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A Stabilized FE-FV scheme for the compressible Navier-Stokes-Fourier system

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- Part I: Introduction
- Part II: Numerical method
- Part III: Stability
- Part IV: Test

Introduction - - Modelling I

Navier-Stokes-Fourier system for the compressible, viscous and heat conducting flow.

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0. \quad (1a)$$

$$(\rho \mathbf{u})_t + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) = -\nabla p + \operatorname{div} \mathbb{S}(\nabla \mathbf{u}). \quad (1b)$$

$$c_v((\rho \theta)_t + \operatorname{div}(\rho \theta \mathbf{u})) + \operatorname{div} \mathbf{q}(\theta, \nabla \theta) = \mathbb{S}(\nabla \mathbf{u}) : \nabla \mathbf{u} - \theta \frac{\partial p(\rho, \theta)}{\partial \theta} \operatorname{div} \mathbf{u}. \quad (1c)$$

$\rho, p, \mathbf{u}, \theta$ are the fluid density, pressure, velocity and temperature.

$c_v > 0$ is the specific heat per volume.

$$\boxed{\partial_t \left(\frac{1}{2} \rho \mathbf{u}^2 + \rho e \right) + \operatorname{div} \left(u \left(\frac{1}{2} \rho \mathbf{u}^2 + \rho e + p \right) \right) = \dots}$$

Boundary condition

$$\mathbf{u}|_{\partial \Omega} = 0, \quad \mathbf{n} \cdot \nabla \theta|_{\partial \Omega} = 0. \quad (1d)$$

Initial values are

$$\rho(\mathbf{x}, 0) = \rho_0 > 0, \quad \mathbf{u}(\mathbf{x}, 0) = \mathbf{u}_0, \quad \theta(\mathbf{x}, 0) = \theta_0 > 0. \quad (1e)$$

$$\mathbb{S} = \mu(\nabla \mathbf{u} + \nabla \mathbf{u}^T - \frac{2}{3}\operatorname{div} \mathbf{u} \mathbf{I}) + \eta \operatorname{div} \mathbf{u} \mathbf{I} = 2\mu \mathbf{D}(\mathbf{u}) + \nu \operatorname{div} \mathbf{u} \mathbf{I},$$

$$\operatorname{div} \mathbb{S}(\nabla \mathbf{u}) = 2\mu \operatorname{div} \mathbf{D}(\mathbf{u}) + \nu \nabla \operatorname{div} \mathbf{u}, \quad \mathbf{D}(\mathbf{u}) = \frac{\nabla \mathbf{u} + \nabla \mathbf{u}^T}{2},$$

$$\mathbb{S}(\nabla \mathbf{u}) : \nabla \mathbf{u} = 2\mu |\mathbf{D}(\mathbf{u})|^2 + \nu |\operatorname{div} \mathbf{u}|^2, \quad \mu > 0, \quad \nu = \eta - \frac{2}{3}\mu > 0.$$

The heat flux \mathbf{q} follows Fourier's law

$$\mathbf{q} = -\kappa(\theta) \nabla \theta = -\nabla \mathcal{K}(\theta), \quad \operatorname{div} \mathbf{q} = -\Delta \mathcal{K}, \quad \mathcal{K}(\theta) = \int_0^\theta \kappa(z) dz.$$

$$p = p^e + p^t, \quad p^e = a\rho^\gamma + b\rho, \quad p^t = \rho\theta, \quad a, b > 0.$$

$$\frac{\partial p}{\partial \theta} = \frac{\partial p^t}{\partial \theta} = \rho.$$

Convergence to weak solution.

