

Tensor methods for higher-dimensional chemical Fokker-Planck equation

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Joint work with Radek Erban and Shuohao Liao

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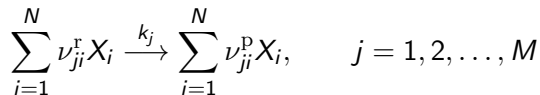


1. Introduction
2. Well-posedness of Fokker-Planck equation
3. Tensor methods
4. Error estimation

General chemical system



Well mixed reactor, N chemical species, M reactions:



Markov proces:

Reaction j occurs in time interval $(t, t + dt)$ with probability $\alpha_j(\mathbf{X}(t)) dt$, where

- ▶ $\mathbf{X}(t) = [X_1(t), \dots, X_N(t)]$
- ▶ $X_i(t)$ are numbers of molecules at time t , $i = 1, \dots, N$
- ▶ $\alpha_j(\mathbf{x}) = k_j \prod_{i=1}^N \binom{x_i}{\nu_{ji}^r}$ are propensities, $j = 1, \dots, M$
- ▶ $\nu_{ji} = \nu_{ji}^p - \nu_{ji}^r$, change of X_i during reaction j



Quantity of interest:

$$\rho(\mathbf{x}, t) = \Pr(\mathbf{X}(t) = \mathbf{x})$$

Exact description

- ▶ Stochastic simulations (Gillespie algorithm)
- ▶ Chemical master equation

Approximate description

- ▶ Stochastic simulations (Langevin equation)
- ▶ Chemical Fokker–Planck equation

Chemical Fokker–Planck equation



$$\frac{\partial p}{\partial t}(\mathbf{x}, t) = \sum_{i=1}^N \sum_{k=1}^N \frac{\partial^2}{\partial x_i \partial x_k} \left[\left(\frac{1}{2} \sum_{j=1}^M \nu_{ji} \nu_{jk} \alpha_j(\mathbf{x}) \right) p(\mathbf{x}, t) \right] - \sum_{i=1}^N \frac{\partial}{\partial x_i} \left[\left(\sum_{j=1}^M \nu_{ji} \alpha_j(\mathbf{x}) \right) p(\mathbf{x}, t) \right]$$

Vector notation

$$\frac{\partial p}{\partial t}(\mathbf{x}, t) = \operatorname{div} [\operatorname{div}(\mathcal{A}(\mathbf{x})p(\mathbf{x}, t)) - \mathbf{b}(\mathbf{x})p(\mathbf{x}, t)]$$

► $\mathcal{A} \in \mathbb{R}^{N \times N}$, $\mathcal{A}_{ik}(\mathbf{x}) = \frac{1}{2} \sum_{j=1}^M \nu_{ji} \nu_{jk} \alpha_j(\mathbf{x})$

► $\mathbf{b} \in \mathbb{R}^N$, $b_i(\mathbf{x}) = \sum_{j=1}^M \nu_{ji} \alpha_j(\mathbf{x})$



2. Well-posedness of chemical Fokker-Planck equation



$$\frac{\partial p}{\partial t} = \text{div} [\text{div}(\mathcal{A}p) - \mathbf{b}p]$$

Fokker-Planck operator is elliptic if $\mathcal{A}(\mathbf{x})$ is symmetric positive definite.

Notice: $\mathcal{A}_{ik}(\mathbf{x}) = \frac{1}{2} \sum_{j=1}^M \nu_{ji} \alpha_j(\mathbf{x}) \nu_{jk} \implies \mathcal{A} = \frac{1}{2} \boldsymbol{\nu}^T \boldsymbol{\alpha} \boldsymbol{\nu}$

- ▶ $\boldsymbol{\nu} \in \mathbb{R}^{M \times N}$, ν_{ji} are stoichiometric coefficients
- ▶ $\boldsymbol{\alpha} \in \mathbb{R}^{M \times M}$, $\boldsymbol{\alpha} = \text{diag}(\alpha_1, \dots, \alpha_M)$

Lemma.

Let N columns of $\boldsymbol{\nu} \in \mathbb{R}^{M \times N}$ be linearly independent.

