

# On singular limits for inviscid fluid flows

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# Singular limits

## Primitive system

$$\partial_t \mathbf{U}_\varepsilon + \mathcal{B}_\varepsilon \mathbf{U}_\varepsilon + \mathcal{A}(\mathbf{U}_\varepsilon) = 0, \quad \mathbf{U}_\varepsilon(0) = \mathbf{U}_{0,\varepsilon}$$

## Decomposition

$$\mathbf{U}_\varepsilon = \mathbf{V}_\varepsilon + \underbrace{\mathbf{W}_\varepsilon}_{\text{oscillatory}}, \quad \mathcal{B}_\varepsilon \mathbf{V}_\varepsilon = 0,$$

## Oscillatory part

$$\partial_t \mathbf{W}_\varepsilon + \mathcal{B}_\varepsilon \mathbf{W}_\varepsilon = 0$$

$$\partial_t \mathbf{V}_\varepsilon + \mathcal{A}(\mathbf{V}_\varepsilon + \mathbf{W}_\varepsilon) = 0$$

## Target system

$$\partial_t \mathbf{U} + \mathcal{A}(\mathbf{U}) = 0, \quad \mathbf{U}(0, \cdot) = \mathbf{U}_0$$

## Ill vs. well prepared initial data

$$\mathbf{U}_{0,\varepsilon} = \mathbf{V}_{0,\varepsilon} + \mathbf{W}_{0,\varepsilon}, \quad \mathbf{V}_{0,\varepsilon} \rightarrow \mathbf{U}, \quad \mathbf{W}_0 \begin{cases} \rightarrow 0 \\ \rightarrow \mathbf{W}_{0,\varepsilon} \end{cases}$$

# Primitive system

## Euler system - standard variables

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0,$$

$$\partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \nabla_x p(\varrho, \vartheta) = \varrho \nabla_x F$$

$$\begin{aligned} \partial_t \left( \frac{1}{2} \varrho |\mathbf{u}|^2 + \varrho e(\varrho, \vartheta) \right) + \operatorname{div}_x \left[ \left( \frac{1}{2} \varrho |\mathbf{u}|^2 + \varrho e(\varrho, \vartheta) \right) \mathbf{u} \right] + \operatorname{div}_x(p(\varrho, \vartheta) \mathbf{u}) \\ = \varrho \nabla_x F \cdot \mathbf{u} \end{aligned}$$

# Conservative variables - scaling

## Polytropic EOS

$$\mathbf{m} = \rho \mathbf{u}, \quad E_\varepsilon = \frac{1}{2} \frac{|\mathbf{m}|^2}{\rho} + \frac{1}{\varepsilon^2} \rho e, \quad p = (\gamma - 1) \rho e = (\gamma - 1) \left( E_\varepsilon - \frac{|\mathbf{m}|^2}{\rho} \right)$$

## Euler system - conservative variables

$$\begin{aligned} \partial_t \rho + \operatorname{div}_x \mathbf{m} &= 0, \\ \partial_t \mathbf{m} + \operatorname{div}_x \left( \frac{\mathbf{m} \otimes \mathbf{m}}{\rho} \right) + \frac{1}{\varepsilon^2} \nabla_x p &= \frac{1}{\varepsilon^2} \rho \nabla_x F \\ \partial_t E_\varepsilon + \operatorname{div}_x \left[ \left( E_\varepsilon + \frac{1}{\varepsilon^2} p \right) \frac{\mathbf{m}}{\rho} \right] &= \frac{1}{\varepsilon^2} \nabla_x F \cdot \mathbf{m} \end{aligned}$$

# Solution class for the primitive system

## Classical solutions

Existence on a short time interval the length of which may depend on  $\varepsilon$ .  
Results of this type by Klainerman, Majda, Schochet, Alazard and many others

## Weak solutions

Global existence not known. Problem is ill posed e.g. in  $L^\infty$  for some initial data

# Why to go measure-valued?

## Motto: The larger (class) the better

- Universal limits of `numerical` schemes
- Limits of more complex physical systems - vanishing viscosity/heat conductivity limit
- Global existence for “any” data

## Weak-strong uniqueness

A (DMV) solution coincides with a smooth solution with the same initial data as long as the latter solution exists