E. Feireisl. Dynamics of viscous compressible fluids. Oxford University Press, Oxford, 2004.

E. Feireisl, T. Karper, A. Novotny. A convergent numerical method for the NavierStokesFourier system. Preprint.

$$\operatorname{div} \mathbb{S}(\nabla \mathbf{u}) = \mu \Delta \mathbf{u} + \lambda \nabla \operatorname{div} \mathbf{u}, \quad \mathbb{S}(\nabla \mathbf{u}) : \nabla \mathbf{u} = \mu |\nabla \mathbf{u}|^2 + \lambda |\operatorname{div} \mathbf{u}|^2$$

$$\operatorname{div} \mathbb{S}(\nabla \mathbf{u}) = 2\mu \operatorname{div} \mathbf{D}(\mathbf{u}) + \nu \nabla \operatorname{div} \mathbf{u}, \quad \mathbb{S}(\nabla \mathbf{u}) : \nabla \mathbf{u} = 2\mu |\mathbf{D}(\mathbf{u})|^2 + \nu |\operatorname{div} \mathbf{u}|^2$$

Numerical implementation

$$\text{Stability} \quad \frac{d}{dt} E \leq 0.$$

$$E = Ke + c_v \mathcal{E} + H.$$

$$Ke = \frac{1}{2} \rho \mathbf{u}^2, \quad \mathcal{E} = \rho \theta = p^t, \quad H = \rho \int_1^\rho \frac{p^e(z)}{z^2}$$

$$H'' \geq 0.$$

$$\left(\frac{H}{\rho}\right)' = -\frac{H}{\rho^2} + \frac{H'}{\rho}$$

$$p^e = \rho H' - H.$$

Functional spaces

Piecewise linear Crouzeix-Raviart element for velocity.

$$V_{0,h} \equiv \{\mathbf{v}_h \in L^2(\Omega_h); \quad \mathbf{v}_h|_K \in \mathcal{P}^1(K), \forall K \in \Omega_h; \\ \int_{\Gamma} [\![\mathbf{v}_h]\!] = 0, \forall \Gamma \in \mathcal{E}^{int}; \quad \int_{\Gamma} \mathbf{v}_h = 0, \forall \Gamma \in \mathcal{E}^{ext}\}.$$

Piecewise constant element for density, pressure and temperature

$$Q_h \equiv \{\phi_h \in L^2(\Omega_h); \phi_h|_K \in \mathcal{P}^0(K), K \in \Omega_h\}.$$

- \mathcal{E} edges
- $\mathcal{E}^{ext} = \mathcal{E} \cap \partial\Omega$ exterior edges
- $\mathcal{E}^{int} = \mathcal{E} \setminus \mathcal{E}^{ext}$ interior edges
- K, L element
- $\Gamma = K \cap L$
- $\mathbf{n}_{\Gamma,K}$ be the outer normal, pointing from K to L

Upwind flux

$$\mathcal{F}^{up}(f, \mathbf{u})|_{\Gamma} = \begin{cases} f_K & \text{if } s_{\Gamma,K} \geq 0, \\ f_L & \text{else,} \end{cases}$$

where $s_{\Gamma,K}^n = \mathbf{u}^n|_{\Gamma} \cdot \mathbf{n}_{\Gamma,K} = s_{\Gamma,K}^{n,+} + s_{\Gamma,K}^{n,-}$,

$$s_{\Gamma,K}^{n,+} = \max\{0, s_{\Gamma,K}^n\} \geq 0, \quad s_{\Gamma,K}^{n,-} = \min\{0, s_{\Gamma,K}^n\} \leq 0.$$

$$s_{\Gamma,L}^{n,-} = -s_{\Gamma,K}^{n,+}, \quad s_{\Gamma,K}^{n,-} = -s_{\Gamma,L}^{n,+}, \quad s_{\Gamma,L}^n = -s_{\Gamma,K}^n.$$