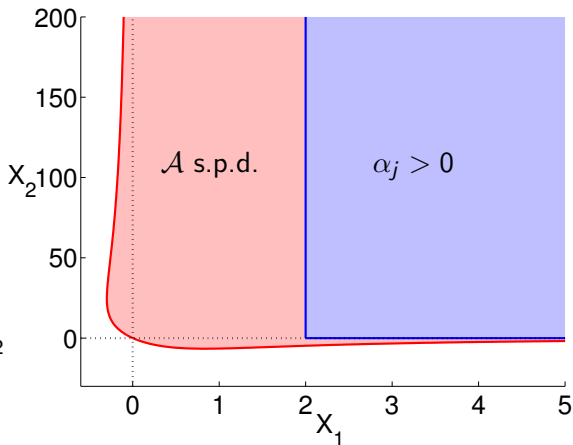
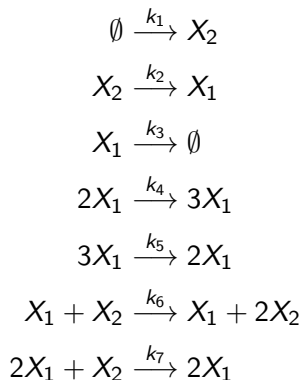
Let $\alpha_j > 0$ for all $j = 1, 2, \dots, M$.

Then $\mathcal{A} = \frac{1}{2} \boldsymbol{\nu}^T \boldsymbol{\alpha} \boldsymbol{\nu}$ is symmetric positive definite.

Proof. Let $\mathbf{y} \in \mathbb{R}^N$, $\mathbf{y} \neq 0$. Then $\mathbf{z} = \boldsymbol{\nu} \mathbf{y} \neq 0$. Thus,

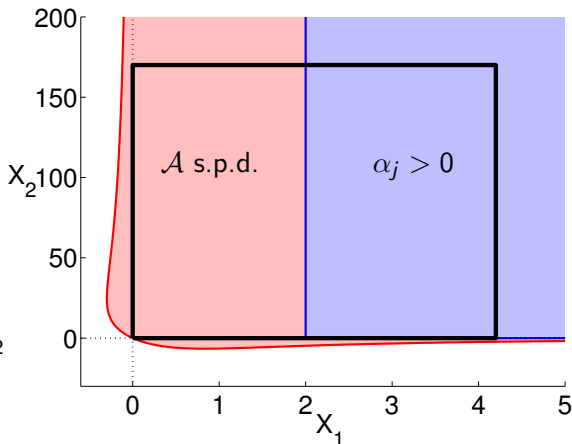
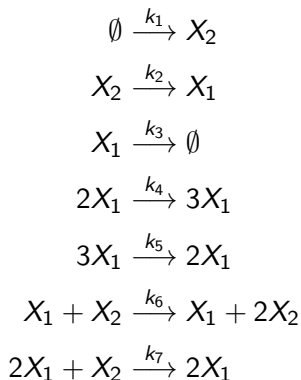
$$\mathbf{y}^T \mathcal{A} \mathbf{y} = \frac{1}{2} \mathbf{y}^T \boldsymbol{\nu}^T \boldsymbol{\alpha} \boldsymbol{\nu} \mathbf{y} = \frac{1}{2} \mathbf{z}^T \boldsymbol{\alpha} \mathbf{z} > 0.$$

Example: SNIPER



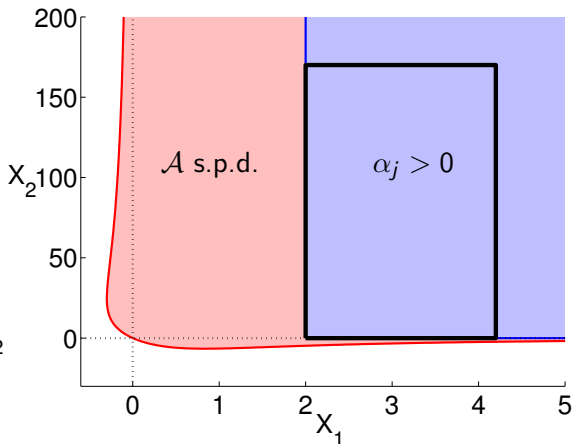
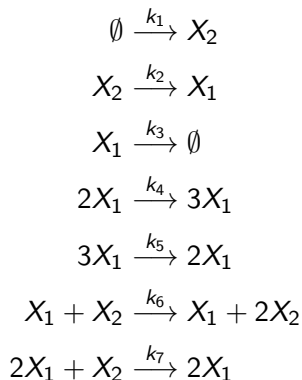
[R. Erban, S.J. Chapman, I.G. Kevrekidis, T. Vejchodsky, SIAM J. Appl. Math. 2009.]

