

# Asymptotic preserving error estimates for numerical solutions of compressible Navier-Stokes equations in the low Mach number regime

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Let  $\rho, \mathbf{u}, p$  be the fluid density, velocity, pressure.

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0. \quad (1a)$$

$$(\rho \mathbf{u})_t + \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) + \frac{1}{\varepsilon^2} \nabla p = \operatorname{div} \mathbb{S}(\nabla \mathbf{u}). \quad (1b)$$

$$\rho(\mathbf{x}, 0) = \rho_0 > 0, \quad \mathbf{u}(\mathbf{x}, 0) = \mathbf{u}_0. \quad (1c)$$

$$\mathbb{S} = \mu(\nabla \mathbf{u} + \nabla \mathbf{u}^T - \frac{2}{3} \operatorname{div} \mathbf{u} \mathbf{1}) + \eta \operatorname{div} \mathbf{u}, \quad \mu > 0, \quad \eta \geq 0.$$

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$$p \in C^2(0, \infty) \cap C^1[0, \infty), p(0) = 0, p'(\rho) > 0, \text{ for all } \rho > 0, \quad (1d)$$
$$\lim_{\rho \rightarrow \infty} \frac{p'(\rho)}{\rho^{\gamma-1}} = p_\infty > 0, \quad \lim_{\rho \rightarrow 0^+} \frac{p'(\rho)}{\rho^\alpha} = p_0 > 0, \quad \gamma \geq 1, \alpha \leq 1.$$

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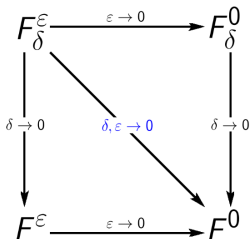
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$$p = \rho^\gamma$$

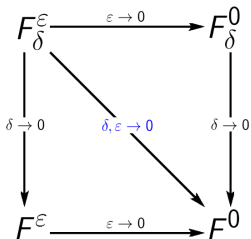
## Asymptotic preserving<sup>[3]</sup>



$$\bar{\rho}(\partial_t \mathbf{V} + \mathbf{V} \cdot \nabla_x \mathbf{V}) + \nabla_x \pi = \mu \Delta \mathbf{V}, \operatorname{div}_x \mathbf{V} = 0, \bar{\rho} > 0, \quad (2a)$$

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$$\partial_t^\ell \mathbf{V} \in C^\ell([0, T]; W^{k-\ell, 2}(\Omega; \mathbb{R}^3)), \ell = 0, 1, 2,$$

$$\partial_t^j \pi \in C^j([0, T]; W^{k-1-j}(\Omega)), j = 0, 1, k \geq 4,$$

$$\mathbf{V}_0 \in W^{k, 2}(\Omega; \mathbb{R}^3), \operatorname{div}_x \mathbf{V}_0 = 0.$$

Piecewise linear Crouzeix-Raviart element for velocity.

$$V_{0,h} \equiv \{ \mathbf{v}_h \in L^2(\Omega_h); \mathbf{v}_h|_K \in \mathcal{P}^1(K), \forall K \in \Omega_h; \\ \int_{\Gamma} [[\mathbf{v}_h]] = 0, \forall \Gamma \in \mathcal{E}_{int}; \int_{\Gamma} \mathbf{v}_h = 0, \forall \Gamma \in \partial\mathcal{T} \}.$$

Piecewise constant element for density, pressure and temperature

$$Q_h \equiv \{ \phi_h \in L^2(\Omega_h); \phi_h|_K \in \mathcal{P}^0(K), K \in \Omega_h \}.$$

- Element  $K, L$ , interface  $\Gamma = K \cap L$ .
- $\mathbf{n}_{\Gamma,K}$  be the outer normal, pointing from  $K$  to  $L$

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**Upwind flux**

$$\mathcal{F}^{up}(f, \mathbf{u})|_{\Gamma} = \begin{cases} f_K & \text{if } s_{\Gamma,K} \geq 0, \\ f_L & \text{else,} \end{cases} \quad s_{\Gamma,K}^n = \mathbf{u}^n|_{\Gamma} \cdot \mathbf{n}_{\Gamma,K} = s_{\Gamma,K}^{n,+} + s_{\Gamma,K}^{n,-}$$