# Entropy

## Gibbs' relation

$$\vartheta Ds = De + pD\left(\frac{1}{\varrho}\right)$$

## Entropy equation (inequality)

$$\partial_t(\varrho s) + \operatorname{div}_x(\varrho s \mathbf{u}) \boxed{\geq} 0$$

$$\partial_t(\varrho s) + \operatorname{div}_x(s \mathbf{m}) \geq 0.$$

## Thermodynamic stability - standard variables

$$\frac{\partial p}{\partial \varrho} > 0, \quad \frac{\partial e}{\partial \vartheta} > 0$$

## Thermodynamic stability - conservative variables

$$\mathcal{S} : (\varrho, \mathbf{m}, E) \mapsto \varrho s(\varrho, \mathbf{m}, E) \text{ concave function}$$

# Entropy renormalization

## Entropy in the polytropic case

$$s = S\left(\frac{p}{\rho^\gamma}\right) = S\left((\gamma - 1) \frac{E - \frac{1}{2} \frac{|m|^2}{\rho}}{\rho^\gamma}\right)$$

## Renormalization

$$S_\chi = \chi \circ s, \quad \partial_t(\rho S_\chi) + \operatorname{div}_x(\rho S_\chi \mathbf{u}) \geq 0$$

## Levels of renormalization

- “Isentropic”

$$\chi''(s) \leq 0$$

- Standard dissipative

$$\chi'(s) \geq 0, \quad \chi''(s) \leq 0$$

- Vanishing viscosity

$$\chi(s) = s$$



# Relative energy

## Ballistic free energy

$$H_{\Theta}(\varrho, \vartheta) = \varrho e(\varrho, \vartheta) - \Theta \varrho s(\varrho, \vartheta)$$

## Relative energy in standard variables

$$\begin{aligned} & \mathcal{E}(\varrho, \vartheta, \mathbf{u} \mid \tilde{\varrho}, \tilde{\vartheta}, \tilde{\mathbf{u}}) \\ &= \frac{1}{2} \varrho |\mathbf{u} - \tilde{\mathbf{u}}|^2 + H_{\tilde{\vartheta}}(\varrho, \vartheta) - \frac{\partial H(\tilde{\varrho}, \tilde{\vartheta})}{\partial \varrho} (\varrho - \tilde{\varrho}) - H_{\tilde{\vartheta}}(\tilde{\varrho}, \tilde{\vartheta}) \end{aligned}$$

## Relative energy in the conservative variables

$$\begin{aligned} & \mathcal{E}(\varrho, E, \mathbf{m} \mid \tilde{\varrho}, \tilde{E}, \tilde{\mathbf{m}}) \\ &= -\tilde{\vartheta} \left[ S(\varrho, E, \mathbf{m}) - \nabla_{\varrho, E, \mathbf{m}} S(\tilde{\varrho}, \tilde{E}, \tilde{\mathbf{m}}) \cdot (\varrho - \tilde{\varrho}, E - \tilde{E}, \mathbf{m} - \tilde{\mathbf{m}}) \right. \\ & \quad \left. - S(\tilde{\varrho}, \tilde{E}, \tilde{\mathbf{m}}) \right], \quad e(\tilde{\varrho}, \tilde{\vartheta}) = \tilde{E} - \frac{1}{2} \frac{|\tilde{\mathbf{m}}|^2}{\tilde{\varrho}} \end{aligned}$$