Example: SNIPER



[R. Erban, S.J. Chapman, I.G. Kevrekidis, T. Vejchodsky, SIAM J. Appl. Math. 2009.]

Example: SNIPER



[R. Erban, S.J. Chapman, I.G. Kevrekidis, T. Vejchodsky, SIAM J. Appl. Math. 2009.]

Time-dependent chemical Fokker-Planck equation



$$\begin{aligned}\frac{\partial p}{\partial t}(\mathbf{x}, t) &= \operatorname{div} [\operatorname{div}(\mathcal{A}(\mathbf{x})p(\mathbf{x}, t)) - \mathbf{b}(\mathbf{x})p(\mathbf{x}, t)] \quad \text{for } \mathbf{x} \in \Omega, \quad t > 0 \\ 0 &= [\operatorname{div}(\mathcal{A}(\mathbf{x})p(\mathbf{x}, t)) - \mathbf{b}(\mathbf{x})p(\mathbf{x}, t)] \cdot \mathbf{n} \quad \text{for } \mathbf{x} \in \partial\Omega, \quad t > 0 \\ p(\mathbf{x}, 0) &= p_0 \quad \text{for } \mathbf{x} \in \Omega\end{aligned}$$

- ▶ Choice of Ω ?
- ▶ No flux boundary conditions \Rightarrow conservativity

Stationary chemical Fokker-Planck equation



$$\begin{aligned}\operatorname{div} [\operatorname{div}(\mathcal{A}(\mathbf{x})\rho(\mathbf{x})) - \mathbf{b}(\mathbf{x})\rho(\mathbf{x})] &= 0 && \text{for } \mathbf{x} \in \Omega \\ [\operatorname{div}(\mathcal{A}(\mathbf{x})\rho(\mathbf{x})) - \mathbf{b}(\mathbf{x})\rho(\mathbf{x})] \cdot \mathbf{n} &= 0 && \text{for } \mathbf{x} \in \partial\Omega \\ \int_{\Omega} \rho(\mathbf{x}) \, d\mathbf{x} &= 1\end{aligned}$$

- ▶ No flux boundary conditions \Rightarrow 'unique' nontrivial solution based on adjoint problem and Fredholm alternative
- ▶ Eigenvalue problem

Zero Dirichlet boundary conditions



Equation:

$$\begin{aligned}\operatorname{div}[\operatorname{div}(\mathcal{A}(\mathbf{x})p(\mathbf{x})) - \mathbf{b}(\mathbf{x})p(\mathbf{x})] &= 0 && \text{for } \mathbf{x} \in \Omega \\ p(\mathbf{x}) &= 0 && \text{for } \mathbf{x} \in \partial\Omega\end{aligned}$$

$$\implies p(\mathbf{x}) = 0$$

Eigenvalue problem:

$$\begin{aligned}\operatorname{div}[\operatorname{div}(\mathcal{A}(\mathbf{x})p(\mathbf{x})) - \mathbf{b}(\mathbf{x})p(\mathbf{x})] &= \lambda_1 p(\mathbf{x}) && \text{for } \mathbf{x} \in \Omega \\ p(\mathbf{x}) &= 0 && \text{for } \mathbf{x} \in \partial\Omega\end{aligned}$$

$$\int_{\Omega} p(\mathbf{x}) \, d\mathbf{x} = 1$$



Eigenvalue problem:

$$\begin{aligned} \operatorname{div} [\operatorname{div}(\mathcal{A}(\mathbf{x})p(\mathbf{x})) - \mathbf{b}(\mathbf{x})p(\mathbf{x})] &= \lambda_1 p(\mathbf{x}) && \text{for } \mathbf{x} \in \Omega \\ p(\mathbf{x}) &= 0 && \text{for } \mathbf{x} \in \partial\Omega \end{aligned}$$

$$\int_{\Omega} p(\mathbf{x}) \, d\mathbf{x} = 1$$

- ▶ Cover Ω by a mesh
- ▶ Replace derivatives by finite differences
- ▶ $\implies \mathbf{A}\mathbf{p} = \tilde{\lambda}_1\mathbf{p}$

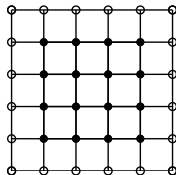
Curse of dimensionality



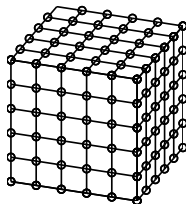
▶ $N = 1 \Rightarrow n$ nodes



▶ $N = 2 \Rightarrow n^2$ nodes



▶ $N = 3 \Rightarrow n^3$ nodes



▶ In general n^N nodes.