**Jump, average on element & edge**

$$\llbracket f \rrbracket_{\Gamma} = f_L - f_K. \quad \hat{f}_K = \frac{1}{|K|} \int_K f dx. \quad \{f\}_{\Gamma} = \frac{1}{2}(f_K + f_L).$$



## A convergent scheme (Karper [4])

Find  $\{\rho_h^n, \mathbf{u}_h^n\}_{n=1}^{n_T} \subset (Q_h \times V_{0,h})$  such that for any  $(\phi_h, \mathbf{v}_h) \in (Q_h \times V_{0,h})$

$$\sum_{K \in \mathcal{T}} |K| \frac{\rho_K^n - \rho_K^{n-1}}{\Delta t} \phi_h - \sum_{K \in \mathcal{T}} \sum_{\Gamma \in \partial K} |\sigma| \rho_\sigma^{n,up} (\mathbf{u}_\sigma^n \cdot \mathbf{n}_{\sigma,K}) \phi_h = 0, \quad (3a)$$

$$\begin{aligned} \sum_{K \in \mathcal{T}} \int_K \frac{\rho_h^n \hat{\mathbf{u}}_h^n - \rho_h^{n-1} \hat{\mathbf{u}}_h^{n-1}}{\Delta t} \mathbf{v}_h - \sum_{K \in \mathcal{T}} \sum_{\Gamma \in \partial K} |\sigma| (\rho \hat{\mathbf{u}})_\sigma^{n,up} (\mathbf{u}_\sigma^n \cdot \mathbf{n}_{\sigma,K}) \cdot \mathbf{v}_h \\ - \frac{1}{\varepsilon^2} \sum_{K \in \mathcal{T}} \rho_h^n \sum_{\Gamma \in \partial K} \operatorname{div}_h \mathbf{v}_h + \mu \sum_{K \in \mathcal{T}} \int_K \nabla_h \mathbf{u}_h^n : \nabla_h \mathbf{v}_h \\ + \left(\frac{\mu}{3} + \eta\right) \sum_{K \in \mathcal{T}} \int_K \operatorname{div}_h \mathbf{u}_h^n \operatorname{div}_h \mathbf{v}_h = 0. \quad (3b) \end{aligned}$$

Relative energy inequality (Feireisl et al.<sup>[1]</sup>)

$$\mathcal{E}(\rho, \mathbf{u}|z, \mathbf{V})|_0^\tau + \int_0^\tau \int_\Omega (\mathbb{S}(\nabla \mathbf{u}) - \mathbb{S}(\nabla \mathbf{V})) : (\nabla \mathbf{u} - \nabla \mathbf{V}) dx dt \leq \int_0^\tau \mathcal{R}(\rho, \mathbf{u}, z, \mathbf{V}) dt$$

$$\mathcal{E}(\rho, \mathbf{u}|z, \mathbf{V}) = \int_\Omega (\rho |\mathbf{u} - \mathbf{V}|^2 + E(\rho|z)) dx$$

$$E(\rho|z) = H(\rho) - H'(z)(\rho - z) - H(z), H(\rho) = \rho \int_1^\rho \frac{\rho(s)}{s^2} ds$$

$H(\rho)$  convex for  $\rho \in (0, \infty)$ ,  $E(\rho, z) \geq 0$  and  $E(\rho, z) = 0 \Leftrightarrow \rho = z$ .

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$$\mathcal{E}_\varepsilon(\rho, \mathbf{u}|z, \mathbf{V}) = \int_\Omega (\rho |\mathbf{u} - \mathbf{V}|^2 + \frac{1}{\varepsilon^2} E(\rho|z)) dx$$

## Theorem

Let  $p$  satisfy (1d) with  $\gamma \geq 3/2$ . Let  $\{\rho^n, \mathbf{u}^n\}_{0 \leq n \leq N}$  be a family of numerical solutions constructed by the scheme (3) and the mesh be regular, and initial data  $(\rho^0, \mathbf{u}^0)$  obey

$$\mathcal{E}_\varepsilon(\rho_\varepsilon^0, \mathbf{u}_\varepsilon^0 | \bar{\rho}, \mathbf{V}(0)) \leq E_0 < \infty, M_0/2 \leq \int_{\mathcal{T}} \rho_\varepsilon^0 dx \leq 2M_0, M_0 = \bar{\rho} |\mathcal{T}| \quad (4)$$

