

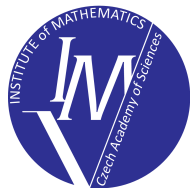
# Convergence of a finite difference MAC scheme for the compressible Navier-Stokes

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$$\begin{aligned}\rho_t + \operatorname{div}(\rho \mathbf{u}) &= 0. \\ (\rho \mathbf{u})_t + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) &= -\nabla p + \nabla \cdot \mathbb{S}\end{aligned}\tag{1}$$

$\rho$  : density

$\mathbf{u}$  : velocity

$p$  : pressure,  $p = a\rho^\gamma$

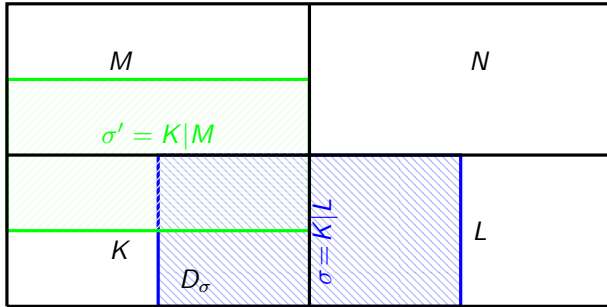
$\mathbb{S}$  : viscous stress,  $\mathbb{S} = \mu \nabla \mathbf{u} + \lambda \operatorname{div} \mathbf{u}$ ,  $\mu > 0$ ,  $\lambda \geq 0$

**Finite Volume-Finite Element** by T. Karper, 2013,  $\gamma > 3$

- E. Feireisl, R. Hošek, D. Maltese, A. Novotný, 2017  
bounded numerical solution
- E. Feireisl & M. Lukáčová-Medvid'ová, 2017  
dissipative measure-valued solution,  $\gamma \in (1, 2)$
- R. Hosěk & BS, 2018; H. Mizerová & BS, in preparation

**Our interests:** Simple FD, stability and convergence





- Elements:  $\mathcal{M}$
- Faces:  $\mathcal{E}$
- Dual Faces:  $\tilde{\mathcal{E}}$

$$\rho, p \in L_{\mathcal{M}}$$

$$\mathbf{u} = (u_1, \dots, u_d) \in H_{\mathcal{E}} = (H_{\mathcal{E}}^{(1)}, \dots, H_{\mathcal{E}}^{(d)})$$

$$D_t \rho_h^n + \operatorname{div}_{\mathcal{M}}^{\text{up}}(\rho_h^n \mathbf{u}_h^n) = 0$$

$$D_t(\widehat{\rho}_h^n(i) u_{i,h}^n) + \operatorname{div}_{\mathcal{E}(i)}^{\text{up}}(\rho_h^n \mathbf{u}_h^n u_{i,h}^n) - \mu \Delta_{\mathcal{E}}^{(i)} u_{i,h}^n - \lambda \bar{\delta}_{\mathcal{E}}^{(i)} \operatorname{div}_{\mathcal{M}} \mathbf{u}_h^n + \bar{\delta}_{\mathcal{E}}^{(i)} p(\rho_h^n) = 0 \quad (2)$$

$$\sigma = \overrightarrow{K|L}, K = [\overrightarrow{\sigma''|\sigma}], \epsilon = \sigma|\sigma' \in \tilde{\mathcal{E}}^{(i)}$$

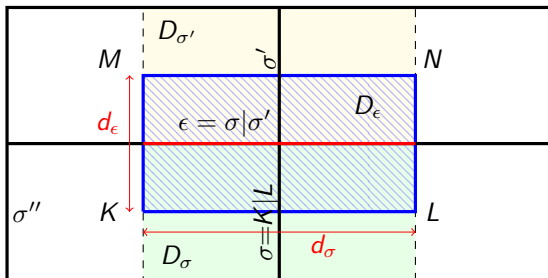
$$D_t f = \frac{f(t) - f(t - \Delta t)}{\Delta t}$$

$$\left(\Delta_{\mathcal{E}}^{(i)} v_i\right)_{D_\sigma} = \frac{1}{|D_\sigma|} \sum_{\epsilon \in \tilde{\mathcal{E}}(D_\sigma)} \frac{|\epsilon|}{d_\epsilon} (v_{\sigma'} - v_\sigma)$$

$$\left(\bar{\delta}_{\mathcal{M}}^{(i)} v_i\right)_K = \frac{v_{\sigma''} - v_\sigma}{h_i}$$

$$\widehat{\rho}^{(i)}(\mathbf{x}) = \sum_{\sigma \in \mathcal{E}^{(i)}} \rho_{D_\sigma} \chi_{D_\sigma}, \quad \rho_{D_\sigma} = \frac{\rho_K |K| + \rho_L |L|}{|D_\sigma|}$$

$$\left(\bar{\delta}_{\mathcal{E}}^{(i)} r\right)_{D_\sigma} = \frac{r_L - r_K}{d_\sigma}$$