3. Tensor methods

Tensor product of matrices



$$A = \begin{bmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{bmatrix} \quad B = \begin{bmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{bmatrix}$$

$$A \otimes B = \begin{bmatrix} a_{11} \begin{bmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{bmatrix} & a_{12} \begin{bmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{bmatrix} \\ a_{21} \begin{bmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{bmatrix} & a_{22} \begin{bmatrix} b_{11} & b_{12} \\ b_{21} & b_{22} \end{bmatrix} \end{bmatrix} = \begin{bmatrix} a_{11}b_{11} & a_{11}b_{12} & a_{12}b_{11} & a_{12}b_{12} \\ a_{11}b_{21} & a_{11}b_{22} & a_{12}b_{21} & a_{12}b_{22} \\ a_{21}b_{11} & a_{21}b_{12} & a_{22}b_{11} & a_{22}b_{12} \\ a_{21}b_{21} & a_{21}b_{22} & a_{22}b_{21} & a_{22}b_{22} \end{bmatrix}$$

Tensor product of vectors



$$\mathbf{a} = \begin{bmatrix} a_1 \\ a_2 \end{bmatrix} \quad \mathbf{b} = \begin{bmatrix} b_1 \\ b_2 \\ b_3 \end{bmatrix}$$

$$\mathbf{a} \otimes \mathbf{b} = \begin{bmatrix} a_1 \begin{bmatrix} b_1 \\ b_2 \\ b_3 \end{bmatrix} \\ a_2 \begin{bmatrix} b_1 \\ b_2 \\ b_3 \end{bmatrix} \end{bmatrix} = \begin{bmatrix} a_1 b_1 \\ a_1 b_2 \\ a_1 b_3 \\ a_2 b_1 \\ a_2 b_2 \\ a_2 b_3 \end{bmatrix}$$

$$(\mathbf{a} \otimes \mathbf{b} \otimes \mathbf{c})_{[ijk]} = a_i b_j c_k$$

Tensor product of vectors



Notice:

$$\mathbf{a}_1 \in \mathbb{R}^n, \mathbf{a}_2 \in \mathbb{R}^n, \dots, \mathbf{a}_N \in \mathbb{R}^n \implies \mathbf{a}_1 \otimes \mathbf{a}_2 \otimes \dots \otimes \mathbf{a}_N \in \mathbb{R}^{n^N}$$

Tensor methods for N dimensional Laplacian



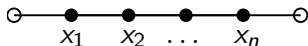
$$\begin{aligned} -\Delta u(\mathbf{x}) &= f(\mathbf{x}) & \text{for } \mathbf{x} \in \Omega &= (0, 1)^N \\ u(\mathbf{x}) &= 0 & \text{for } \mathbf{x} \in \partial\Omega \end{aligned}$$



$$-u''(x) = f(x) \quad \text{for } x \in \Omega = (0, 1)$$

$$u(0) = u(1) = 0$$

- ▶ x_i form a partition of $(0, 1)$



- ▶ $i = 1, 2, \dots, n$

- ▶ $\mathbf{f} \in \mathbb{R}^n$, $\mathbf{f}_i = f(x_i)$

- ▶ $\mathbf{u} \in \mathbb{R}^n$, $u(x_i) \approx u_i$

- ▶ $-u''(x_i) \approx \frac{-u_{i-1} + 2u_i - u_{i+1}}{h^2}$

- ▶ $-u''(x) \approx D\mathbf{u}$

- ▶ $D\mathbf{u} = \mathbf{f}$

$$D = \frac{1}{h^2} \begin{bmatrix} 2 & -1 & 0 & \cdots & 0 \\ -1 & 2 & -1 & \cdots & 0 \\ 0 & -1 & 2 & \cdots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & 2 \end{bmatrix}$$



$$-\frac{\partial^2 u}{\partial x_1^2}(x_1, x_2) - \frac{\partial^2 u}{\partial x_2^2}(x_1, x_2) = f(x_1, x_2) \quad \text{for } (x_1, x_2) \in \Omega = (0, 1)^2$$