Jump

$$[f]_{\Gamma} = f_L - f_K.$$

Average on K

$$\hat{f}_K = \frac{1}{|K|} \int_K f dx.$$

Average on edge

$$\{f\}_{\Gamma} = \frac{1}{2}(f_K + f_L).$$

Implicit nonlinear FV-FE Scheme

Find $\{(\rho_h^{n+1}, \mathbf{u}_h^{n+1}, \theta_h^{n+1})\}_{n=0}^{n_T-1} \subset (Q_h \times V_{0,h} \times Q_h)$ such that

$$\begin{aligned} & \sum_{K \in \Omega_h} \int_K \frac{\rho_h^{n+1} - \rho_h^n}{\Delta t} \phi_h - \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \mathcal{F}^{up}(\rho_h^{n+1}, \mathbf{u}_h^{n+1}) [\![\phi_h]\!] \\ & + h^{\alpha} \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} [\![\rho_h^{n+1}]\!] [\![\phi_h]\!] = 0, \quad (2a) \end{aligned}$$

for any $\phi_h \in Q_h$.

$$0 < \alpha < 1.$$

$$\begin{aligned}
 & \sum_{K \in \Omega_h} \int_K \frac{\mathbf{q}_h^{n+1} - \mathbf{q}_h^n}{\Delta t} \mathbf{v}_h - \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \mathcal{F}^{up}(\mathbf{q}_h^{n+1}, \mathbf{u}_h^{n+1}) [\![\hat{\mathbf{v}}_h]\!] \\
 & - \sum_{K \in \Omega_h} \int_K p_h^{n+1} \operatorname{div}_h \mathbf{v}_h + 2\mu \sum_{K \in \Omega_h} \int_K \mathbf{D}(\mathbf{u}_h^{n+1}) \mathbf{D}(\mathbf{v}_h) + \nu \sum_{K \in \Omega_h} \int_K \operatorname{div}_h \mathbf{u}_h^{n+1} \operatorname{div}_h \mathbf{v}_h \\
 & + h^\alpha \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} [\![\rho_h^{n+1}]\!] \{ \hat{\mathbf{u}}_h^{n+1} \} [\![\hat{\mathbf{v}}_h]\!] + 2\mu \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \frac{1}{h} [\![\mathbf{u}_h^{n+1}]\!] [\![\mathbf{v}_h]\!] = 0,
 \end{aligned} \tag{2b}$$

for any $\mathbf{v}_h \in V_{0,h}$.

\mathbf{q}_h is the momentum, piecewise constant for all $K \in \Omega_h$

$$\mathbf{q}_K = \rho_K \hat{\mathbf{u}}_K, \quad \hat{\mathbf{u}}_K = \frac{1}{|K|} \int_K \mathbf{u}_h$$

$$\begin{aligned}
 & c_v \sum_{K \in \Omega_h} \int_K \frac{\mathcal{E}_h^{n+1} - \mathcal{E}_h^n}{\Delta t} \phi_h - c_v \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \mathcal{F}^{up}(\mathcal{E}_h^{n+1}, \mathbf{u}_h^{n+1}) [\![\phi_h]\!] \\
 & + \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \frac{1}{d_{\Gamma}} [\![\mathcal{K}(\theta_h^{n+1})]\!] [\![\phi_h]\!] = \sum_{K \in \Omega_h} \int_K (2\mu |\mathbf{D}(\mathbf{u}_h^{n+1})|^2 + \nu |\operatorname{div} \mathbf{u}_h^{n+1}|^2) \phi_h \\
 & + 2\mu \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \frac{1}{h} [\![\mathbf{u}_h^{n+1}]\!]^2 - \sum_{K \in \Omega_h} \int_K \mathcal{E}_h^{n+1} \operatorname{div}_h \mathbf{u}_h^{n+1} \phi_h, \quad (2c)
 \end{aligned}$$

for any $\phi_h \in Q_h$.

Iterative solver, linear or nonlinear ?

