \left\{ (g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))^2(t), (\mu_k \|w_1 - w_2\|)^2 \right\} \tag{63}$$

with the goal to prove that

$$\dot{V}_k(t) \leq 0 \quad \text{a.e.} \tag{64}$$

Assume that there exists a Lebesgue point $t \in (0, T)$ of both $\dot{\xi}_{r_k}^1, \dot{\xi}_{r_k}^2$ such that $\dot{V}_k(t) > 0$. Then we have

$$|g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2)|(t) > \mu_k \|w_1 - w_2\|, \tag{65}$$

$$\frac{d}{dt} (g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) (g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) > 0. \tag{66}$$

It follows from (61)–(62) that

$$(g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) ((q_1 - q_2) - (\xi_{r_k}^1 - \xi_{r_k}^2))(t) \geq 0, \tag{67}$$

$$(g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) ((\xi_{r_j}^1 - \xi_{r_j}^2) - (\xi_{r_k}^1 - \xi_{r_k}^2))(t) \geq 0, \quad j = 1, \dots, k - 1. \tag{68}$$

As a consequence of (68) and the monotonicity of g_k , we have

$$(g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) (\xi_{r_j}^1 - \xi_{r_j}^2)(t) \geq 0 \quad \text{for } j = 1, \dots, k - 1, \tag{69}$$

and the monotonicity of g_j yields in turn

$$(g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) (g_j(\xi_{r_j}^1) - g_j(\xi_{r_j}^2))(t) \geq 0 \quad \text{for } j = 1, \dots, k - 1. \tag{70}$$

Combining (67) with (70) we obtain

$$(g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) \left(\left(q_1 + b \sum_{j=1}^{k-1} g_j(\xi_{r_j}^1) \right) - \left(q_2 + b \sum_{j=1}^{k-1} g_j(\xi_{r_j}^2) \right) - (\xi_{r_k}^1 - \xi_{r_k}^2) \right) (t) \geq 0, \tag{71}$$

which is nothing but (cf. (55))

$$(g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) ((w_1 - w_2)(t) - b(t)(g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2))(t) - (\xi_{r_k}^1 - \xi_{r_k}^2)(t)) \geq 0. \tag{72}$$

Using the monotonicity of g_k and inequality $\mu_k |\xi_{r_k}^1 - \xi_{r_k}^2| \geq |g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2)|$, we thus have like in (52)

$$|g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2)|(t) \leq \frac{\mu_k}{1 + b(t)\mu_k} |w_1 - w_2|(t) \leq \mu_k \|w_1 - w_2\| \tag{73}$$

in contradiction with (65). Hence, (64) holds, $V_k(t)$ is nonincreasing, and we have by (63) for every $t \in [0, T]$ that

$$|g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2)|(t) \leq \max \{ |g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2)|(0), \mu_k \|w_1 - w_2\| \}. \tag{74}$$

The monotonicity of g_k and of the initial value mapping in (27) now yield

$$|g_k(\xi_{r_k}^1) - g_k(\xi_{r_k}^2)|(0) \leq \mu_k |\xi_{r_k}^1 - \xi_{r_k}^2|(0) \leq \mu_k |q_1 - q_2|(0) \leq \mu_k \|w_1 - w_2\|(0),$$

and we conclude from (56) and (74) that

$$\|q_1 - q_2\| \leq \rho_{k-1} (1 + \bar{b}\mu_k) \|w_1 - w_2\| = \rho_k \|w_1 - w_2\|, \tag{75}$$

which completes the induction argument. □

We are now ready to prove Theorem 3.4.

Proof. (Theorem 3.4) Let b, q_1, q_2, w_1, w_2 be as in (32). We choose $R > \max\{\|q_1\|, \|q_2\|\}$, so that $\xi_r^i(t) = 0$ for all $r \geq R, t \in [0, T]$, and $i = 1, 2$. For every $\delta > 0$, we choose a partition $0 = r_0 < r_1 < \dots < r_m = R$ with $r_j - r_{j-1} < \delta$. Our goal is to approximate \mathcal{P} by \mathcal{P}_m of the form (40) by putting

$$g_j(v) = \int_{r_{j-1}}^{r_j} g(r, v) \, dr \quad \text{for } v \in \mathbb{R}, \quad j = 1, \dots, m.$$

We have for $i = 1, 2$ that

$$|\mathcal{P}[q_i] - \mathcal{P}_m[q_i]| = \left| \int_0^\infty g(r, \xi_r^i) \, dr - \sum_{j=1}^m g_j(\xi_{r_j}^i) \right| = \left| \sum_{j=1}^m \int_{r_{j-1}}^{r_j} (g(r, \xi_r^i) - g(r, \xi_{r_j}^i)) \, dr \right|.$$

It follows from Lemma 3.5 that $|g(r, \xi_r^i) - g(r, \xi_{r_j}^i)| \leq \mu(r)(r_j - r_{j-1})$ for $r \in [r_{j-1}, r_j]$, hence

$$|\mathcal{P}[q_i] - \mathcal{P}_m[q_i]|(t) \leq \delta \int_0^R \mu(r) \, dr \leq M\delta.$$