# Upwind flux

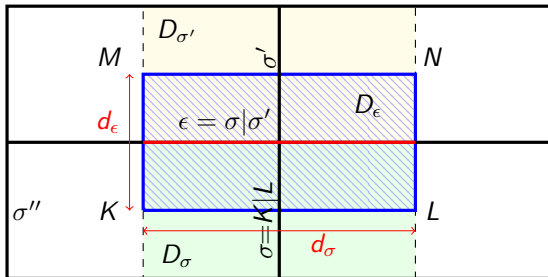
$$\operatorname{div}_{\mathcal{M}}^{\text{up}}(\rho \mathbf{u})_K = \frac{1}{|K|} \sum_{\sigma \in \mathcal{E}(K)} F_{\sigma,K}(\rho, \mathbf{u}), \quad \operatorname{div}_{\mathcal{E}^{(i)}}^{\text{up}}(\rho \mathbf{u} u_i)_{D_\sigma} = \frac{1}{|D_\sigma|} \sum_{\epsilon \in \tilde{\mathcal{E}}(D_\sigma)} F_{\epsilon,\sigma}(\rho, \mathbf{u}) u_\epsilon$$

$$F_{\sigma,K}(\rho, \mathbf{u}) = |\sigma| \rho_\sigma^{\text{up}} u_{\sigma,K}, \quad \forall \sigma = K|L \in \mathcal{E},$$

$$\rho_\sigma^{\text{up}} = \begin{cases} \rho_K & \text{if } u_{\sigma,K} \geq 0, \\ \rho_L & \text{otherwise.} \end{cases} \quad u_{\sigma,K} = u_\sigma \mathbf{e}^{(i)} \cdot \mathbf{n}_{\sigma,K}, \quad \sigma \in \mathcal{E}^{(i)}(K)$$

$$F_{\epsilon,\sigma} = \frac{1}{2} \begin{cases} F_{\sigma,K}(\rho, \mathbf{u}) \mathbf{n}_{\epsilon,D_\sigma} \cdot \mathbf{n}_{\sigma,K} + F_{\sigma',K}(\rho, \mathbf{u}) \mathbf{n}_{\epsilon,D_\sigma} \cdot \mathbf{n}_{\sigma',K} & \epsilon \parallel \sigma \\ F_{\tau,K}(\rho, \mathbf{u}) + F_{\tau',L}(\rho, \mathbf{u}) & \epsilon \perp \sigma \end{cases}$$

$$u_\epsilon \equiv u_{i,\epsilon} = \frac{u_\sigma + u_{\sigma'}}{2}$$



## Lemma 1

Let  $\rho_h^{n-1} \in L_{\mathcal{M}}$ ,  $\mathbf{u}_h^{n-1} \in H_{\mathcal{E}}$  be given;  $\rho_K^{n-1} > 0$  for all  $K \in \mathcal{M}$ . Then the numerical scheme (2) admits a solution

$$\rho_h^n \in L_{\mathcal{M}}, \rho_K^n > 0 \text{ for all } K \in \mathcal{M}, \mathbf{u}_h^n \in H_{\mathcal{E}}.$$

Moreover, it satisfies the discrete conservation of mass

$$\sum_{K \in \mathcal{M}} \int_K \rho_K^n = \sum_{K \in \mathcal{M}} \int_K \rho_K^{n-1}.$$

## Theorem 2 (Energy estimate)

Let  $(\rho_h, \mathbf{u}_h)$  be a solution of (2). Then for any  $n = 1, \dots, N$ , there exist  $\rho^{n-1, n} \in [\min(\rho_h^{n-1}, \rho_h^n), \max(\rho_h^{n-1}, \rho_h^n)]$  and  $\bar{\rho}_\sigma^n \in [\min(\rho_K^n, \rho_L^n), \max(\rho_K^n, \rho_L^n)]$  for any  $\sigma = K|L \in \mathcal{E}$  such that

$$\begin{aligned} & \frac{1}{\Delta t} \int_{\Omega} \left( \mathcal{H}(\rho_h^n) + \frac{1}{2} \rho_h^n |\mathbf{u}_h^n|^2 \right) dx - \frac{1}{\Delta t} \int_{\Omega} \left( \mathcal{H}(\rho_h^{n-1}) + \frac{1}{2} \rho_h^{n-1} |\mathbf{u}_h^{n-1}|^2 \right) dx \\ & + \mu \|\nabla_h \mathbf{u}_h^n\|_{L^2(\Omega)}^2 + \lambda \|\operatorname{div}_{\mathcal{M}}^{\text{up}} \mathbf{u}_h^n\|_{L^2(\Omega)}^2 + \frac{1}{2\Delta t} \int_{\Omega} \rho_h^{n-1} |\mathbf{u}_h^n - \mathbf{u}_h^{n-1}|^2 dx \\ & + \int_{\Omega} \frac{1}{2\Delta t} \mathcal{H}''(\rho^{n-1, n}) (\rho_h^n - \rho_h^{n-1})^2 dx + \frac{1}{2} \sum_{\sigma \in \mathcal{E}^{(i)}} |\sigma| |\mathbf{u}_{\sigma, K}^n| \mathcal{H}''(\bar{\rho}_\sigma^n) [\rho_h]_{\sigma}^2 = 0 \end{aligned}$$

where  $\mathcal{H}(\rho) = \frac{1}{\gamma-1} \rho(\rho)$ .