$$u(x_1, x_2) = 0 \quad \text{for } \mathbf{x} \in \partial\Omega$$

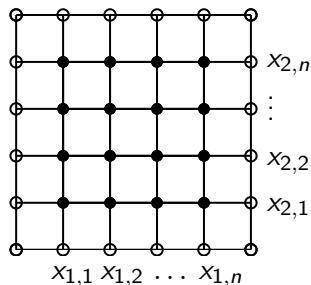
If $u(x_1, x_2) = q_1(x_1)q_2(x_2)$ then

$$-\Delta u(x_1, x_2) = -q_1''(x_1)q_2(x_2) - q_1(x_1)q_2''(x_2)$$

$$-\Delta u \approx (D\mathbf{q}_1) \otimes \mathbf{q}_2 + \mathbf{q}_1 \otimes (D\mathbf{q}_2)$$

$$= \underbrace{(D \otimes I + I \otimes D)}_A \underbrace{(\mathbf{q}_1 \otimes \mathbf{q}_2)}_{\mathbf{u}}$$

$$\Rightarrow A\mathbf{u} = \mathbf{f}$$



Laplacian in 3D

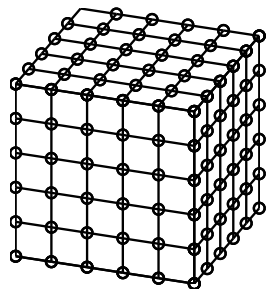


$$\blacktriangleright A = D \otimes I \otimes I + I \otimes D \otimes I + I \otimes I \otimes D$$

$$\blacktriangleright \mathbf{A}\mathbf{u} = \mathbf{f}$$

$$u(x_1, x_2, x_3) \approx \sum_{r=1}^R q_1^r(x_1) q_2^r(x_2) q_3^r(x_3)$$

$$\mathbf{u} \approx \sum_{r=1}^R \mathbf{q}_1^r \otimes \mathbf{q}_2^r \otimes \mathbf{q}_3^r$$



$$\begin{aligned} \mathbf{A}\mathbf{u} &= (D \otimes I \otimes I + I \otimes D \otimes I + I \otimes I \otimes D) \left(\sum_{r=1}^R \mathbf{q}_1^r \otimes \mathbf{q}_2^r \otimes \mathbf{q}_3^r \right) \\ &= \sum_{r=1}^R (D\mathbf{q}_1^r) \otimes \mathbf{q}_2^r \otimes \mathbf{q}_3^r + \mathbf{q}_1^r \otimes (D\mathbf{q}_2^r) \otimes \mathbf{q}_3^r + \mathbf{q}_1^r \otimes \mathbf{q}_2^r \otimes (D\mathbf{q}_3^r) \end{aligned}$$



Tensor truncation:

$$\sum_{r=1}^R \mathbf{q}_1^r \otimes \mathbf{q}_2^r \otimes \cdots \otimes \mathbf{q}_N^r \approx \sum_{s=1}^S \mathbf{v}_1^s \otimes \mathbf{v}_2^s \otimes \cdots \otimes \mathbf{v}_N^s \quad \text{with } S < R$$

Canonical decomposition:

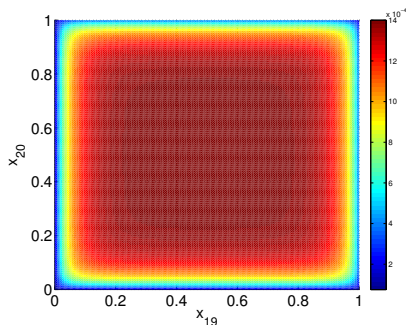
$$\mathbf{u} \approx \underbrace{\sum_{r=1}^R \mathbf{q}_1^r \otimes \mathbf{q}_2^r \otimes \cdots \otimes \mathbf{q}_N^r}_{nNR \text{ entries}}, \quad \text{where } \mathbf{q}_d^r \in \mathbb{R}^n, d = 1, 2, \dots, N$$

Numerical stability \Rightarrow Tensor train format

Example: 20D Laplacian



$$\begin{aligned} -\Delta u(\mathbf{x}) &= 1 & \text{for } \mathbf{x} \in \Omega = (0, 1)^{20} \\ u(\mathbf{x}) &= 0 & \text{for } \mathbf{x} \in \partial\Omega \end{aligned}$$



$$\text{Plot of } \tilde{u}(x_{19}, x_{20}) = \int_0^1 \cdots \int_0^1 u(x_1, \dots, x_{20}) dx_1 \cdots dx_{18}$$