Moreover, suppose that  $[\Pi, \mathbf{V}]$  is a classical solution to (2) emanating from the initial data  $\mathbf{V}_0 \in W^{k,2}(\Omega; R^3)$ ,  $\operatorname{div} \mathbf{V}_0 = 0$ ,  $k \geq 4$ . Then there exists a positive number independent of  $h, \Delta t, \varepsilon$  such that

$$\begin{aligned} \sup_{1 \leq n \leq N} \mathcal{E}_\varepsilon(\rho_\varepsilon^n, \mathbf{u}_\varepsilon^n | \bar{\rho}, \mathbf{V}(t_n, \cdot)) + \Delta t \sum_{1 \leq n \leq N} \int_{\mathcal{T}} |\nabla_h \mathbf{u}^n - \nabla_x \mathbf{V}(t_n, \cdot)|^2 \\ \leq c \left( \sqrt{\Delta t} + h^a + \varepsilon + \mathcal{E}_\varepsilon(\rho_\varepsilon^0, \mathbf{u}_\varepsilon^0 | \bar{\rho}, \mathbf{V}_0) \right) \quad (5) \end{aligned}$$

where  $a = \min\{\frac{2\gamma-3}{\gamma}, 1\}$ .

- Step 1: uniform estimates (Karper<sup>[4]</sup>, Gallouët et al.<sup>[2]</sup>)

$$\boxed{(3b)|_{\mathbf{v}_h=\mathbf{u}_K^n} + (3a)|_{\phi_h=\frac{-|\hat{\mathbf{u}}_K^n|^2}{2}} + (3a)|_{\phi_h=H'(\rho_K^n)}}$$

sum up for each time step  $\implies$

$$\begin{aligned} & \sum_{K \in \mathcal{T}} |K| \left( \frac{1}{2} \rho_K^m |\hat{\mathbf{u}}_K^m| + \frac{1}{\varepsilon^2} E(\rho_K^m | \bar{\rho}) \right) - \sum_{K \in \mathcal{T}} |K| \left( \frac{1}{2} \rho_K^0 |\hat{\mathbf{u}}_K^0| + \frac{1}{\varepsilon^2} E(\rho_K^0 | \bar{\rho}) \right) \\ & + \Delta t \sum_{n=1}^m \sum_{K \in \mathcal{T}} \int_K \left( \mu \int_K |\nabla_x \mathbf{u}^n|^2 + \left( \frac{\mu}{3} + \eta \right) \int_K |\operatorname{div}_x \mathbf{u}^n|^2 \right) + D_i^m = 0 \end{aligned}$$

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Remark: not valid for linear scheme yet

- Step 2: Relative energy inequality (Gallouët et al.<sup>[2]</sup>)

$$(3b)|_{\mathbf{v}_h=\mathbf{u}_K^n} + (3a)|_{\phi_h=\frac{-|\hat{\mathbf{u}}_K^n|^2}{2}} + (3a)|_{\phi_h=H'(\rho_K^n)}$$

$$- (3b)|_{\mathbf{v}_h=\mathbf{v}_K^n} + (3a)|_{\phi_h=\frac{-|\hat{\mathbf{v}}_K^n|^2}{2}} + (3a)|_{\phi_h=H'(z_K^{n-1})}$$

sum up for each time step  $\implies$

$$\int_{\Omega} \left( \rho^m |\hat{\mathbf{u}}^m - \hat{\mathbf{v}}_h^m|^2 + \frac{1}{\varepsilon^2} E(\rho^m | \bar{\rho}) \right) dx - \int_{\Omega} \left( \rho^0 |\hat{\mathbf{u}}^0 - \hat{\mathbf{v}}_h^0|^2 + \frac{1}{\varepsilon^2} E(\rho^0 | \bar{\rho}) \right) dx$$

$$+ \Delta t \sum_{n=1}^m \sum_{K \in \Omega} \left( \int_K |\nabla_x(\mathbf{u}^n - \mathbf{v}_h^n)|^2 + \left( \frac{\mu}{3} + \eta \right) \int_K |\operatorname{div}(\mathbf{u}^n - \mathbf{v}_h^n)|^2 \right) \leq \sum_{i=1}^3 S_i(\rho_h) + R_1^m + G^m$$

$$|R_1^m| \leq c(\sqrt{\Delta t} + h^a), \quad |G^m| \leq c\Delta t \sum_{n=1}^m \mathcal{E}_{\varepsilon}(\rho^n, \mathbf{u}^n | \bar{\rho}, \mathbf{V}^n).$$