We rewrite (32) as

$$q_i(t) + b(t)\mathcal{P}_m[q_i](t) = w_i(t) + (\mathcal{P}_m[q_i] - \mathcal{P}[q_i])(t) \quad \forall t \in [0, T], \quad i = 1, 2. \tag{76}$$

From Proposition 3.6 it follows that

$$\|q_1 - q_2\| \leq \rho_m (\|w_1 - w_2\| + 2M\delta)$$

with $\rho_m = \prod_{j=1}^m (1 + \bar{b}\mu_j)$, $\mu_j = \int_{r_{j-1}}^{r_j} \mu(r) \, dr$. Note that $\log \rho_m = \sum_{j=1}^m \log(1 + \bar{b}\mu_j) \leq \sum_{j=1}^m \bar{b}\mu_j \leq \bar{b}M$. Hence,

$$\|q_1 - q_2\| \leq e^{\bar{b}M} (\|w_1 - w_2\| + 2M\delta).$$

We can choose δ arbitrarily small, and the assertion follows. □

We now prove the Lipschitz continuous dependence of the inverse mapping on b under suitable assumptions.

Proposition 3.7. *Let Hypotheses 3.3 (i), (ii) hold. Let $b_1, b_2 \in C[0, T]$*

be such that $0 \leq b_i(t) \leq \bar{b}$ for all $t \in [0, T]$, $i = 1, 2$, and let $w_1, w_2 \in C[0, T]$ be given. Let $q_1, q_2 \in C[0, T]$ be solutions of the equations

$$q_i(t) + b_i(t)\mathcal{P}[q_i](t) = w_i(t) \quad \forall t \in [0, T], \quad i = 1, 2. \tag{77}$$

Then we have

$$\|q_1 - q_2\| \leq e^{\bar{b}M} (\|w_1 - w_2\| + M_1 \|b_1 - b_2\|). \tag{78}$$

Proof. We have

$$q_1(t) + b_1(t)\mathcal{P}[q_1](t) = w_1(t),$$

$$q_2(t) + b_1(t)\mathcal{P}[q_2](t) = w_2(t) - (b_2(t) - b_1(t))\mathcal{P}[q_2](t), \tag{79}$$

and Theorem 3.4 together with (29) yields

$$\|q_1 - q_2\| \leq e^{\bar{b}M} (\|w_1 - w_2\| + \|\mathcal{P}[q_2]\| \|b_1 - b_2\|) \leq e^{\bar{b}M} (\|w_1 - w_2\| + M_1 \|b_1 - b_2\|),$$

which we wanted to prove. \square

4 Longitudinal oscillations of a piezoelectric beam

We keep Hypothesis 3.3 on the Preisach operator \mathcal{P} , and assume moreover that the mapping $(r, v) \mapsto v \frac{\partial g}{\partial v}(r, v)$ belongs to $L^1((0, \infty) \times \mathbb{R})$, and there exists $\mu_1 \in L^1(0, \infty)$ such that

$$|v| \frac{\partial g}{\partial v}(r, v) \leq \mu_1(r) \quad \text{a.e.}$$

This guarantees that the potential operator \mathcal{U} of the form (30) is bounded and Lipschitz continuous. The constitutive function $f(\varepsilon)$ in (12)–(14) will be assumed to possess the following properties.

Hypothesis 4.1. *The function $f : \mathbb{R} \rightarrow \mathbb{R}$ has Lipschitz continuous derivative f' and the functions $\varepsilon \mapsto 1/f(\varepsilon)$ and $\varepsilon \mapsto \varepsilon/f(\varepsilon)$ are Lipschitz continuous as well.*

Under these hypotheses, we reformulate the constitutive equation (17) in the form

$$\sigma = v\varepsilon_t + \underline{c}\varepsilon + \mathcal{W}[\varepsilon] \tag{80}$$

with a Lipschitz continuous operator $\mathcal{W} : C[0, T] \rightarrow C[0, T]$. Indeed, Eq. (16) can be rewritten as

$$q + \frac{1}{\kappa f(\varepsilon)} \mathcal{P}[q] = \frac{1}{\kappa f(\varepsilon)} (r - e\varepsilon), \quad q = \frac{E}{f(\varepsilon)}. \tag{81}$$

The functions $\varepsilon \mapsto 1/f(\varepsilon)$, $\varepsilon \mapsto \varepsilon/f(\varepsilon)$ are Lipschitz continuous by Hypothesis 4.1. Hence, by Proposition 3.7, the mapping $\varepsilon \mapsto q$ is Lipschitz continuous, and (17) can be written as

$$\sigma = v\varepsilon_t + \underline{c}\varepsilon - \frac{e}{\kappa} (r - e\varepsilon - \mathcal{P}[q]) + f'(\varepsilon) \mathcal{U}[q], \tag{82}$$

so that (80) holds with

$$\mathcal{W}[\varepsilon] = -\frac{e}{\kappa} (r - e\varepsilon - \mathcal{P}[q]) + f'(\varepsilon) \mathcal{U}[q].$$

This enables us to state the PDE problem (15) in the form

$$\rho u_{tt} - v u_{xxt} - \underline{c} u_{xx} = \mathcal{W}[u_x]_x \tag{83}$$

and we couple it with boundary conditions

$$u(0, t) = 0, (v u_{xt} + \underline{c} u_x + \mathcal{W}[u_x])(\ell, t) = s(t), \tag{84}$$

and initial conditions

$$u(x, 0) = u^0(x), u_t(x, 0) = u^1(x). \tag{85}$$

In variational form, the problem reads

$$\rho \int_0^\ell u_{tt} \phi \, dx + \int_0^\ell (\underline{c}u_x + \nu u_{xt} + \mathcal{W}[u_x]) \phi_x \, dx = s(t)\phi(\ell) \quad a.e. \forall \phi \in X, \quad (86)$$

where we set $X = \{\phi \in W^{1,2}(0, \ell) : \phi(0) = 0\}$.