- Picard

Let $\mathbf{w} = (\rho, \mathbf{u}, \theta)$. Given the solution $\mathbf{w}^{n,\ell}$ ($\mathbf{w}^{n,0} = \mathbf{w}^n$), and solve a linear system to get $\mathbf{w}^{n,\ell+1}$ for $\ell = 1, 2, \dots$, until $\|\mathbf{w}^{n,\ell+1}\| - \|\mathbf{w}^{n,\ell}\| < \text{tol}$. $\mathbf{w}^{n+1} = \mathbf{w}^{n,\ell+1}$.

$$\begin{aligned} & \sum_{K \in \Omega_h} \int_K \frac{\rho_h^{n,\ell+1} - \rho_h^n}{\Delta t} \phi_h + h^\alpha \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_\Gamma [\![\rho_h^{n,\ell+1}]\!] [\!\phi_h]\!] \\ &= \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_\Gamma \mathcal{F}^{up}(\rho_h^{n,\ell}, \mathbf{u}_h^{n,\ell}) [\!\phi_h]\!]. \end{aligned}$$

- inexact-Newton

Jacobian Free Newton Krylov(JFNK)

Positivity preserving

Lemma 1

Suppose that $\rho_h^{n+1} \in Q_h(\Omega_h)$ satisfies (2a), where $\rho_h^n > 0$ in Ω_h and $\mathbf{u}_h^{n+1} \in V_{0,h}(\Omega_h)$. Then

$$\rho_h^{n+1} > 0, \text{ in } \Omega_h. \quad (3)$$

Stability of total energy

Theorem 2

Let $(u_h^{n+1}, \rho_h^{n+1}, \theta_h^{n+1})$ be the solution of the Scheme (2) and the density is initially positive. Then the total energy is dissipative in time

$$E_{\Omega_h}^{n+1} \leq E_{\Omega_h}^n. \quad (4)$$

Stability - - Step 1. density scheme I

$$\begin{aligned} \sum_{K \in \Omega_h} \int_K \frac{\rho_h^{n+1} - \rho_h^n}{\Delta t} \phi_h - \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \mathcal{F}^{up}(\rho_h^{n+1}, \mathbf{u}_h^{n+1}) [\![\phi_h]\!] \\ + h^{\alpha} \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} [\![\rho_h^{n+1}]\!] [\![\phi_h]\!] = 0, \end{aligned}$$

$$\phi_h = -\frac{1}{2} |\hat{\mathbf{u}}_h^{n+1}|^2 \longrightarrow T_{11} + T_{12} + T_{13} = 0.$$

$$T_{11} := -\frac{1}{2\Delta t} \sum_{K \in \Omega_h} |K| (\rho_K^{n+1} - \rho_K^n) |\hat{\mathbf{u}}_K^{n+1}|^2,$$

$$T_{12} := -\frac{1}{2} \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \mathcal{F}^{up}(\rho_h^{n+1}, \mathbf{u}_h^n) |\hat{\mathbf{u}}_K^{n+1}|^2,$$

$$T_{13} := -\frac{1}{2} h^{\alpha} \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} [\![\rho_h^{n+1}]\!] [\![|\hat{\mathbf{u}}_h^{n+1}|^2]\!].$$

Stability - - Step 1. density scheme II

$$\phi_h = H'(\rho_h^{n+1}) \longrightarrow T_{14} + T_{15} \geq 0.$$

$$\begin{aligned} T_{14} &:= \frac{1}{\Delta t} \sum_{K \in \Omega_h} |K| (\rho_K^{n+1} - \rho_K^n) H'(\rho_K^{n+1}) \\ &\geq \frac{1}{\Delta t} \sum_{K \in \Omega_h} |K| \left(H(\rho_K^{n+1}) - H(\rho_K^n) \right) = \frac{1}{\Delta t} (H_{\Omega_h}^{n+1} - H_{\Omega_h}^n), \\ T_{15} &:= \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \mathcal{F}^{up}(\rho_h^{n+1}, \mathbf{u}_h^n) H'(\rho_h^{n+1}). \end{aligned}$$

$$H(\rho^n) = H(\rho^{n+1}) + H'(\rho^{n+1})(\rho^n - \rho^{n+1}) + \frac{H''(\eta_0)}{2}(\rho^n - \rho^{n+1})^2$$

$$\eta_0 \in \text{co}\{\rho^n, \rho^{n+1}\}, \quad \text{co}\{a, b\} = [\min(a, b), \max(a, b)].$$