- Step 3: Control of  $S_i$

$$S_1 = \Delta t \sum_{n=1}^m \sum_{K \in \mathcal{T}} \int_K \mu \nabla_x \mathbf{v}_h^n : \nabla_x (\mathbf{v}_h^n - \mathbf{u}^n) dx,$$

$$S_2 = \Delta t \sum_{n=1}^m \sum_{K \in \mathcal{T}} |K| \bar{\rho} \frac{\mathbf{v}_{h,K}^n - \mathbf{v}_{h,K}^{n-1}}{\Delta t} \cdot (\mathbf{v}_{h,K}^n - \mathbf{u}_K^n) dx,$$

$$S_3 = \Delta t \sum_{n=1}^m \sum_{K \in \mathcal{T}} \sum_{\Gamma \in \partial K} |\sigma| \bar{\rho} (\hat{\mathbf{v}}_{h,\sigma}^{n,up} - \hat{\mathbf{u}}_{\sigma}^{n,up}) \cdot (\mathbf{v}_{h,\sigma}^n - \mathbf{v}_{h,K}^n) \hat{\mathbf{v}}_{h,\sigma}^{n,up} \cdot \mathbf{n}_{\sigma,K} dx.$$

Navier-Stokes (2) test with  $(\mathbf{V}^n - \mathbf{u}^n) \implies$

$$\sum_{i=1}^3 S_i + R_2^m = 0, \quad |R_2^m| \leq c(h^b + \Delta t + \varepsilon), \quad b = \min\left\{\frac{5\gamma - 6}{2\gamma}, 1\right\}.$$



# Error Estimates

$$\begin{aligned} & \sup_{1 \leq n \leq N} \mathcal{E}_\varepsilon \left( \rho_\varepsilon^n, \mathbf{u}_\varepsilon^n | \bar{\rho}, \mathbf{V}(t_n, \cdot) \right) + \Delta t \sum_{1 \leq n \leq N} \int_{\mathcal{T}} |\nabla_h \mathbf{u}^n - \nabla_x \mathbf{V}(t_n, \cdot)|^2 \\ & \leq c \left( \sqrt{\Delta t} + h^a + \varepsilon + \mathcal{E}_\varepsilon(\rho_\varepsilon^0, \mathbf{u}_\varepsilon^0 | \bar{\rho}, \mathbf{V}_0) \right), \quad a = \min \left\{ \frac{2\gamma - 3}{\gamma}, 1 \right\} \end{aligned}$$

$$\begin{aligned} & \sup_{1 \leq n \leq N} \mathcal{E}_\varepsilon \left( \rho_\varepsilon^n, \mathbf{u}_\varepsilon^n | \bar{\rho}, \mathbf{V}(t_n, \cdot) \right) + \Delta t \sum_{1 \leq n \leq N} \int_{\mathcal{T}} |\nabla_h \mathbf{u}^n - \nabla_x \mathbf{V}(t_n, \cdot)|^2 \\ & \leq c \left( \sqrt{\Delta t} + h^a + \varepsilon + \mathcal{E}_\varepsilon(\rho_\varepsilon^0, \mathbf{u}_\varepsilon^0 | \bar{\rho}, \mathbf{V}_0) \right), \quad a = \min \left\{ \frac{2\gamma - 3}{\gamma}, 1 \right\} \end{aligned}$$

## Remark

*Theorem holds also in the 2D case for any  $0 \leq a < \frac{2\gamma-2}{\gamma}$  if  $\gamma \in (1, 2]$  and  $a = 1$  if  $\gamma > 2$ . Note that in this case the limit system (2) admits global-in-time smooth solutions as long as the initial data are regular.*