Theorem 4.2. *Let $r \in C[0, T]$, $s \in L^2(0, T)$, $u^0 \in X$, and $u^1 \in L^2(0, \ell)$ be given, Then Problem (85)–(86) admits a unique solution $u \in C([0, T]; X)$ such that $u_{xt} \in L^2((0, \ell) \times (0, T))$, $u_t \in C([0, T]; L^2(\Omega))$, $u_{tt} \in L^2(0, T; X')$.*

Proof. For $v \in C([0, T]; X)$ such that $v_{xt} \in L^2((0, \ell) \times (0, T))$ and $v(x, 0) = u^0(x)$, we find u with the desired regularity as the solution of the linear problem

$$\rho \int_0^\ell u_{tt} \phi \, dx + \int_0^\ell (\underline{c}u_x + \nu u_{xt} + \mathcal{W}[v_x]) \phi_x \, dx = s(t)\phi(\ell) \quad a.e. \forall \phi \in X \quad (87)$$

with initial conditions (85). We now prove that the mapping $v \mapsto u$ is a contraction in the space

$$Y = \{v \in C([0, T]; X) : v_{xt} \in L^2((0, \ell) \times (0, T)), v(x, 0) = u^0(x), v(0, t) = 0\},$$

endowed with a suitable norm defined below in (91). Let $v, \hat{v} \in Y$ be given, and let u, \hat{u} be the corresponding solutions. We test the difference of Eqs. (87) for u and \hat{u}

$$\rho \int_0^\ell (u_{tt} - \hat{u}_{tt}) \phi \, dx + \int_0^\ell (\underline{c}(u_x - \hat{u}_x) + \nu(u_{xt} - \hat{u}_{xt}) + \mathcal{W}[v_x] - \mathcal{W}[\hat{v}_x]) \phi_x \, dx = 0 \quad a.e. \forall \phi \in X$$

by $\phi = u_t - \hat{u}_t$ and obtain, using the Lipschitz continuity of \mathcal{W} , that

$$\begin{aligned} & \frac{\rho}{2} \frac{d}{dt} \int_0^\ell (u_t - \hat{u}_t)^2 \, dx + \frac{\underline{c}}{2} \frac{d}{dt} \int_0^\ell (u_x - \hat{u}_x)^2 \, dx + \nu \int_0^\ell (u_{xt} - \hat{u}_{xt})^2 \, dx \\ &= - \int_0^\ell (\mathcal{W}[v_x] - \mathcal{W}[\hat{v}_x])(u_{xt} - \hat{u}_{xt}) \, dx \\ &\leq C \int_0^\ell |u_{xt} - \hat{u}_{xt}| \max_{\tau \in [0, t]} |v_x(x, \tau) - \hat{v}_x(x, \tau)| \, dx \end{aligned} \quad (88)$$

with some constant $C > 0$. We have $\max_{\tau \in [0, t]} |v_x(x, \tau) - \hat{v}_x(x, \tau)| \leq \int_0^t |v_{xt}(x, \tau) - \hat{v}_{xt}(x, \tau)| \, d\tau$, so that the right hand side (88) can be further estimated from above, using Hölder's and Young's inequalities, by the expression

$$\frac{\nu}{2} \int_0^\ell (u_{xt} - \hat{u}_{xt})^2 \, dx + \frac{C^2 t}{2\nu} \int_0^t \int_0^\ell (v_{xt} - \hat{v}_{xt})^2(x, \tau) \, dx \, d\tau.$$

We thus obtain from (88) that

$$\frac{1}{\nu} \frac{d}{dt} \int_0^\ell (\rho(u_t - \hat{u}_t)^2 + \underline{c}(u_x - \hat{u}_x)^2) \, dx + \int_0^\ell (u_{xt} - \hat{u}_{xt})^2 \, dx \leq \tilde{C} t \int_0^t \int_0^\ell (v_{xt} - \hat{v}_{xt})^2(x, \tau) \, dx \, d\tau, \quad (89)$$

with $\tilde{C} = \frac{C^2}{\nu^2}$. This is an inequality of the form

$$\dot{\alpha}(t) + \beta(t) \leq \tilde{C} t \int_0^t \delta(\tau) \, d\tau \quad (90)$$

with nonnegative functions

$$\alpha = \frac{1}{\nu} \int_0^\ell (\rho(u_t - \hat{u}_t)^2 + \underline{c}(u_x - \hat{u}_x)^2) \, dx, \quad \beta = \int_0^\ell (u_{xt} - \hat{u}_{xt})^2 \, dx, \quad \delta = \int_0^\ell (v_{xt} - \hat{v}_{xt})^2 \, dx.$$