Stability - - Step 2. momentum scheme

$$\begin{aligned}
& \sum_{K \in \Omega_h} \int_K \frac{\mathbf{q}_h^{n+1} - \mathbf{q}_h^n}{\Delta t} \mathbf{v}_h - \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \mathcal{F}^{up}(\mathbf{q}_h^{n+1}, \mathbf{u}_h^{n+1}) [\![\hat{\mathbf{v}}_h]\!] \\
& - \sum_{K \in \Omega_h} \int_K \rho_h^{n+1} \operatorname{div}_h \mathbf{v}_h + 2\mu \sum_{K \in \Omega_h} \int_K \mathbf{D}(\mathbf{u}_h^{n+1}) \mathbf{D}(\mathbf{v}_h) + \nu \sum_{K \in \Omega_h} \int_K \operatorname{div}_h \mathbf{u}_h^{n+1} \operatorname{div}_h \mathbf{v}_h \\
& + h^\alpha \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} [\![\rho_h^{n+1}]\!] \{ \hat{\mathbf{u}}_h^{n+1} \} [\![\hat{\mathbf{v}}_h]\!] + 2\mu \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \frac{1}{h} [\![\mathbf{u}_h^{n+1}]\!] [\![\mathbf{v}_h]\!] = 0,
\end{aligned}$$

$$\mathbf{v}_h = \mathbf{u}_h^{n+1} \longrightarrow T_{21} + T_{22} + T_{23} + T_{24} + T_{25} + T_{26} = 0.$$

$$T_{21} := \frac{1}{\Delta t} \sum_{K \in \Omega_h} |K| (\rho_K^{n+1} |\hat{\mathbf{u}}_K^{n+1}|^2 - \rho_K^n \hat{\mathbf{u}}_K^n \hat{\mathbf{u}}_K^{n+1}),$$

$$T_{22} := \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \mathcal{F}^{up}(\mathbf{q}_h, \mathbf{u}_h)^{n+1} \hat{\mathbf{u}}_K^{n+1}, \quad T_{23} := - \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} (p^e)_K^{n+1} s_{\Gamma}^{n+1}$$

$$T_{24} := \sum_{K \in \Omega_h} \int_K 2\mu |\mathbf{D}(\mathbf{u}_h^{n+1})|^2 + \nu |\operatorname{div} \mathbf{u}_h^{n+1}|^2 + 2\mu \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \frac{1}{h} [\![\mathbf{u}_h^{n+1}]\!]^2,$$

$$T_{25} := h^\alpha \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} [\![\rho_h^{n+1}]\!] \{ \hat{\mathbf{u}}_h^{n+1} \} [\![\hat{\mathbf{u}}_h^{n+1}]\!], \quad T_{26} := - \sum_{K \in \Omega_h} \int_K (p^t)_h^{n+1} \operatorname{div} \mathbf{u}_h^{n+1}$$

Term $T_{11} + T_{21}$

$$\begin{aligned} & T_{11} + T_{21} \\ &= -\frac{1}{2\Delta t} \sum_{K \in \Omega_h} |K| (\rho_K^{n+1} - \rho_K^n) |\hat{\mathbf{u}}_K^{n+1}|^2 + \frac{1}{\Delta t} \sum_{K \in \Omega_h} |K| (\rho_K^{n+1} |\hat{\mathbf{u}}_K^{n+1}|^2 - \rho_K^n \hat{\mathbf{u}}_K^n \hat{\mathbf{u}}_K^{n+1}) \\ &= \frac{1}{\Delta t} \sum_{K \in \Omega_h} |K| \left(\frac{1}{2} \rho_K^{n+1} |\hat{\mathbf{u}}_K^{n+1}|^2 - \frac{1}{2} \rho_K^n |\hat{\mathbf{u}}_K^n|^2 + \frac{1}{2} \rho_K^n (\hat{\mathbf{u}}_K^{n+1} - \hat{\mathbf{u}}_K^n)^2 \right) \\ &\geq \frac{1}{\Delta t} (K e_{\Omega_h}^{n+1} - K e_{\Omega_h}^n). \end{aligned}$$