Chemical Fokker-Planck equation

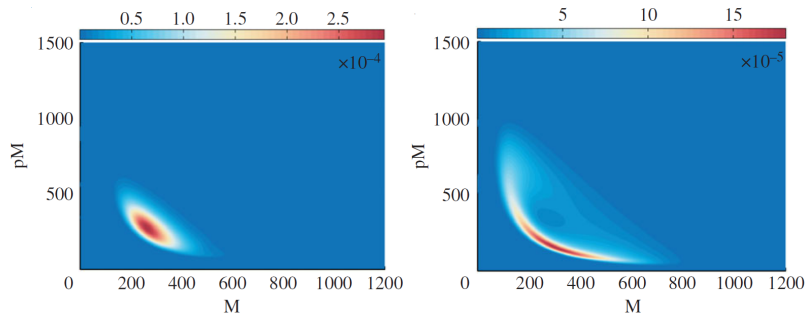


$$\begin{aligned} \operatorname{div} [\operatorname{div}(\mathcal{A}(\mathbf{x})\rho(\mathbf{x})) - \mathbf{b}(\mathbf{x})\rho(\mathbf{x})] &= \lambda_1 \rho(\mathbf{x}) & \text{for } \mathbf{x} \in \Omega \\ \rho(\mathbf{x}) &= 0 & \text{for } \mathbf{x} \in \partial\Omega \end{aligned}$$

$$\int_{\Omega} \rho(\mathbf{x}) \, d\mathbf{x} = 1$$

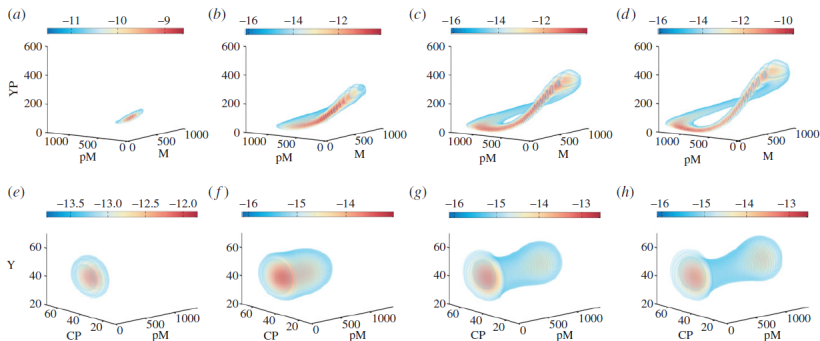
- ▶ $\mathcal{A}_{ik}(\mathbf{x}) = \frac{1}{2} \sum_{j=1}^M \nu_{ji} \nu_{jk} k_j \prod_{i=1}^N \binom{x_i}{\nu_{ji}^r}$
- ▶ $b_i(\mathbf{x}) = \sum_{j=1}^M \nu_{ji} k_j \prod_{i=1}^N \binom{x_i}{\nu_{ji}^r} \implies \mathbf{A}\mathbf{p} = \tilde{\lambda}_1 \mathbf{p}$
- ▶ TT-toolbox by Ivan Oseledets et al.
- ▶ Unknown parameter as an additional dimension
- ▶ Multilevel approach

Bifurcation analysis of the stochastic cell cycle model



- ▶ 6 chemical species + 1 parameter dimension
- ▶ [S. Liao, T. Vejchodsky, R. Erban:
J. R. Soc. Interface 12 (108), 2015, 20150233]

Bifurcation analysis of the stochastic cell cycle model



► 6 chemical species + 1 parameter dimension

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4. Error estimation

Type of errors



- ▶ Error of Fokker-Planck equation
- ▶ Error caused by the domain truncation
- ▶ Error of artificial boundary conditions
- ▶ Discretization error
- ▶ Tensor error
- ▶ Algebraic error

Future work:

- ▶ estimates of individual errors
- ▶ equilibrate individual errors



- ▶ Tensor methods can break the curse of dimensionality.
- ▶ Parametric chemical Fokker-Planck equation
 - ▶ Bifurcation analysis
 - ▶ Parameter estimation
 - ▶ Sensitivity analysis
- ▶ Tensor methods can be combined with the FEM, DG, etc.
- ▶ Tensor methods can solve the chemical master equation.

Thank you for your attention

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