We multiply (90) by $e^{-\tilde{C}t^2}$ and obtain

$$\frac{d}{dt} \left(e^{-\tilde{C}t^2} \left(\alpha(t) + \frac{1}{2} \int_0^t \delta(\tau) \, d\tau \right) \right) + 2\tilde{C}te^{-\tilde{C}t^2} \alpha(t) + e^{-\tilde{C}t^2} \beta(t) \leq \frac{1}{2} e^{-\tilde{C}t^2} \delta(t).$$

We now integrate the above inequality from 0 to T and obtain

$$\int_0^T e^{-\tilde{C}t^2} \beta(t) \, dt \leq \frac{1}{2} \int_0^T e^{-\tilde{C}t^2} \delta(t) \, dt.$$

We conclude that the mapping $v \mapsto u$ is a contraction in Y endowed with norm

$$\|v\| = \left(\int_0^T \int_0^\ell e^{-\tilde{c}t^2} |v_{xt}(x,t)|^2 dx dt + \int_0^\ell |v(x,0)|^2 dx \right)^{\frac{1}{2}}, \quad (91)$$

which implies existence and uniqueness of solutions. \square

5 Thin structures under uniaxial loading

To some extent, the results from Section 4 can be extended to a spatially three dimensional setting, as long as the setup still allows to assume that the third component of the dielectric displacement is constant in the polarization direction. Namely we first of all lift the model (12)–(14) to the spatially 3-d setting by prescribing a fixed polarization direction p , onto which we project the electric field. For simplicity we assume that p is constant. Now σ , ε , D , E are tensor and vector valued functions, respectively, and \underline{c} , e , κ are constant 4th, 3rd, and 2nd order material tensors (\underline{c} and κ are symmetric positive definite), but the internal variable q is a scalar valued function of the strain ε and of the projected electric field $p \cdot E$. Moreover, \mathcal{P} is still a scalar Preisach hysteresis operator with counterclockwise hysteresis potential \mathcal{U} , i.e., we assume (6) to hold. For some strictly positive scalar valued function $f : \mathbb{R}^6 \rightarrow \mathbb{R}^+$ satisfying Hypothesis 4.1 (with the preimage space \mathbb{R} replaced by \mathbb{R}^6), we consider the analog of (12)–(14)

$$\sigma = \underline{c}\varepsilon - eE + \mathcal{U}[q]Df(\varepsilon) \quad (92)$$

$$D = e^T \varepsilon + \kappa E + \mathcal{P}[q]p \quad (93)$$

$$F = \frac{1}{2}\varepsilon : (\underline{c}\varepsilon) + \frac{1}{2}E \cdot (\kappa E) + f(\varepsilon)\mathcal{U}[q] \quad (94)$$

$$q = \frac{p \cdot E}{f(\varepsilon)},$$

where $Df(\varepsilon)$ is the gradient of f , $Df : \mathbb{R}^6 \rightarrow \mathbb{R}^6$. It is readily checked that this ensures thermodynamic admissibility, i.e., the higher dimensional analog of (7).

We now consider Eqs. (1)–(2) in the form

$$\rho u_{tt} - \nabla_s^T \sigma = 0 \quad (95)$$

$$-\nabla \cdot D = 0 \quad (96)$$

$$\varepsilon = \nabla_s u \quad (97)$$

$$E = \nabla \phi, \quad (98)$$

where u is the mechanical displacement, ϕ the negative of the electric potential, and ∇_s the symmetric gradient.

Without loss of generality (and actually consistently with the usual notation) we assume that the polarization direction is parallel to the z axis, i.e.,

$$p = e_z,$$

and that the tensor of dielectric coefficients takes the form $\kappa = \begin{pmatrix} \kappa^{xy} & 0 \\ 0 & \kappa^z \end{pmatrix}$ with some positive definite 2×2 matrix κ^{xy} and $\kappa^z > 0$.

Our main restriction is the assumption that the z component of the dielectric displacement does not change in z direction

$$\frac{\partial}{\partial z} D^z = 0, \quad (99)$$

that the domain takes the cylindrical form

$$\Omega = \Omega^{xy} \times (0, \ell)$$

and that the boundary conditions

$$D \cdot n = 0 \text{ on } \{z\} \times \partial\Omega^{xy} \quad \forall z \in (0, \ell) \quad (100)$$

$$D \cdot n = -D^z = -r \text{ on } \{0\} \times \Omega^{xy}, \quad D \cdot n = D^z = r \text{ on } \{\ell\} \times \Omega^{xy} \quad (101)$$

hold, i.e., a current is prescribed, whose average over the boundary vanishes. (Note that therewith the electric potential can only be expected to be unique up to a constant, but the electric field will still be unique.) Assumption (99) is realistic, e.g., when dealing with a thin structure extended in the xy plane and excited by imposing some prescribed normal current via a pairs of opposite electrodes on top and bottom. Thus we have $D^z(x, y, z, t) = r(x, y, t)$ for all $(x, y, z) \in \Omega$, $t \in (0, T)$. Therewith, the combination of (93) with (96) can be split as follows:

$$\nabla^{xy} \cdot \left((e^T \varepsilon)^{xy} + \kappa^{xy} E^{xy} \right) = 0 \quad (102)$$

$$(e^T \varepsilon)^z + \kappa^z E^z + \mathcal{P}[q] = r \quad (103)$$

where the latter equation can be cast in a form convenient for application of the results from Section 3

$$q + \frac{1}{\kappa^z f(\varepsilon)} \mathcal{P}[q] = \frac{r - (e^T \varepsilon)^z}{\kappa^z f(\varepsilon)}$$

with $q = \frac{E^z}{f(\varepsilon)}$. Hence we can write

$$E^z(x, y, z, \cdot) = \Phi^z[\varepsilon(x, y, z, \cdot)] \quad \forall (x, y, z) \in \Omega$$

with a Lipschitz continuous mapping

$$\Phi^z : C[0, T] \rightarrow C[0, T].$$

On the other hand, for the xy part, testing (102) with ϕ (so that $\nabla^{xy} \phi = E^{xy}$) integrating by parts with respect to (x, y) , and using (100), we obtain the estimate

$$\|E^{xy}(\cdot, \cdot, z, t)\|_{L^2(\Omega^{xy})} \leq \frac{|e|}{\lambda_{\min}(\kappa^{xy})} \|\varepsilon(\cdot, \cdot, z, t)\|_{L^2(\Omega^{xy})} \quad \forall z \in (0, \ell), t \in (0, T),$$

where $\lambda_{\min}(\kappa^{xy})$ is the smallest eigenvalue of κ^{xy} . Thus, by integrating the square of both sides with respect to z

$$\|E^{xy}(\cdot, \cdot, \cdot, t)\|_{L^2(\Omega)} \leq \frac{|e|}{\lambda_{\min}(\kappa^{xy})} \|\varepsilon(\cdot, \cdot, \cdot, t)\|_{L^2(\Omega)} \quad \forall t \in (0, T).$$

By linearity, this provides us with Lipschitz continuity (with constant $L_{\Phi^{xy}} = \frac{|e|}{\lambda_{\min}(\kappa^{xy})}$) of the mapping

$$\Phi^{xy} : L^2(\Omega)^6 \rightarrow L^2(\Omega)^2$$

such that

$$E^{xy}(\cdot, \cdot, \cdot, t) = \Phi^{xy}(\varepsilon(\cdot, \cdot, \cdot, t)) \quad \forall t \in (0, T).$$

Analogously to the proof of Theorem 4.2, under Hypothesis 4.1, we will rewrite (92) in the form

$$\sigma = \nu \varepsilon_t + \underline{c} \varepsilon + \mathcal{W}(\varepsilon) \quad (104)$$

with the operator \mathcal{W} mapping the (tensor valued) function ε of (x, y, z, t) to a (tensor valued) function $\mathcal{W}(\varepsilon)$ of (x, y, z, t) as follows:

$$\begin{aligned} \mathcal{W}(\varepsilon)(x, y, z, t) = & -e \begin{pmatrix} \Phi^{xy}(\varepsilon(\cdot, \cdot, \cdot, t))(x, y, z) \\ \Phi^z[\varepsilon(x, y, z, \cdot)](t) \end{pmatrix} \\ & + \mathcal{U} \left[\frac{\Phi^z[\varepsilon(x, y, z, \cdot)]}{f(\varepsilon(x, y, z, \cdot))} \right] Df(\varepsilon(x, y, z, t)) \end{aligned}$$

To further proceed along the lines of the proof of Theorem 4.2, after elimination of the electric field we can now consider a purely mechanical problem

$$\rho u_{tt} - \nabla_s^T \left(\underline{c} \nabla_s u + \nu \nabla_s u_t + \mathcal{W}(\nabla_s u) \right) = 0 \quad \text{in } \Omega \times (0, T) \quad (105)$$

$$u(x, y, z, 0) = u^0, \quad u_t(x, y, z, 0) = u^1 \quad \text{for all } (x, y, z) \in \Omega \quad (106)$$

$$u = 0 \quad \text{on } \Gamma, \quad n \cdot \left(\nu \varepsilon_t + \mathcal{W}(\nabla_s u) \right) = s \quad \text{on } \partial\Omega \setminus \Gamma, \quad (107)$$

i.e., clamped at the boundary part Γ and traction free or loaded with some surface force on the remainder of the boundary.

The test space X becomes $X = \{\phi \in W^{1,2}(\Omega)^3 : \phi = 0 \text{ on } \Gamma\}$, the ansatz space

$$Y = \{v \in C([0, T]; X) : \nabla_s v_t \in L^2(\Omega \times (0, T))^6, \quad v(t = 0) = u^0, \quad v = 0 \text{ on } \Gamma\},$$

and the fixed point mapping \mathcal{T} maps $v \in Y$ to the unique solution $u \in Y$ of the linear variational problem

$$\rho \int_{\Omega} u_{tt} \phi \, d(x, y, z) + \int_{\Omega} (\underline{c} \nabla_s u + \nu \nabla_s u_t + \mathcal{W}(\nabla_s v)) : \nabla_s \phi \, d(x, y, z) = \int_{\Gamma} s \phi \, dS \quad a.e. \forall \phi \in X \quad (108)$$