Stability - - Upwind flux I

Term $T_{12} + T_{22}$

$$T_{12} := -\frac{1}{2} \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \mathcal{F}^{up}(\rho_h^{n+1}, \mathbf{u}_h^n) |\hat{\mathbf{u}}_K^{n+1}|^2$$

Contribution of T_{12} at an arbitrary edge $\Gamma = K \cap L$

$$\begin{aligned} & -\frac{1}{2} \rho_K^{n+1} s_{\Gamma,K}^{n,+} |\hat{\mathbf{u}}_K^{n+1}|^2 - \frac{1}{2} \rho_L^{n+1} s_{\Gamma,K}^{n,-} |\hat{\mathbf{u}}_K^{n+1}|^2 - \frac{1}{2} \rho_L^{n+1} s_{\Gamma,L}^{n,+} |\hat{\mathbf{u}}_L^{n+1}|^2 - \frac{1}{2} \rho_K^{n+1} s_{\Gamma,L}^{n,-} |\hat{\mathbf{u}}_L^{n+1}|^2 \\ &= -\frac{1}{2} \rho_K^{n+1} s_{\Gamma,K}^{n,+} (|\hat{\mathbf{u}}_K^{n+1}|^2 - |\hat{\mathbf{u}}_L^{n+1}|^2) - \frac{1}{2} \rho_L^{n+1} s_{\Gamma,L}^{n,+} (|\hat{\mathbf{u}}_L^{n+1}|^2 - |\hat{\mathbf{u}}_K^{n+1}|^2). \end{aligned}$$

Contribution of T_{22} at an arbitrary edge $\Gamma = K \cap L$

$$(\rho_K \hat{\mathbf{u}}_K)^{n+1} s_{\Gamma,K}^{n,+} (\hat{\mathbf{u}}_K^{n+1} - \hat{\mathbf{u}}_L^{n+1}) + (\rho_L \hat{\mathbf{u}}_L)^{n+1} s_{\Gamma,L}^{n,+} (\hat{\mathbf{u}}_L^{n+1} - \hat{\mathbf{u}}_K^{n+1}).$$

Term $T_{12} + T_{22}$

$$\begin{aligned} T_{12} + T_{22} &= \sum_{\Gamma \in \mathcal{E}^{int}} \int_{\Gamma} s_{\Gamma,K}^{n,+} \rho_K^{n+1} \left(\frac{1}{2} |\hat{\mathbf{u}}_L^{n+1}|^2 + \frac{1}{2} |\hat{\mathbf{u}}_K^{n+1}|^2 - \hat{\mathbf{u}}_L^{n+1} \hat{\mathbf{u}}_K^{n+1} \right) \\ &\quad + \sum_{\Gamma \in \mathcal{E}^{int}} \int_{\Gamma} s_{\Gamma,L}^{n,+} \rho_L^{n+1} \left(\frac{1}{2} |\hat{\mathbf{u}}_L^{n+1}|^2 + \frac{1}{2} |\hat{\mathbf{u}}_K^{n+1}|^2 - \hat{\mathbf{u}}_L^{n+1} \hat{\mathbf{u}}_K^{n+1} \right) \\ &= \frac{1}{2} \sum_{\Gamma \in \mathcal{E}^{int}} \int_{\Gamma} |\mathbf{u}_{\Gamma}^n \cdot \mathbf{n}| \rho^{n+1,up} (\hat{\mathbf{u}}_K^{n+1} - \hat{\mathbf{u}}_L^{n+1})^2 \\ &\geq 0. \end{aligned}$$