## Relative energy inequality

$$\begin{aligned} & \left[ \int_{\Omega} \mathcal{E}(\varrho, \mathbf{m}, E | \tilde{\varrho}, \tilde{\vartheta}, \tilde{\mathbf{u}}) \, dx \right]_{t=0}^{t=\tau} \\ & \leq - \int_0^{\tau} \int_{\Omega} \left[ \varrho s(\varrho, \mathbf{m}, E) \partial_t \tilde{\vartheta} + s(\varrho, \mathbf{m}, E) \mathbf{m} \cdot \nabla_x \tilde{\vartheta} \right] \, dx dt \\ & \quad + \int_0^{\tau} \int_{\Omega} \left[ (\varrho \tilde{\mathbf{u}} - \mathbf{m}) \cdot \partial_t \tilde{\mathbf{u}} + \frac{(\varrho \tilde{\mathbf{u}} - \mathbf{m}) \otimes \mathbf{m}}{\varrho} : \nabla_x \tilde{\mathbf{u}} \right] \, dx dt \\ & \quad - (\gamma - 1) \int_0^{\tau} \int_{\Omega} \left[ \left( E - \frac{1}{2} \frac{|\mathbf{m}|^2}{\varrho} \right) \operatorname{div}_x \tilde{\mathbf{u}} \right] \, dx dt \\ & \quad + \int_0^{\tau} \int_{\Omega} \left[ \varrho \partial_t \tilde{\vartheta} s(\tilde{\varrho}, \tilde{\vartheta}) + \mathbf{m} \cdot \nabla_x \tilde{\vartheta} s(\tilde{\varrho}, \tilde{\vartheta}) \right] \, dx dt \\ & \quad + \int_0^{\tau} \int_{\Omega} \left[ (\tilde{\varrho} - \varrho) \frac{1}{\tilde{\varrho}} \partial_t p(\tilde{\varrho}, \tilde{\vartheta}) - \mathbf{m} \cdot \frac{1}{\tilde{\varrho}} \nabla_x p(\tilde{\varrho}, \tilde{\vartheta}) \right] \, dx dt \end{aligned}$$

# Dissipative measure-valued (DMV) solutions

## Parameterized measure

$$\underbrace{\mathcal{F}}_{\text{phase space}} = \left\{ \varrho \geq 0, \mathbf{m} \in R^3, E \in [0, \infty) \right\}, \quad \underbrace{Q_T}_{\text{physical space}} = (0, T) \times \Omega$$
$$\{\mathcal{V}_{t,x}\}_{(t,x) \in Q_T}, Y_{t,x} \in \mathcal{P}(\mathcal{F})$$

## Field equations

$$\partial_t \langle \mathcal{V}_{t,x}; \varrho \rangle + \operatorname{div}_x \langle \mathcal{V}_{t,x}; \mathbf{m} \rangle = 0$$

$$\partial_t \langle \mathcal{V}_{t,x}; \mathbf{m} \rangle + \operatorname{div}_x \left\langle \mathcal{V}_{t,x}; \frac{\mathbf{m} \otimes \mathbf{m}}{\varrho} \right\rangle + \nabla_x \langle \mathcal{V}_{t,x}; p \rangle = D_x \mu_C$$

$$\partial_t \int_{\Omega} \langle \mathcal{V}_{t,x}; E \rangle dx + \mathcal{D} = 0, \quad \partial_t \langle \mathcal{V}_{t,x}; \varrho s \rangle + \operatorname{div}_x \langle \mathcal{V}_{t,x}; s \mathbf{m} \rangle \geq 0$$