with initial conditions (106), provided $s \in L^2(0, T; H^{-1/2}(\Gamma))$. Again, to show contractivity, for $v, \hat{v} \in Y$, we test the difference between the equations for $u = \mathcal{T}v$ and $\hat{u} = \mathcal{T}\hat{v}$ with $u_t - \hat{u}_t$ and use Young's inequality to end up with the estimate

$$\begin{aligned} \frac{\rho}{2} \frac{d}{dt} \|u_t - \hat{u}_t\|_{L^2(\Omega)}^2 + \frac{1}{2} \frac{d}{dt} \|\sqrt{c}(\nabla_s u - \nabla_s \hat{u})\|_{L^2(\Omega)}^2 + \frac{\nu}{2} \|\nabla_s u_t - \nabla_s \hat{u}_t\|_{L^2(\Omega)}^2 \\ \leq \frac{1}{2\nu} \|\mathcal{W}(\nabla_s v) - \mathcal{W}(\nabla_s \hat{v})\|_{L^2(\Omega)}^2. \end{aligned}$$

It remains to show that \mathcal{W} obeys the Lipschitz condition

$$\max_{\tau \in [0, t]} \|\mathcal{W}(\varepsilon) - \mathcal{W}(\hat{\varepsilon})\|_{L^2(\Omega)}^2 \leq C t \int_0^t \|\varepsilon_t - \hat{\varepsilon}_t\|_{L^2(\Omega)}^2 \, d\tau$$

for some constant $C > 0$ and any $\varepsilon, \hat{\varepsilon} \in L^2(\Omega \times (0, T))^6$ such that $\varepsilon_t, \hat{\varepsilon}_t \in L^2(\Omega \times (0, T))^6$ and $\varepsilon(t = 0) = \hat{\varepsilon}(t = 0)$. For the terms containing Φ^z and \mathcal{U} , this follows analogously to the 1-d case from the respective Lipschitz continuity. For the term containing Φ^{xy} we have

$$\begin{aligned} \max_{\tau \in [0, t]} \int_{\Omega} |\Phi^{xy}(\varepsilon(\cdot, \cdot, \cdot, \tau))(x, y, z) - \Phi^{xy}(\hat{\varepsilon}(\cdot, \cdot, \cdot, \tau))(x, y, z)|^2 \, d(x, y, z) \\ \leq L_{\Phi^{xy}}^2 \max_{\tau \in [0, t]} \|\varepsilon - \hat{\varepsilon}\|_{L^2(\Omega)}^2 = L_{\Phi^{xy}}^2 \max_{\tau \in [0, t]} \left\| \int_0^{\tau} (\varepsilon_t - \hat{\varepsilon}_t) \, d\tau_1 \right\|_{L^2(\Omega)}^2 \\ \leq L_{\Phi^{xy}}^2 t \int_0^t \|\varepsilon_t - \hat{\varepsilon}_t\|_{L^2(\Omega)}^2 \, d\tau. \end{aligned}$$

The rest of the proof goes exactly like in the 1-d case.

Therewith we have

Corollary 5.1. *Let $r \in C([0, T]; L^2(\Omega^{x,y}(0) \cup \Omega^{x,y}(\ell)))$, $s \in L^2(0, T; H^{-1/2}(\Gamma))$, $u^0 \in X$, and $u^1 \in L^2(\Omega)$ be given. Then Problem (92)–(98) with boundary conditions (99)–(101), (107) and initial conditions (106) admits a unique solution $(E^{xy}, E^z, u) \in C([0, T]; L^2(\Omega)) \times L^2(\Omega; C[0, T]) \times C([0, T]; X)$ such that $\nabla_s u_t \in L^2(\Omega \times (0, T))^6$, $u_t \in C([0, T]; L^2(\Omega)^3)$, $u_{tt} \in L^2(0, T; X')$, where $X = \{\phi \in W^{1,2}(\Omega)^3 : \phi = 0 \text{ on } \Gamma\}$.*

Proof. Starting from the existence, uniqueness, and regularity of u , which we have already obtained along the lines of the proof of Theorem 4.2, we use Lipschitz continuity of the mappings Φ^{xy} and Φ^z to obtain assertions on E^{xy} , E^z . In the latter case, the estimate for obtaining the claimed regularity looks as follows:

$$\begin{aligned} \int_{\Omega} \max_{t \in [0, T]} |E^z(x, y, z, t)|^2 \, d(x, y, z) &\leq \int_{\Omega} L_{\Phi^z} \max_{t \in [0, T]} |\nabla_s u(x, y, z, t)|^2 \, d(x, y, z) \\ &\leq \int_{\Omega} L_{\Phi^z} \left| u^0(x, y, z) + \int_0^T \nabla_s u_t(x, y, z, t) \, dt \right|^2 \, d(x, y, z) \\ &\leq 2L_{\Phi^z} \left(\|u^0\|_{L^2(\Omega)}^2 + T \|\nabla_s u_t\|_{L^2(\Omega \times (0, T))}^2 \right) \quad \square \end{aligned}$$

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Book Review

Gert Naue: Was ist Fluid-Turbulenz? Turbulente Oszillation und Konvektion. TWENTYSIX-Verlag, Norderstedt 2016, 420 S., 47 Tab., 210 Abb., Hardcover, EUR 58.99, ISBN 9783740710552, e-book EUR 5.49, ISBN 9783740736408.

