

§ 1. Kinematics of Fluids

Let $(T_1, T_2) \subset \mathbb{R}^1$ be a time interval

$\Omega_t \subset \mathbb{R}^3$ domain occupied by the fluid
at time $t \in (T_1, T_2)$

Hypothesis ; exactly one fluid particle passes through each point $x \in \Omega_t$ at any time t

Lagrangian Description of the Flow

consider the motion of each particular fluid particle

then

the trajectory of particles

(1.1) $x = \varphi(X, t)$

$(x_i = \varphi_i(X, t), i = 1, 2, 3)$

X. the reference determining the particle under consideration (usually is described in some coord. system by parameters X_1, X_2, X_3 which, together with time are called Lagrangian coordinates

(1.1) determines the position of the particle given by the reference X at time t .

Sometimes we use a more detailed description of the motion of fluid particles

$$(1.2) \quad X = \Psi(X, t_0; t)$$

which determines, at time t , the position x of the particle passing through the point X at time t_0 . Then

$$(1.3) \quad X = \Psi(X, t_0; t_0)$$

provided the references are identical with the coordinates of particles at time t_0 .

The Lagrangian description is used

if we study the flow of a piece of fluid

formed by the same particles at each

time instant and filling a domain

$\Omega(t) \subset \mathbb{R}^3$ at time t .

The velocity

$$(1.4) \vec{v}(X, t) = \frac{\partial \varphi}{\partial t}(X, t)$$

The acceleration

$$(1.5) \vec{a}(X, t) = \frac{\partial^2 \varphi}{\partial t^2}(X, t)$$

Eulerian Description

= based on the determination of the velocity $v(x, t)$ of the fluid particle passing through the point x at time t .

We can write

$$(1.6) \quad v(x, t) = \vec{v}(X, t) = \frac{\partial \varphi}{\partial t}(X, t),$$

where $x = \varphi(X, t)$.

Since it is not important what the fluid particle considered is given by the reference x , we will use the simplified notation

$$\tilde{\varphi}(t) = \varphi(x, t)$$

Then $\xrightarrow{(1.7)}$ $v(x, t) = \frac{d\tilde{\varphi}}{dt}(t)$, $x = \tilde{\varphi}(t) = (\tilde{\varphi}_1(t), \tilde{\varphi}_2(t), \tilde{\varphi}_3(t))$

time t , together with coordinates x_1, x_2, x_3

as Eulerian coordinates.

We suppose that the domain of definition of the mapping $v = v(t)$ is the set M :

$$(1.8) \quad \begin{cases} v: M \rightarrow \mathbb{R}^3 \\ M = \{x, t; x \in \Omega_t, t \in (T_1, T_2)\} \subset \mathbb{R}^4 \end{cases}$$

where Ω_t depends continuously on time t ,

the set M is open.

assume

$$(1.9) \quad v \in [C^1(M)]^3$$

Acceleration: $a(x, t) = \frac{d^2 \tilde{\varphi}}{dt^2}(t) = \frac{d}{dt} [v(\tilde{\varphi}(t), t)] =$

$$\boxed{v_x = \tilde{\varphi}(t)}$$

$$= \frac{\partial v}{\partial t} (\tilde{\varphi}(t), t) + \sum_{i=1}^3 \frac{\partial v}{\partial x_i} (\tilde{\varphi}(t), t) \frac{d\tilde{\varphi}^i(t)}{dt}$$

Hence

$$(1.9) \quad a(x,t) = \frac{\partial v}{\partial t}(x,t) + \sum_{i=1}^3 v_i(x,t) \frac{\partial v}{\partial x_i}(x,t),$$

$$(1.10) \quad a = \frac{\partial v}{\partial t} + (v \cdot \nabla) v$$

$$(1.11) \quad \frac{d}{dt} = \underbrace{\frac{\partial}{\partial t}}_{\text{local derivative}} + \underbrace{v \cdot \text{grad}}_{\text{convective derivative}} \quad \text{material derivative}$$

Remark:

In the Lagrangian conception this quantity is viewed as a function $F(x,t)$ determining the value of the quantity considered, found for the particle given by the reference x , at time t .

If we use the Eulerian description, the quantity is represented by a function F dependent on x and t . Then $F(x,t)$ denotes

The value of the quantity at the point x and time t .

Describing the path of a fluid particle by the equation $x = \Psi(x, t) = \tilde{\Psi}(t)$, we can write

$$(1.12) \quad \vec{F}(x, t) = \tilde{F}(t) = F(\Psi(x, t), t) = F(\tilde{\Psi}(t), t)$$

We express the rate of change of the quantity F bound to this particle as the time derivative of the function \tilde{F} .

$$(1.13) \quad \frac{d\tilde{F}}{dt}(t) = \frac{\partial F}{\partial t}(x, t)$$

$$= \frac{\partial F}{\partial t}(\tilde{\Psi}(t), t) + \sum \frac{\partial F}{\partial x_i}(\tilde{\Psi}(t), t) \frac{d\tilde{\Psi}_i}{dt}(t)$$

$$= \frac{\partial F}{\partial t}(x, t) + v(x, t) \cdot \text{grad} F(x, t)$$

⇒ rate of change of the quantity F
= the material derivative of F

⇒ the local derivative $\frac{\partial F}{\partial t} \Leftarrow$

from dependence of F on time t

⇒ $v \cdot \nabla F \dots \Leftarrow$ transport of the quantity F by the moving fluid

The material derivative is called the derivative along a trajectory of a fluid particle.

Relation between the Lagrangian and Eulerian Description of Flow

Lagr. descr. \rightarrow Euler. descr.

$$(1.14) \quad v(x,t) = \vec{v}(x,t) = \frac{\partial \Phi}{\partial t}(x,t); \quad \Phi(x,t) \left\{ \begin{array}{l} \text{mapping} \\ \text{denote } \Phi(x,t) \end{array} \right.$$

$$(1.14) \quad x = \Phi(X,t) \Leftrightarrow X = \Phi(x,t)$$

$$(1.15) \quad v(x,t) = \vec{v}(\Phi(x,t), t)$$

In case the position of the particle at time t_0 is known for its reference X and the Lagrangian description in the form

$$x = \Phi(X, t_0; t), \quad \text{then mapping}$$

$$\Phi(x,t) \text{ is the inverse mapping to } \Phi(\cdot, t_0; t).$$

The transition from the Eulerian description to Lagrangian is equivalent to the determination of the paths of fluid particles on the basis of a given velocity field $v(x, t)$. The trajectory of the fluid particle passing through a point $x \in M$ at time $t_0 \in (T_1, T_2)$ is given as the solution of the initial value problem

$$(1.16) \quad \frac{dx}{dt} = v(x, t), \quad x(t_0) = x.$$

1. Theorem:

Under assumption $v \in [C^1(M)]^3$ the following statements hold:

- (1) For each $(x, t_0) \in M$ the problem (1.15) has exactly one maximal solution $\varphi(x, t_0; t)$ (defined for t from a certain interval $(\alpha_{x, t_0}, \beta_{x, t_0})$).
- (2) The mapping φ has continuous first order partial derivatives

with respect to x_1, x_2, x_3, t_0, t and

L9.

Continuous derivative $\frac{\partial^2 \varphi}{\partial t \partial x_i}, \frac{\partial^2 \varphi}{\partial t_0 \partial x_i}$ $i=1, 2, 3, 4$

in domain of domain of definition

$\{(x_i, t_0, t); (x_i, t_0) \in M, t \in (\alpha_{x_i, t_0}, \beta_{x_i, t_0})\}$.