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$$\begin{aligned} p &= \rho^\gamma, z = (1 + \varepsilon^2 \Pi)^{1/\gamma} \\ e_\varepsilon &= \sup_{1 \leq n \leq N} \mathcal{E}_\varepsilon \left( \rho_\varepsilon^n, \mathbf{u}_\varepsilon^n | \bar{\rho}, \mathbf{V}(t_n, \cdot) \right) \\ e_{\nabla \mathbf{u}} &= \left\| \nabla_h \mathbf{u}^n - \nabla_x \mathbf{V}(t_n, \cdot) \right\|_{L^2(L^2)} \\ e_{\mathbf{u}} &= \left\| \mathbf{u}^n - \mathbf{V}(t_n, \cdot) \right\|_{L^2(L^2)} \\ e_\rho &= \left\| \rho - z \right\|_{L^2(L^2)}, \\ e_p &= \left\| p - 1 - \varepsilon^2 \Pi \right\|_{L^2(L^2)}. \end{aligned}$$

# Experiment - I

$$\begin{aligned}u_1(x, y, 0) &= \sin^2(\pi x) \sin(2\pi y) \\u_2(x, y, 0) &= -\sin(2\pi x) \sin^2(\pi y) \\ \rho(x, y, 0) &= 1 - \frac{\varepsilon^2}{2} \tanh(y - 0.5)\end{aligned}$$

**Table :** Error with respect to the numerical solution of NS<sup>[5]</sup>

$h$	$e_\varepsilon$	EOC	$e_{\nabla \mathbf{u}}$	EOC	$e_{\mathbf{u}}$	EOC	$e_\rho$	EOC	$e_p$	EOC
1/8	1.12e-03	-	4.91e-01	-	1.82e-03	-	3.65e-04	-	5.11e-04	-
1/16	3.74e-04	1.58	2.55e-01	0.95	1.18e-03	0.63	7.90e-05	2.21	1.11e-04	2.20
1/32	1.09e-04	1.78	2.29e-01	0.16	7.87e-04	0.58	1.50e-05	2.40	2.09e-05	2.41
1/64	1.91e-05	2.51	1.26e-01	0.86	3.24e-04	1.28	3.31e-06	2.18	4.64e-06	2.17
1/128	4.50e-06	2.09	4.41e-02	1.51	1.41e-04	1.20	8.73e-07	1.92	1.22e-06	1.93
1/256	1.14e-06	1.98	1.54e-02	1.52	6.32e-05	1.16	2.06e-07	2.08	2.88e-07	2.08

**(a)**  $\varepsilon = h, \mu = 0.01$

$h$	$e_\varepsilon$	EOC	$e_{\nabla \mathbf{u}}$	EOC	$e_{\mathbf{u}}$	EOC	$e_\rho$	EOC	$e_p$	EOC
1/8	5.42e-03	-	4.66e-01	-	2.99e-03	-	8.81e-04	-	1.23e-03	-
1/16	1.34e-03	2.02	1.73e-01	1.43	2.17e-03	0.46	1.79e-04	2.30	2.51e-04	2.29
1/32	3.66e-04	1.87	6.58e-02	1.39	1.17e-03	0.89	3.78e-05	2.24	5.30e-05	2.24
1/64	1.06e-04	1.79	2.37e-02	1.47	6.13e-04	0.93	8.21e-06	2.20	1.15e-05	2.20
1/128	2.96e-05	1.84	8.26e-03	1.52	2.78e-04	1.14	1.81e-06	2.18	2.54e-06	2.18
1/256	7.96e-06	1.89	2.97e-03	1.48	1.31e-04	1.09	4.09e-07	2.15	5.73e-07	2.15

**(b)**  $\varepsilon = h, \mu = 1$

## Exact solution of unsteady Taylor vortex

$$\begin{aligned}
 V_1(x, y, t) &= \sin(2\pi x) \cos(2\pi y) e^{-8\pi^2 \mu t} \\
 V_2(x, y, t) &= -\cos(2\pi x) \sin(2\pi y) e^{-8\pi^2 \mu t} \\
 \Pi(x, y, t) &= \frac{1}{4} (\cos(4\pi x) + \cos(4\pi y)) e^{-16\pi^2 \mu t}
 \end{aligned}
 \quad \left| \quad
 \begin{aligned}
 \rho(x, y, 0) &= 1 + \varepsilon^2 \Pi(x, y, 0) \\
 u_1(x, y, 0) &= V_1(x, y, 0) \\
 u_2(x, y, 0) &= V_2(x, y, 0)
 \end{aligned}$$