Gegenstand des Buches ist ein neues Konzept zur Beschreibung der Turbulenz in Fluiden. Der Autor hat sich viele Jahrzehnte mit den Fragen der Turbulenz auseinandergesetzt. Zwar anerkennend, dass auf der Basis der auf Prandtl zurückgehenden Ansätze und deren Weiterentwicklung praktische Aufgaben, die mit turbulenten Strömungen zusammenhängen, meist ausreichend gut beschreibbar sind und damit einer Berechnung zugeführt werden können, wurde der theoretische Hintergrund kritisch hinterfragt. Nach gründlichem Durchdenken wurde die etablierte Herangehensweise jedoch als nicht befriedigend angesehen, wenn es um ein tieferes Verständnis der Turbulenz geht. Deshalb reifte im Laufe vieler Jahre ein neuer Ansatz, der in diesem umfangreichen Buch detailliert vorgestellt wird, um damit eine neue Turbulenztheorie zu begründen, die Anwendbarkeit an Beispielen zu zeigen und zur Weiterentwicklung der Theorie anzuregen. Die im Buch enthaltenen umfangreichen Berechnungen basieren auf dem Softwarepaket MATHEMATICA von Wolfram.

Das Buch enthält 10 Kapitel und einen Anhang. Im Kapitel 1 wird auf Charakteristika laminarer und turbulenter Strömungen eingegangen. Anstelle der Beschreibung mit einer Turbulenzviskosität soll basierend auf der Selbstähnlichkeit und der optimalen Selbstorganisation ein neues Konzept entwickelt werden, dessen Grundzüge schon hier grob skizziert werden. Gleichzeitig erfolgt der Hinweis, dass eine spezielle Mathematik dafür notwendig ist. Bevor darauf jedoch näher eingegangen wird, erfolgt im Kapitel 2 ein kurzer Überblick über die konventionellen Modelle der Turbulenzmodellierung und die in einem früheren Buch des Autors entwickelte Theorie der Trajektionsprozesse, die einen alternativen Zugang und eine andere Interpretation der bei Turbulenzmodellen notwendigen Zusatzterme ermöglicht. Auch im Kapitel 3 stehen noch allgemeine Grundlagen im Vordergrund. Es werden unterschiedliche Geschwindigkeitsbegriffe eingeführt und dabei der Bogen zur Wellenausbreitung gespannt, da im Verlauf des Buches eine Formulierung des Impulssatzes auf der Basis von den Turbulenzstrukturen zugeordneten Welleneigenschaften eingeführt wird. Diese neue Transporteigenschaft wird Tur-

bulenzgeschwindigkeit genannt, die zusätzlich zu der an die Masse gebundenen Konvektionsgeschwindigkeit auftritt. In diesem Zusammenhang werden den Turbulenzstrukturen auch u. a. Dipol- und/oder Quadrupol-Strukturen zugeschrieben. Weiterhin wird herausgearbeitet, dass durch die Koexistenz von Konvektionsgeschwindigkeit und einer an Wellenpakete gebundenen Gruppengeschwindigkeit Phasenverschränkungen auftreten, die für die Ausbildung der Strömungseigenschaften verantwortlich sind. Die dafür notwendigen mathematischen Ansätze werden zunächst in allgemeiner Form aufgeführt.

Ausführlich wird im Kapitel 4 auf die optimale Selbstorganisation der Turbulenz eingegangen. Dabei werden konkurrierende und sequentielle Vorgänge unterschieden, für deren Beschreibung die Trajektionsdifferentialgleichung und die Informationsentropie herangezogen werden. Zusätzlich ist für die Charakterisierung der wechselwirkenden Wellenpakete deren Brechungsverhalten einzubeziehen. Dies führt, wie schon vor einigen Jahren vom Autor herausgearbeitet wurde, zu Gleichungen für die Turbulenzgeschwindigkeit sowie zur Vorhersage kritischer Kennzahlen von Strömungen und des Druckverlustes turbulenter Rohrströmungen für alle Rauigkeitsbereiche. Da dies für die folgenden Betrachtungen wichtig ist, wird in kompakter Form noch einmal darauf eingegangen. Außerdem werden Aussagen zu Freistrahlen getroffen. Es wird gezeigt, dass auch hierfür mittels Brechungsgesetz und Entropiebetrachtungen, wie schon bei den Rohrströmungen, die experimentellen Befunde für den Ausbreitungskegel, das Geschwindigkeitsprofil und die Turbulenzparameter ohne die Einbeziehung von empirischen Konstanten wiedergegeben werden können. Darauf aufbauend erfolgt die Ausdehnung der Betrachtungen auf das Verhalten von Freistrahlen im Querwind, wobei einerseits die bisher verwendeten empirischen Parameter mit sehr guter Genauigkeit vorausgesagt werden können und andererseits die Ursachen für das Verhalten theoretisch erklärt werden können. Den Abschluss dieses Abschnitts bilden Ausführungen zu Turbulenzstrukturen in Wellenpaketen, die das Ergebnis von Wellenpaaren mit speziellen Eigenschaften sind. Die Wellenpakete können zur Ausbildung von Singularitäten mit einer Fluidgrenze führen, in deren Inneren Brechungsvorgänge auftreten. Die Strukturen der Singularitäten und ihre Eigenschaften werden ausführlich diskutiert bis hin zu den Dissipationsmechanismen im

Inneren von Turbulenzstrukturen, wobei aber derzeit noch Fragen offen bleiben.