Term $T_{13} + T_{25}$

$$T_{13} + T_{25}$$

$$\begin{aligned} &= -\frac{1}{2} h^\alpha \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} [\![\rho_h^{n+1}]\!] [\!|[\hat{\mathbf{u}}_h^{n+1}]|^2]\!] + h^\alpha \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} [\![\rho_h^{n+1}]\!] \{[\hat{\mathbf{u}}_h^{n+1}]\} [\![\hat{\mathbf{u}}_h^{n+1}]\!] \\ &= 0, \end{aligned}$$

as for all $\Gamma = K \cap L$

$$\begin{aligned} \{[\hat{\mathbf{u}}_h^{n+1}]\}_{\Gamma} [\![\hat{\mathbf{u}}_h^{n+1}]\!]_{\Gamma} &= \frac{\hat{\mathbf{u}}_L^{n+1} + \hat{\mathbf{u}}_K^{n+1}}{2} (\hat{\mathbf{u}}_L^{n+1} - \hat{\mathbf{u}}_K^{n+1}) = \frac{|\hat{\mathbf{u}}_L^{n+1}|^2 - |\hat{\mathbf{u}}_K^{n+1}|^2}{2} \\ &= \frac{1}{2} [\!|[\hat{\mathbf{u}}_h^{n+1}]|^2]\!]_{\Gamma}. \end{aligned}$$

Term $T_{23} + T_{15}$

$$\begin{aligned}
 T_{23} + T_{15} &= \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} -(p^e)_K^{n+1} s_{\Gamma}^{n+1} + \mathcal{F}^{up}(\rho_h^{n+1}, \mathbf{u}_h^n) H'(\rho_h^{n+1}) \\
 &= \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} s_{\Gamma}^{n+1} \left(H(\rho_K^{n+1}) - \rho_K^{n+1} H'(\rho_K^{n+1}) + \rho^{n+1, up} H'(\rho_K^{n+1}) \right) \\
 &= \sum_{\Gamma \in \mathcal{E}^{int}} \int_{\Gamma=K \cap L} s_{\Gamma, K}^{n+1, +} \left(H(\rho_K^{n+1}) - H(\rho_L^{n+1}) + H'(\rho_L^{n+1})(\rho_L^{n+1} - \rho_K^{n+1}) \right) \\
 &\quad + \sum_{\Gamma \in \mathcal{E}^{int}} \int_{\Gamma=K \cap L} s_{\Gamma, L}^{n+1, +} \left(H(\rho_L^{n+1}) - H(\rho_K^{n+1}) + H'(\rho_K^{n+1})(\rho_K^{n+1} - \rho_L^{n+1}) \right) \\
 &= \sum_{\Gamma \in \mathcal{E}^{int}} \int_{\Gamma=K \cap L} \left(s_{\Gamma, K}^{n+1, +} \frac{H''(\eta_1)}{2} + s_{\Gamma, L}^{n+1, +} \frac{H''(\eta_2)}{2} \right) (\rho_K^{n+1} - \rho_L^{n+1})^2 \geq 0
 \end{aligned}$$

$$\eta_1, \eta_2 \in \text{co}\{\rho_K^{n+1}, \rho_L^{n+1}\}$$

Stability - - Kinetic and elastic energy

$$0 = \sum_{i=1}^3 T_{1i} + \sum_{i=1}^6 T_{2i} \geq T_{24} + T_{26} + \frac{1}{\Delta t} (K e_{\Omega_h}^{n+1} - K e_{\Omega_h}^n + H_{\Omega_h}^{n+1} - H_{\Omega_h}^n)$$

$$T_{24} := \sum_{K \in \Omega_h} \int_K 2\mu |\mathbf{D}(\mathbf{u}_h^{n+1})|^2 + \nu |\operatorname{div} \mathbf{u}_h^{n+1}|^2 + 2\mu \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \frac{1}{h} [\![\mathbf{u}_h^{n+1}]\!]^2,$$

$$T_{26} := - \sum_{K \in \Omega_h} \int_K \mathcal{E}_h^{n+1} \operatorname{div} \mathbf{u}_h^{n+1}$$

$$E_h = K e_h + H_h + c_v \mathcal{E}_h$$

$$T_{24} + T_{26} + \frac{c_v}{\Delta t} \sum_{K \in \Omega_h} \int_K (\mathcal{E}_h^{n+1} - \mathcal{E}_h^n) = ?$$