## Compatibility

$$\int_0^T \int_{\Omega} |\mu_C| dx dt \leq C \int_0^T \mathcal{D} dt$$

## Relative energy inequality for (DMV) solutions

$$\begin{aligned}
 & \left[ \int_{\Omega} \langle \mathcal{V}_{t,x}; \mathcal{E}(\varrho, \mathbf{m}, E | \tilde{\varrho}, \tilde{\vartheta}, \tilde{\mathbf{u}}) \rangle dx \right]_{t=0}^{t=\tau} + \mathcal{D}(\tau) \\
 & \leq - \int_0^{\tau} \int_{\Omega} \left[ \langle \mathcal{V}_{t,x}; \varrho s(\varrho, \mathbf{m}, E) \rangle \partial_t \tilde{\vartheta} + \langle \mathcal{V}_{t,x}; s(\varrho, \mathbf{m}, E) \mathbf{m} \rangle \cdot \nabla_x \tilde{\vartheta} \right] dx dt \\
 & \quad + \int_0^{\tau} \int_{\Omega} \left[ \langle \mathcal{V}_{t,x}; \varrho \tilde{\mathbf{u}} - \mathbf{m} \rangle \cdot \partial_t \tilde{\mathbf{u}} + \left\langle \mathcal{V}_{t,x}; \frac{(\varrho \tilde{\mathbf{u}} - \mathbf{m}) \otimes \mathbf{m}}{\varrho} \right\rangle : \nabla_x \tilde{\mathbf{u}} \right] dx dt \\
 & \quad - (\gamma - 1) \int_0^{\tau} \int_{\Omega} \left[ \left\langle \mathcal{V}_{t,x}; E - \frac{1}{2} \frac{|\mathbf{m}|^2}{\varrho} \right\rangle \operatorname{div}_x \tilde{\mathbf{u}} \right] dx dt \\
 & \quad + \int_0^{\tau} \int_{\Omega} \left[ \langle \mathcal{V}_{t,x}; \varrho \rangle \partial_t \tilde{\vartheta} s(\tilde{\varrho}, \tilde{\vartheta}) + \langle \mathcal{V}_{t,x}; \mathbf{m} \rangle \cdot \nabla_x \tilde{\vartheta} s(\tilde{\varrho}, \tilde{\vartheta}) \right] dx dt \\
 & \quad + \int_0^{\tau} \int_{\Omega} \left[ \langle \mathcal{V}_{t,x}; \tilde{\varrho} - \varrho \rangle \frac{1}{\tilde{\varrho}} \partial_t p(\tilde{\varrho}, \tilde{\vartheta}) - \langle \mathcal{V}_{t,x}; \mathbf{m} \rangle \cdot \frac{1}{\tilde{\varrho}} \nabla_x p(\tilde{\varrho}, \tilde{\vartheta}) \right] dx dt \\
 & \quad + \int_0^{\tau} \int_{\Omega} \nabla_x \tilde{\mathbf{u}} : d\mu_C,
 \end{aligned}$$

# Limits of Euler flows with strong stratification

## Scaled Euler system

$$\begin{aligned}\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) &= 0 \\ \partial_t(\varrho \mathbf{u}) + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \frac{1}{\varepsilon^2} \nabla_x \rho(\varrho, \vartheta) &= \frac{1}{\varepsilon^2} \varrho \nabla_x F, \\ \partial_t \left( \frac{1}{2} \varrho |\mathbf{u}|^2 + \frac{1}{\varepsilon^2} \varrho e(\varrho, \vartheta) \right) + \operatorname{div}_x \left[ \left( \frac{1}{2} \varrho |\mathbf{u}|^2 + \frac{1}{\varepsilon^2} \varrho e(\varrho, \vartheta) \right) \mathbf{u} \right] \\ + \operatorname{div}_x \left( \frac{1}{\varepsilon^2} \rho(\varrho, \vartheta) \mathbf{u} \right) &= \frac{1}{\varepsilon^2} \varrho \nabla_x F \cdot \mathbf{u}.\end{aligned}$$

## Geometry

$\Omega = \mathcal{T}^2 \times (0, 1)$ ,  $\mathcal{T}^2 = [0, 1] \times [0, 1]$  – the two dimensional torus

$$\mathbf{u} \cdot \mathbf{n}|_{\partial\Omega} = 0$$

# Isothermal case

## Boyle-Marriot EoS - stationary problem

$$p = \rho\vartheta, \quad F = F(z) = -z, \quad \nabla_x p = \rho \nabla_x F$$

## Isothermal case

$$\nabla_x(\rho_s \bar{\Theta}) = -\rho_s \nabla_x F, \quad \rho_s = \exp\left(-\frac{z}{\bar{\Theta}}\right), \quad \bar{\Theta} > 0$$

## Well-prepared initial data

$$\rho_{0,\varepsilon} = \rho_s + \varepsilon \rho_{0,\varepsilon}^{(1)}, \quad \vartheta_{0,\varepsilon} = \bar{\Theta} + \varepsilon \vartheta_{0,\varepsilon}^{(1)}, \quad \mathbf{u}_{0,\varepsilon}$$

$$\|\rho_{0,\varepsilon}^{(1)}\|_{L^\infty(\Omega)} + \|\vartheta_{0,\varepsilon}^{(1)}\|_{L^\infty(\Omega)} + \|\mathbf{u}_{0,\varepsilon}\|_{L^\infty(\Omega; \mathbb{R}^N)} \leq c,$$

$$\rho_\varepsilon^{(1)} \rightarrow 0, \quad \vartheta_{0,\varepsilon}^{(1)} \rightarrow 0, \quad \mathbf{u}_{0,\varepsilon} \rightarrow \mathbf{U}_0 \text{ in } L^1(\Omega) \text{ as } \varepsilon \rightarrow 0,$$

$$\mathbf{U}_0 \in W^{k,2}(\Omega; \mathbb{R}^3), \quad k > 3, \quad \mathbf{U}_0 = [U_0^1, U_0^2, 0], \quad \operatorname{div}_h \mathbf{U}_0 = 0.$$