## Lemma 3

Let  $\rho_h^n, \mathbf{u}_h^n$  be the solution to the numerical scheme (2) and suppose  $\rho_h^n < \bar{\rho}$ . Then

$$\int_{\Omega} \partial_h^t \rho_h^n \phi dx - \int_{\Omega} \rho_h^n \mathbf{u}_h^n \cdot \nabla_x \phi dx = \mathcal{O}(h)$$

$$\begin{aligned} \int_{\Omega} \partial_h^t (\rho_h \mathbf{u}_h)^n \cdot \mathbf{v} dx - \int_{\Omega} \rho_h^n \mathbf{u}_h^n \otimes \mathbf{u}_h^n : \nabla_x \mathbf{v} dx - \int_{\Omega} p(\rho_h^n) \operatorname{div}_x \mathbf{v} dx \\ + \mu \int_{\Omega} (\nabla_h \mathbf{u}_h^n) : \nabla_x \mathbf{v} dx = \mathcal{O}(h) \end{aligned}$$





## Definition 4 (DMVS<sup>1</sup>)

We say that a parameterized measure  $\{\nu_{t,x}\}_{(t,x) \in (0,T) \times \Omega}$ ,

$$\nu \in L_{weak}^{\infty} \left( (0, T) \times \Omega; \mathcal{P}([0, \infty) \times \mathbb{R}^N) \right)$$

is a dissipative measure-valued solution of the Navier-Stokes system in  $(0, T) \times \Omega$ , if the following holds for a.a.  $\tau \in (0, T)$ , for any  $\psi \in C^1((0, T) \times \Omega; \mathbb{R}^d)$

$$\begin{aligned} \left[ \int_{\Omega} \langle \nu_{\tau,x}; \rho \rangle \psi(\tau, \cdot) dx \right]_{t=0}^{t=\tau} &= \int_0^{\tau} \int_{\Omega} [\langle \nu_{t,x}; \rho \rangle \partial_t \psi + \langle \nu_{t,x}; \rho \mathbf{u} \rangle \cdot \nabla_x \psi] dx dt \\ \left[ \int_{\Omega} \langle \nu_{\tau,x}; \rho \mathbf{u} \rangle \psi(\tau, \cdot) dx \right]_{t=0}^{t=\tau} &= \int_0^{\tau} \int_{\Omega} [\langle \nu_{t,x}; \rho \mathbf{u} \rangle \partial_t \psi + \langle \nu_{t,x}; \rho \mathbf{u} \otimes \mathbf{u} + p(\rho) \mathbf{I} \rangle : \nabla_x \psi] dx dt \\ &\quad - \int_0^{\tau} \int_{\Omega} \mathbb{S} : \nabla_x \psi dx dt, + \int_0^{\tau} \int_{\Omega} \mathcal{R} : \nabla_x \psi dx dt \\ \left[ \int_{\Omega} \langle \nu_{\tau,x}; \mathbf{E} \rangle \psi(\tau, \cdot) dx \right]_{t=0}^{t=\tau} &+ \int_0^{\tau} \int_{\Omega} \mathbb{S} : \nabla_x \psi dx dt + \mathcal{D}(\tau) \leq 0, \end{aligned}$$

where

$$\int_0^{\tau} \|\mathcal{R}\|_{\mathcal{M}(\Omega)} dt \leq \int_0^{\tau} \mathcal{D}(\tau) dt$$

<sup>1</sup>Feireisl et.al. Dissipative measure-valued solutions to the compressible Navier–Stokes system. Calc. Vari. Partial Differ. Equ. 2016

## Theorem 5

Let  $[\rho_h, \mathbf{u}_h]$  be a family of numerical solutions obtained by the scheme (2) with suitable initial data and suppose  $\rho_h^n < \bar{\rho}$ .

Then any Young measure  $\nu_{t,x}$  generated by the  $[\rho_h, \mathbf{u}_h]$  for  $h \rightarrow 0$  represents a dissipative measure-valued solution of NS (1).

Remark: Applying the weak-strong uniqueness (Feireisl et.al. 2016) we conclude the convergence to the strong solution.

Thank you for your attention!