Stability - - Step 3. $\phi_h = 1$ for temperature scheme

$$\begin{aligned}
 & c_v \sum_{K \in \Omega_h} \int_K \frac{\mathcal{E}_h^{n+1} - \mathcal{E}_h^n}{\Delta t} \phi_h - c_v \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \mathcal{F}^{up}(\mathcal{E}_h^{n+1}, \mathbf{u}_h^{n+1}) [\![\phi_h]\!] \\
 & + \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \frac{1}{d_{\Gamma}} [\![\mathcal{K}(\theta_h^{n+1})]\!] [\![\phi_h]\!] = \sum_{K \in \Omega_h} \int_K (2\mu |\mathbf{D}(\mathbf{u}_h^{n+1})|^2 + \nu |\operatorname{div} \mathbf{u}_h^{n+1}|^2) \phi_h \\
 & + 2\mu \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \frac{1}{h} [\![\mathbf{u}_h^{n+1}]\!]^2 - \sum_{K \in \Omega_h} \int_K \mathcal{E}_h^{n+1} \operatorname{div}_h \mathbf{u}_h^{n+1} \phi_h,
 \end{aligned}$$

$$T_{31} + T_{32} + T_{33} = 0.$$

$$T_{31} := \frac{c_v}{\Delta t} \sum_{K \in \Omega_h} \int_K (\mathcal{E}_h^{n+1} - \mathcal{E}_h^n) = \frac{c_v}{\Delta t} (\mathcal{E}_{\Omega_h}^{n+1} - \mathcal{E}_{\Omega_h}^n),$$

$$T_{32} := - \sum_{K \in \Omega_h} \int_K (2\mu |\mathbf{D}(\mathbf{u}_h^{n+1})|^2 + \nu |\operatorname{div} \mathbf{u}_h^{n+1}|^2) - 2\mu \sum_{K \in \Omega_h} \sum_{\Gamma \in \partial K} \int_{\Gamma} \frac{1}{h} [\![\mathbf{u}_h^{n+1}]\!]^2,$$

$$T_{33} := \sum_{K \in \Omega_h} \int_K \mathcal{E}_h^{n+1} \operatorname{div}_h \mathbf{u}_h^{n+1}.$$

$$\mu = \nu = 1.0, a = 1.0, b = 1.0, \gamma = 3.0, c_v = 1.4, \alpha = 0.83.$$

Time step is

$$\Delta t = \text{CFL} \frac{\min(h_K)}{\max(|U|) + c}, \quad c = \sqrt{\frac{\partial p}{\partial \rho}} = \sqrt{a\gamma\rho^{\gamma-1} + b + \theta}, \quad \text{CFL} = 0.6.$$

Boundary conditions is set as periodic.

density

temperature

velocity

Table : Convergence results of Poiseulle flow

h	$\ \rho\ _{L^\infty(L^\gamma)}$	EOC	$\ \rho\ _{L^1(L^1)}$	EOC	$\ u\ _{L^2(L^2)}$	EOC	$\ u\ _{L^2(H^1)}$	EOC	$\ \theta\ _{L^2(L^2)}$	EOC
1/8	1.87e-02	—	1.00e-02	—	1.18e-01	—	9.70e-01	—	1.36e-02	—
1/16	8.07e-03	1.21	3.74e-03	1.42	3.33e-02	1.83	4.58e-01	1.08	4.79e-03	1.51
1/32	3.89e-03	1.05	1.70e-03	1.14	1.05e-02	1.67	2.25e-01	1.03	2.16e-03	1.15
1/64	2.03e-03	0.94	9.47e-04	0.84	3.84e-03	1.45	1.12e-01	1.01	1.06e-03	1.03
1/128	1.07e-03	0.92	5.60e-04	0.76	1.58e-03	1.28	5.61e-02	1.00	5.24e-04	1.02