## Target problem - isothermal case

### Limit velocity

$$\mathbf{U} = \mathbf{U}(t, x_h, z), \quad x_h \in \mathcal{T}^2, \quad z \in (0, 1)$$

### Incompressible Euler system in 2D

$$\partial_t \mathbf{U} + \mathbf{U} \cdot \nabla_h \mathbf{U} + \nabla_x \Pi = 0, \quad \operatorname{div}_h \mathbf{U} = 0 \text{ in } \mathcal{T}^2, \quad z \text{ fixed}$$

### Stratified initial data

$$\mathbf{U}(0, x) = \mathbf{U}_0(x_h, z) = [U_0^1(x_h, z), U_0^2(x_h, z), 0]$$

## Singular limit (MV) $\rightarrow$ strong

### Convergence to the target system

Let  $\{\mathcal{V}_{t,x}^\varepsilon\}_{(t,x) \in (0,T) \times \Omega}$ ,  $\mathcal{D}^\varepsilon$  be a family of dissipative measure-valued solutions to the scaled system scaled Euler system, with the well prepared initial data

$$\mathcal{V}_{0,x}^\varepsilon = \delta_{\varrho_{0,\varepsilon}, \varrho_{0,\varepsilon}} \mathbf{u}_{0,\varepsilon, c_V \varrho_{0,\varepsilon}} \vartheta_{0,\varepsilon}.$$

Then

$$\mathcal{D}^\varepsilon \rightarrow 0 \text{ in } L^\infty(0, T),$$

and

$$\mathcal{V}^\varepsilon \rightarrow \delta_{\varrho_s, \varrho_s} \mathbf{u}_{c_V \varrho_s} \bar{\Theta} \text{ in } L^\infty(0, T; L^1(\Omega; \mathcal{M}^+(\mathcal{F})_{\text{weak-}(\ast)})),$$

where  $[\varrho_s, \bar{\Theta}]$  is the static state and  $\mathbf{U}$  is the unique solution to the incompressible 2D Euler system



# Isentropic case

## Stationary problem

$$p = \varrho \vartheta, \quad F = F(z) = -z$$

$$s_s(\varrho_s, \vartheta_s) = \bar{s}, \quad \exp((\gamma - 1)\bar{s}) \nabla_x(\varrho_s^\gamma) = -\varrho_s \nabla_x F$$

$$p_s = \exp((\gamma - 1)\bar{s}) \varrho_s^\gamma$$

## Target problem - isentropic case

### Anelastic Euler system

$$\partial_t \mathbf{U} + \mathbf{U} \cdot \nabla \mathbf{U} + \nabla_x \Pi = 0, \quad \operatorname{div}(\varrho_s \mathbf{U}) = 0, \quad x \in \Omega$$

## Singular limit (MV) $\rightarrow$ strong

Let

$$\mathbf{U}_0 \in W^{k,2}(\Omega; \mathbb{R}^N), k \geq N, \operatorname{div}_x(\varrho_s \mathbf{U}_0) = 0, \mathbf{U}_0 \cdot \mathbf{n}|_{\partial\Omega} = 0.$$

Suppose that the anelastic Euler system with initial datum  $\mathbf{U}_0$  admits a unique strong solution  $\mathbf{U}$  defined on a maximal time interval  $[0, T_{\max})$ .