Kapitel 5 ist dem Oszillationsmodell der Turbulenz gewidmet, wodurch schon vorangegangene Aussagen zu den Welleneigenschaften der Strömung tiefgründiger betrachtet werden. Es geht zunächst um die Bestimmung der ungeborenen kontravarianten Turbulenzgeschwindigkeit und deren Zuordnung zu den Raumkoordinaten, wozu zylindrische, sphärische und in späteren Kapiteln auch hyperbolische Polarkoordinaten herangezogen werden. Für die Turbulenz werden die Oszillations-, Amplituden-, Polar-, Phasen-, Koexistenz-, Leistungs- und Energie- sowie die Faltungsform hergeleitet, diskutiert und in Verbindung mit einigen speziellen Strömungsvorgängen gebracht. Zum Abschluss wird die geometrische Verschränkungsform am Beispiel des Freistrahls diskutiert. Im Kapitel 6 werden die Prinzipien der Turbulenz (Fraktal-, Spiegelungs-, Verschränkungs-, Brechungs- und Faltungsprinzip) postuliert. Die konkreten Erscheinungsformen hängen zwar von den unterschiedlichen Rand- und Nebenbedingungen ab, die Prinzipien werden aber als allgemeingültige Wesensmerkmale der Turbulenz angesehen. Sie beschreiben z. B. auch Effekte wie die lokale Koexistenz von turbulenter und laminarer Strömung bzw. von Intermittenz und ermöglichen die Rückführbarkeit turbulenter Prozesse auf Basiszahlen.

In den Kapiteln 7 und 8, die jeweils ca. 100 Seiten umfassen, erfolgt die Anwendung der vorgestellten Theorie auf erzwungen- bzw. frei-turbulente Strömungen und der Vergleich der theoretischen Ergebnisse mit experimentellen Befunden. Es werden einerseits Strömungen in glatten/rauen Rohren, mit Drall, Querschnittsänderungen, Intermittenz, Einlaufeffekten, die Zirkulation in Ringspalten und Rührbehältern sowie andererseits Trennschichten, Nachläufe und unterschiedliche Freistrahllkonstellationen ausführlich behandelt. Außerdem werden das Auftreten von Solitonen in Kanälen, Strömungsablösungen sowie der Zusammenhang zwischen Dissipation und Turbulenzstruktur betrachtet. Dabei werden nur die zuvor abgeleiteten theoretisch begründeten Parameter herangezogen und damit Aussagen zu den Geschwindigkeitsprofilen, Turbu-

lenzcharakteristika, Druckverlusten u. a. m. gemacht, die durchweg in sehr guter Übereinstimmung mit den experimentellen Ergebnissen sind und zu ihrem besseren Verständnis beitragen.

Kapitel 9 beinhaltet eine detaillierte Diskussion des bisher als universell angesehen Wandgesetzes bei erzwungen-turbulenten Strömungen. Wieder werden die experimentellen Befunde analysiert, wobei jedoch auch Messwerte einbezogen werden, die bei der Aufstellung des Wandgesetzes unberücksichtigt blieben. Es zeigt sich auch hier, dass sie mit dem Modellansatz sehr gut beschreibbar sind und auch die früher aussortierten Werte eingebunden werden können.

Obwohl in den vorangegangenen Abschnitten schon zahlreiche unterschiedliche Strömungsprobleme behandelt wurden, bleiben noch viele Probleme der Turbulenz offen. Einige davon werden im Kapitel 10 angesprochen. Der Autor will so zum Weiterdenken anregen und verbindet damit die Hoffnung, dass andere Wissenschaftler auf der Basis der neuen Theorie diese Lücken schließen. Im Anhang sind MATHEMATICA-Befehle zur Koordinatentransformation und zur Rechnung mit ko- bzw. kontravarianten Vektorfeldern aufgelistet. Außerdem sind dort weitere Daten der Profile in glatten und rauen Rohren für unterschiedliche Reynolds-Zahlen hinterlegt.

Die neue Herangehensweise an die Turbulenz, für die auch andere physikalische Gesetze und mathematische Methoden notwendig sind, erfordert die Bereitschaft, sich darauf einzulassen. Die behandelten Beispiele zeigen jedoch, dass damit auch neue Einblicke in die Turbulenzprozesse möglich sind und empirische durch theoretische Befunde ersetzt werden können. Das Buch kann allen, die an den Fragen der Strömungsmechanik und der Modellierung der dabei auftretenden Erscheinungen interessiert sind, uneingeschränkt empfohlen werden. Mit dem bei diesem Umfang des Buches relativ niedrigen Preis will der Autor einer angestrebten weiten Verbreitung keine Hindernisse in den Weg legen.

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