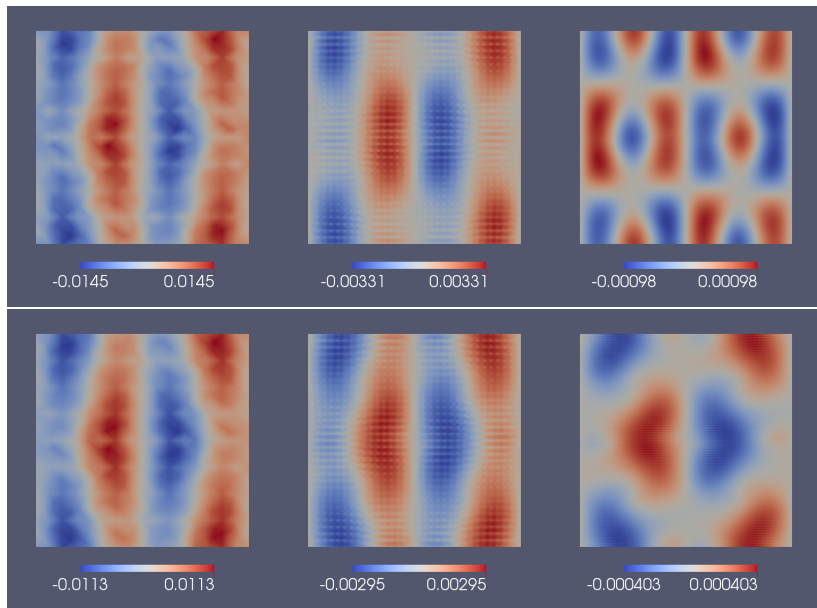
$h$	$e_{\mathcal{E}}$	EOC	$e_{\nabla \mathbf{u}}$	EOC	$e_{\mathbf{u}}$	EOC	$e_p$	EOC	$e_p$	EOC
1/8	3.60e-02	-	3.57e-01	-	7.56e-03	-	3.66e-04	-	4.22e-04	-
1/16	3.04e-03	3.57	1.94e-01	0.88	2.35e-03	1.69	8.67e-05	2.08	9.08e-05	2.22
1/32	2.98e-04	3.35	1.30e-01	0.58	9.44e-04	1.32	1.92e-05	2.17	1.76e-05	2.37
1/64	5.26e-05	2.50	8.35e-02	0.64	3.25e-04	1.54	4.36e-06	2.14	3.45e-06	2.35
1/128	1.46e-05	1.85	4.46e-02	0.90	1.11e-04	1.55	1.08e-06	2.01	1.06e-06	1.70
1/256	3.88e-06	1.91	2.22e-02	1.01	4.05e-05	1.45	2.65e-07	2.03	2.34e-07	2.18

(a)  $\varepsilon = h, \mu = 0.01, \gamma = 3$

$h$	$e_{\mathcal{E}}$	EOC	$e_{\nabla \mathbf{u}}$	EOC	$e_{\mathbf{u}}$	EOC	$e_p$	EOC	$e_p$	EOC
1/8	1.54e-03	-	4.38e-02	-	2.20e-03	-	1.26e-04	-	7.06e-04	-
1/16	3.63e-04	2.08	2.14e-02	1.03	1.04e-03	1.08	3.02e-05	2.06	1.20e-04	2.56
1/32	1.18e-04	1.62	9.99e-03	1.10	4.29e-04	1.28	1.03e-05	1.55	2.41e-05	2.32
1/64	3.95e-05	1.58	4.84e-03	1.05	1.92e-04	1.16	2.70e-06	1.93	5.22e-06	2.21
1/128	1.22e-05	1.69	2.40e-03	1.01	9.26e-05	1.05	6.89e-07	1.97	1.25e-06	2.06
1/256	3.45e-06	1.82	1.20e-03	1.00	4.59e-05	1.01	1.74e-07	1.99	3.10e-07	2.01

(b)  $\varepsilon = h, \mu = 1, \gamma = 3$

Error of velocity component  $u_1 - V_1$  for different mesh sizes and Mach numbers;  $h = \frac{1}{16}, \frac{1}{64}, \frac{1}{256}$  (left to right),  $\varepsilon = 0.1, 0.001$  (top to bottom).



For more details, see Feireisl et al. 2016.

[www.math.cas.cz/fichier/preprints/IM\\_20160916140850\\_71.pdf](http://www.math.cas.cz/fichier/preprints/IM_20160916140850_71.pdf)

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# Thank you for your attention!