Let  $\{\mathcal{V}_{t,x}^\varepsilon\}_{(t,x) \in (0,T) \times \Omega}$ ,  $0 < T < T_{\max}$ , be a family of (DMV) solutions such that

$$\mathcal{V}_{0,x}^\varepsilon \left\{ [\varrho, \mathbf{m}, p] \left| \left| \frac{\varrho - \varrho_s}{\varepsilon} \right| + \left| \frac{\mathbf{m}}{\varrho} - \mathbf{U}_0 \right| + \left| \frac{p - p_s}{\varepsilon} \right| \leq M_\varepsilon(x) \right\} = 1 \text{ for a.a. } x \in \Omega,$$

where

$$\|M_\varepsilon\|_{L^\infty(\Omega)} \leq c \quad \text{and} \quad M_\varepsilon \rightarrow 0 \text{ in } L^1(\Omega) \text{ as } \varepsilon \rightarrow 0.$$

Suppose that the initial entropy satisfies

$$\mathcal{V}_{0,x}^\varepsilon \left\{ [\varrho, \mathbf{m}, p] \left| \bar{s} - \varepsilon^{2+\alpha} \leq s(\varrho, p) \leq \bar{s} + \varepsilon^{2+\alpha} \right. \right\} = 1 \text{ for a.a. } x \in \Omega, \alpha > 0$$

Then

$$\mathcal{D}^\varepsilon \rightarrow 0 \text{ in } L^\infty(0, T) \text{ as } \varepsilon \rightarrow 0,$$

and

$$\mathcal{V}^\varepsilon \rightarrow \delta_{\bar{\varrho}, \bar{\mathbf{u}}, \bar{p}} \text{ in } L^\infty(0, T, L^q(\Omega; \mathcal{M}^+(\mathcal{F}))) \text{ as } \varepsilon \rightarrow 0 \text{ for any } 1 \leq q < \infty$$

### III prepared initial data

#### Isentropic Euler system

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u}) = 0$$

$$\partial_t \varrho + \operatorname{div}_x(\varrho \mathbf{u} \otimes \mathbf{u}) + \frac{1}{\varepsilon^2} \nabla_x p(\varrho) = 0$$

$$\Omega = \mathbb{R}^N, \quad N = 2, 3, \quad \varrho \rightarrow \bar{\varrho} > 0, \quad \mathbf{u} \rightarrow 0 \text{ as } |x| \rightarrow \infty$$

#### III prepared initial data

$$\varrho(0, \cdot) = \varrho_{0,\varepsilon}, \quad \varrho_{0,\varepsilon} \approx \bar{\varrho} + \varepsilon \boxed{s_\varepsilon} \rightarrow s \text{ in } L^1 \cap L^\infty(\mathbb{R}^N)$$

$$\mathbf{u}(0, \cdot) = \mathbf{u}_{0,\varepsilon} \rightarrow \mathbf{U}_0 + \boxed{\nabla_x \Phi} \text{ in } L^2 \cap L^\infty(\mathbb{R}^N, \mathbb{R}^N)$$

# Target system, acoustic waves

## Incompressible Euler system

$$\partial_t \mathbf{U} + \operatorname{div}_x(\mathbf{U} \otimes \mathbf{U}) + \nabla_x \Pi = 0, \operatorname{div}_x \mathbf{U} = 0 \text{ in } R^N$$

## Acoustic waves

$$\varepsilon \partial_t s_\varepsilon + \operatorname{div}_x(\bar{\rho} \nabla_x \Phi) = 0$$

$$\varepsilon \partial_t \nabla_x \Phi + \frac{p'(\bar{\rho})}{\bar{\rho}} \nabla_x s = 0$$

## Dispersive estimates (Strichartz's estimates)

### Strichartz estimates

$$\begin{aligned} & \|s(\tau, \cdot)\|_{L^p(R^3)}^2 + \|\nabla_x \Phi(\tau, \cdot)\|_{L^p(R^3; R^3)}^2 \\ & \leq c \left(1 + \frac{\tau}{\varepsilon}\right)^{(N-1)\left(\frac{1}{p} - \frac{1}{q}\right)} \left[ \|s_0\|_{W^{k,q}(R^3)}^2 + \|\nabla_x \Phi_0\|_{W^{k,q}(R^3; R^3)}^2 \right], \\ & k \geq N \left(\frac{1}{q} - \frac{1}{p}\right), \quad 2 \leq p \leq \infty, \quad \frac{1}{p} + \frac{1}{q} = 1 \end{aligned}$$